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Quantitative Computational Fluid Dynamic Analyses of Particle Deposition on a Transonic Axial Compressor Blade—Part II: Impact Kinematics and Particle Sticking Analysis

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In heavy-duty gas turbines, the microparticles that are not captured by the air filtration system can cause fouling and, consequently, a performance drop of the compressor. This paper presents three-dimensional numerical simulations of the microparticle ingestion ($0\ \mu\text{m}$ – $2\ \mu\text{m}$) on an axial compressor rotor carried out by means of a commercial computational fluid dynamic (CFD) code. Particle trajectory simulations use a stochastic Lagrangian tracking method that solves the equations of motion separately from the continuous phase. The NASA Rotor 37 is considered as a case study for the numerical investigation. The compressor rotor numerical model and the discrete phase model were previously validated by the authors in the first part of this work. The kinematic characteristics (velocity and angle) of the impact of micrometric and submicrometric particles with the blade surface of an axial transonic compressor are shown. The blade zones affected by particle impact were extensively analyzed and reported in the first part of this work, forming the starting point for the analyses shown in this paper. The kinematic analysis showed a high tendency of particle adhesion on the suction side (SS), especially for the particles with a diameter equal to $0.25\ \mu\text{m}$. Fluid dynamic phenomena and airfoil shape play a key role regarding particle impact velocity and angle. This work has the goal of combining, for the first time, the kinematic characteristics of particle impact on the blade with fouling phenomenon by the use of a quantity called sticking probability (SP) adopted from literature. From these analyses, some guidelines for a proper management of the power plant (in terms of filtration and washing strategies) are highlighted. [DOI: 10.1115/1.4028296]

9 Introduction

10 Ambient air is a continuous medium that contains and carries a
11 large number of particles (contaminants). The contaminants in the
12 air are different in composition, size (pollen of $50\ \mu\text{m}$, spores of
13 $3\ \mu\text{m}$ – $10\ \mu\text{m}$, and exhaust particle $< 0.1\ \mu\text{m}$), and quantity [1].

14 The quality and purity of the air entering the turbine is a signifi-
15 cant factor in the performance and life of the gas turbine. There-
16 fore, air inlet filtration systems are employed to remove a
17 significant amount of the contaminants. Fouling of axial compres-
18 sors (caused by particles smaller than $2\ \mu\text{m}$) is a serious operating
19 problem and its control is of critical importance for operators of
20 gas turbine driven power plants, compressor stations, and pump
21 stations.

22 Estimates have cited fouling as being responsible for 70% to
23 85% of all gas performance losses accumulated during operation.
24 Output losses between 2% (under favorable conditions) and
25 15–20% (under adverse conditions) have been experienced [2]. In
26 order to minimize the performance loss of the turbines in the
27 power plant, an adequate filtration system that can limit the inges-
28 tion of contaminants by the power unit is required.

29 Although 99% of the particles in the atmosphere are less than
30 $1\ \mu\text{m}$ in size, 70% of the weight is due to particles which have a
31 diameter greater than $1\ \mu\text{m}$ [3]. In order to capture these different
32 types of particles, filtration systems use many different mecha-
33 nisms. Each filter in fact has various different mechanisms work-
34 ing together to remove the particles. An extensive report on
35 filtration efficiency can be found in Ref. [4] where it can be seen
36 that for the particles with dimensions less than $\approx 2\ \mu\text{m}$, and in
37 more detail, with diameters in the range of $0.1\ \mu\text{m}$ – $1.0\ \mu\text{m}$, con-
38 ventional filtration systems will not entirely prevent these small
39 particles from entering the gas turbine and therefore may cause
40 fouling.

41 The details on how the small particles entering the gas turbine
42 reach the blade surface and stick there are not fully and quantita-
43 tively understood. Particle adhesion on the blade surface is a com-
44 plex phenomenon that includes many aspects that can be
45 summarized as follows:

- 46 – the material of the body in contact (blade and particle): den-
47 sity, ultimate strength, and elastic yield limit (in order to
48 define an elastic or plastic collision);
- 49 – the surface conditions: roughness, presence of added materi-
50 als (water, oil, and grease), presence of electrostatic charges
51 or its generation by contact (bounces or slips) with the
52 particles;
- 53 – the particle size: inertia force and some energy whose effects
54 are directly related to particle size [5];

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- 55 – the impact velocity is directly related to kinetic energy. This
- 56 energy is transferred into the contact zone between two
- 57 bodies and determines deformation, bounce, and sometimes
- 58 particle breakage;
- 59 – the impact angle represents the most important quantity for
- 60 impact kinematics and consequent impact effects (bounce,
- 61 adhesion, or slip).

62 All these phenomena can be studied from many points of view:

63 (i) kinematic (velocity and direction), (ii) dynamic (velocity and

64 mass), and (iii) energy (deformations or breaks). At the same

65 time, it is possible to consider microscopic effects (e.g., atoms

66 attractions and molecular bonds) or macroscopic (e.g., adhesion,

67 rebounds, and deformations).

68 In almost all cases, the phenomena mentioned above occur at

69 the same time during the contact between two bodies and the

70 result of the contact depends on the combination of these effects.

71 Particle sticking on blade surfaces results in an increase of the

72 thickness of the airfoil and the surface roughness. Both of these

73 events change the flow path inside the passage vanes. This leads

74 to in particular: (i) an increment of boundary layer thickness, (ii) a

75 decrement of the flow passage area, and (iii) modifications of 3D

76 fluid dynamic phenomena.

77 In this paper, the kinematic characteristics of the submicrometric

78 and micrometric particles ($0.25\ \mu\text{m}$ – $2.00\ \mu\text{m}$) that impact on

79 an axial compressor blade will be shown and explained. The particle

80 collision results are reported in Ref. [6] and represent the starting

81 point of this work.

82 Literature Review

83 The interaction between two bodies, with or without the action

84 of an external force, has been a subject of study since the 19th

85 century. The first comprehensive study on the mechanical impact

86 between two bodies was conducted by Hertz in 1882 [7]. The

87 classic approach to this type of problem consists in most cases

88 of description of the impact (usually normal impacts) phenomena

89 which involves two bodies (usually sphere–sphere or sphere–

90 surface) made from a ductile material that has a defined yield

91 load. In this case, the deformation of the bodies plays a key role

92 during contact and determines the result of the impact.

93 One of the major contributions to this field has been provided

94 by Johnson, Kendall, and Roberts (JKR) [8]. The JKR model

95 describes the phenomena that occur between two bodies in contact

96 demonstrating that even if there is not an external force maintain-

97 ing two bodies in contact, there is a well-defined contact area at

98 the body interface and it requires a force greater than zero to separate

99 it. Based on the JKR method, many authors have described

100 contact models based on experimental or analytic results in order

101 to make predictions and estimates regarding the impact behaviors

102 of two bodies.

103 Thornton and Ning [9], for example, have formulated a model of

104 rebound/adhesion which takes into account impact velocity, capture

105 velocity, and yield velocity. Capture velocity represents the

106 particle velocity limit below which contact becomes adhesion and

107 above which contact becomes rebound. Yield velocity represents

108 the material particle velocity limit below which impact can be

109 considered elastic and above which a plastic deformation occurs

110 in the contact zone. Yield velocity is a function of the materials

111 (Young's module, Poisson's ratio, and density). The authors also

112 performed a study on the influence of the energy at the interface

113 and demonstrated that for the highest velocity impact the energy

114 interface does not affect the values of restitution coefficients.

115 One of the most important experimental reports was provided

116 by Wall et al. [10]. The authors have performed a number of

117 experiments with ammonium fluorescein microspheres (with

118 diameters equal to $2.58\ \mu\text{m}$, $3.44\ \mu\text{m}$, $4.90\ \mu\text{m}$, and $6.89\ \mu\text{m}$)

119 impacted normally against smooth, flat surfaces of polished

120 molybdenum, silicon, cleaved mica, and a fluorocarbon polymer

121 over an initial velocity up to 100 m/s. The main results shown in

this work can be summarized as follows: (i) at low velocity 122

(<20 m/s), the ratio of rebound to impact velocity was sensitive to 123

target material, decreasing with impact velocity due to the adhesion 124

surface energy, (ii) the kinetic energy recovered in low velocity 125

impacts was found to depend on particle size, (iii) no such particle 126

size dependence was observed for impact velocities near 20 m/s, 127

and (iv) above 40 m/s, the velocity ratio was insensitive to the target 128

material, indicating that the particle has a lower elastic yield limit 129

than the material target. Finally, the authors have highlighted that 130

the plastic deformation was a significant component of energy loss 131

at all impact velocities and the knowledge of interface energy plays 132

a key role for the proper description of particle impact. 133

134 Unfortunately, most of the models and the results reported in

135 literature do not provide a full understanding of the adhesion phe- 136

137 nomena which is responsible for the fouling mechanism. This 138

139 limit is largely due to (i) different particle sizes, (ii) different 140

141 material characteristics (some particle materials do not show the 142

143 elastic yield limit), and (iii) the different impact velocity. In fact, 144

145 if the model reported in Ref. [9] is applied to a metallic micro- 146

147 sized particle with Young's module equal to 72 GPa, Poisson's ratio 148

149 equal to 0.17, and a surface energy equal to $0.2\ \text{J/m}^2$; the capture 150

151 velocity is approximately 1 m/s. By using these results for 152

153 fouling phenomena and considering that the submicrometric and 154

155 micrometric particles follow the streamline with a velocity of 156

157 about 350 m/s, all of these particles must bounce on the blade sur- 158

159 face and the fouling phenomena would not exist. 160

161 Some very interesting results and detailed analysis of micropar- 162

163 ticle adhesion can be found in astrophysics applications related to 164

165 the research of preplanetary dust dynamics. The particles, in most 166

167 cases consisting of submicrometric silica spheres, are the basis of 168

169 the planets' origin. These space-dispersed particles collide with 170

171 each other, and if the impact allows for adhesion, the particles 172

173 generate an agglomerate. The difference between these mecha- 174

175 nisms, not yet fully understood and the fouling phenomenon, is 176

177 due to the type of motion. In fact, the preplanetary particles move 178

179 in cosmic space, characterized by high Knudsen numbers (molec- 180

181 ular motion), while in the case of fouling, and more generally of 182

183 motions in the Earth's atmosphere, the transportation of the par- 184

185 ticles takes place with very small Knudsen numbers (viscous 186

187 motion). This aspect must not diminish the importance of the 188

189 results highlighted in this research field because the experiments 190

191 are conducted only in some cases under vacuum and the results 192

193 are often in line with the more classical theories mentioned above. 194

195 The uniqueness and usefulness of these studies are that the parti- 196

197 cle velocities, materials, and dimensions are in the same range as 198

199 those responsible for the fouling phenomenon. 200

201 In this field of research, one of studies closest to the fouling 202

203 phenomena is Ref. [11], reporting experimental evaluations of 204

205 perfectly spherical and irregular particles impacting a smooth sur- 206

207 face (smooth as the particle surface). Different combinations of 208

209 particle size and materials have been tested. The particle diame- 210

211 ters are very close to $1\ \mu\text{m}$ and in some cases the experiments 212

213 were conducted with submicrometric particles. The material 214

215 (silica in some cases) has a density of about $2000\ \text{kg/m}^3$ (silica). 216

217 The main results reported in this work can be summarized as 218

219 follows: 220

- for impact velocities in the range 1 m/s–10 m/s, the kinetic 178
- energy is typically reduced to one-half for the $1.2\ \mu\text{m}$ and to 179
- one-quarter for the $0.5\ \mu\text{m}$;
- for impact velocities exceeding 10 m/s, the bouncing colli- 180
- sions reduce the kinetic energy by more than one order of 181
- magnitude; 182
- for the $1.2\ \mu\text{m}$ diameter silica spheres, the capture velocity is 183
- independent of the target surface tilt angle (0 deg–60 deg); 184
- electrostatic effects occur during the test and the action of the 185
- electrostatic field is observed up to $40\ \mu\text{m}$ from the surface. 186
- The same effect can be found in Ref. [12]; 187
- experimental results obtained with irregular shaped particles show 188
- a higher capture velocity and, at the same time, a higher SP. 189

190 The SP is the one of the most interesting quantities used in this
191 type of study. The SP was evaluated by a statistical approach. It
192 compares the impact that results in sticking with the total amount
193 of impacts. For the SP, the most common threshold limit is equal
194 to 0.5. If more than 50% of the sampled impact results in sticking
195 between the particle and surface, the final results will be sticking
196 or, vice versa, if less than 50% of the sampled impact results in
197 sticking, the final result will be bounce. With this approach, the
198 authors wish to emphasize that the particle impacts are different
199 from each other and, in order to provide a macroscopic evaluation
200 of the results, a statistic/probabilistic approach is the best way.

201 From the reported literature, it is easy to understand that for the
202 total comprehension of the fouling phenomena it must be known
203 how the contaminants hit the blade surface. In this context, the
204 word *how* refers to the impact velocity and the impact angle for
205 each particle. In this paper, for the first time, we will show the
206 kinematic characteristics of the particle impact on the axial com-
207 pressor blades by means of a CFD numerical simulation.

208 In literature, some interesting experimental and numerical studies
209 on particle ingestion can be found. In particular, research
210 regarding particle erosion and deposition in gas turbines can be
211 found in Refs. [13] and [14], while for the axial compressor some
212 interesting results can be found in Refs. [15] and [16]. We will
213 focus only on the axial compressor studies. In Ref. [15], the
214 authors performed a study of the erosion effects in an axial com-
215 pressor stage. The particles have a diameter equal to 165 μm and
216 the results show the particle trajectories also after the first impact.
217 In the case of the second impact, the erosion magnitude is very
218 low due to the low particle kinetic energy. Ghenaïet [16] studied
219 the particle dynamics and erosion of the front compression stage
220 of a turbofan PW-JT8-D17. Particle trajectory simulations used a
221 stochastic Lagrangian tracking code and the sand particle size
222 varies from 0 μm to 1000 μm . The numerical simulations show a
223 different trajectory for the different particle diameter. The larger
224 particles are affected by inertia and centrifugal forces and after
225 the first impact, they do not follow the airflow stream. The smaller
226 particles, $\approx 10 \mu\text{m}$, tend to follow the flow path closely and are
227 strongly influenced by the flow turbulence, secondary flows, and
228 flow leakage above the blade tip and in due course, induce erosion
229 of the blade tip and shroud. Particles with a diameter less than
230 10 μm were not taken into account for the erosion analyses, since
231 particles of this size do not carry enough energy to cause erosion.

232 In the present work, the authors presented a CFD study for the
233 ultrafine powder ingestion (particle size of 0.25 μm –2.00 μm) by an
234 axial compressor rotor, the NASA Rotor 37. These particle sizes can
235 cause fouling, but are too small to cause erosion. The particle inges-
236 tion was studied by using a CFD commercial code. In particular, in
237 this second part, the authors, beginning with the results reported in
238 Ref. [6], will show the kinematic characteristics of particle impact on
239 the blade surface. This paper includes the following points:

- 240 – a short reference of the numerical model adopted for the con-
241 tinuous and discrete phase (validation and more details can be
242 found in Ref. [6]);
- 243 – a short summary of the particle impact zones on the blade sur-
244 face in order to better understand the following analyses
245 (extensive analyses can be found in Ref. [6]);
- 246 – analyses of particle impact velocity and particle impact angle
247 for the pressure side (PS) and SS;
- 248 – an analysis of the normal and tangential velocity component
249 in order to define the relative impact kinematic characteristics
250 between blade and particles;
- 251 – estimates of the SP up to 1 μm particle diameter in order to
252 define the preferable deposition zones on the blade as a func-
253 tion of the particle diameter.

254 Numerical Model

255 **Continuum Phase.** The reference compressor stage is the
256 NASA Rotor 37 [17]. It is composed of 36 blades but only a single
257 passage vane was modeled. The tip clearance at design speed is

0.356 mm (0.45% of the blade span). All the simulations were per-
formed in a steady multiple frame of reference by using a frozen
stage interface. A multiblock hexahedral grid with a total number
of 1,131,063 elements was used. The numerical simulations were
carried out by means of the commercial CFD code ANSYS FLUENT
13.0 [18]. The standard k - ϵ turbulence model with a standard wall
function was used. The numerical CFD results are in fairly good
agreement with the experimental data. The numerical pressure ratio
 β and the total-to-total efficiency η_{TT} always underestimate the
experimental data but in a very consistent way. The deviation in
terms of mass flow rate at the choked-flow condition is about 1.87%
(all the simulations refer to design speed equal to 17,188 rpm). More
details for the numerical domain and validation are given in Ref. [6].

Discrete Phase. In this paper, the solution approach is based
on a mathematical model with Eulerian conservation equations in
the continuous phase and a Lagrangian frame to simulate a dis-
crete second phase. In this approach, the airflow field is first simu-
lated, and then the trajectories of individual particles are tracked
by integrating a force balance equation on the particle. The force
balance is comprehensive of inertia, drag, and buoyancy term. In
the force balance, there are two contributes due to the shear stress
and diffusion called Saffman's lift force and Brownian force but
these two contributes become important in very few cases. In this
paper, only the Brownian term was neglected. An extensive
description of the force balance can be found in Ref. [6].

The dispersion of particles in the fluid phase can be predicted
by using a stochastic tracking model. This investigation used the
discrete random walk (DRW) model to simulate the stochastic ve-
locity fluctuations in the airflow. The number of trajectories was
selected in order to satisfy the statistical independence, since the
turbulent dispersion is modeled based on a stochastic process. The
inlet/injection surface was made by 1888 uniformly distributed
elements and each analysis of three different injections with 1500
trajectories was carried out.

For the particle-wall interaction boundary conditions, the fol-
lowing conditions have been adopted: (i) ideal adherence condi-
tion (named *trap*) on the blade surfaces and (ii) nonadherence
condition (named *reflect*) on the hub and shroud surfaces. These
conditions allow the evaluation of where and how the contami-
nants encountered the blade surface for the first time, avoiding the
introduction of inaccuracies due to the use of bounce models not
fully representative of the real conditions. The authors have
implemented specific functions and a restitution coefficient for the
near-wall particle behavior. The model functions are defined in
agreement with Ahlert's [19] model and Forder's [20] coeffi-
cients. More details regarding particle-wall interaction can be
found in Ref. [6].

The density particle is equal to 2560 kg/m^3 and the variation of
the particle diameter, d_p , is in the range of 0.25 μm –2.00 μm ,
while the Stokes number (calculated at the inlet of the numerical
model) is in the range of 0.0010–0.0630. All particles are spheri-
cal and nondeformable.

All the analyses refer to injections having particles with the
same diameter, the same material, and thus characterized by
the same Stokes number. On the other hand, the total flow rate of
the discrete phase m_p is linked to the work environment of the com-
pressor and the efficiency of the filtration system. For this reason, a
different value of total flow rate of contaminants was imposed at
the inlet of the compressor. All injections take place on a previously
solved flow field, with the compressor operating at the best effi-
ciency point. All results presented in this paper were obtained from
convergent simulations, with a variation of the residues of the
motion and turbulent equations close to zero. The injection data are
summarized in Table 1 (more details can be found in Ref. [6]).

258 Previous Results

259 In this paragraph, the previous results referring to the particle
260 impact on the NASA Rotor 37 blade are summarized. The entire

Table 1 Characteristics of the injections

Case	1	2	3	4	5
Particle diameter, d_p (μm)	0.25	0.50	1.00	1.50	2.00
Stokes number, St	0.0010	0.0039	0.0158	0.0355	0.0630
Total flow rate, m_p (kg/s)	3.51×10^{-6}	2.46×10^{-5}	8.43×10^{-5}	7.59×10^{-5}	4.50×10^{-5}

Table 2 Particle impact distributions

Case	1	2	3	4	5
η_{hit}	2.38	3.18	4.48	7.99	13.76
$\eta_{hit,PS}$	1.66	2.75	4.31	7.87	13.72
$\eta_{hit,SS}$	0.72	0.43	0.17	0.11	0.04
$\eta_{SIDE,PS}$	69.70	86.40	96.20	98.60	99.70
$\eta_{SIDE,SS}$	30.30	13.60	3.80	1.40	0.30

Table 4 Mass contaminant on the blade (kg)

d_p (μm)	Operating time t			
	1 h	1 day	2 days	1 week
0.25	0.001	0.022	0.043	0.151
0.50	0.008	0.203	0.404	1.421
1.00	0.453	10.861	21.721	76.025
1.50	0.086	2.054	4.107	14.375
2.00	0.067	1.601	3.203	11.210

$$M_c = \overline{AR} A_b t \tag{1}$$

analysis is reported in Ref. [6]. Only a portion of particles injected from the inlet surface of the numerical model has an impact with the blade surface, and due to the imposed surface condition (trap), the contact results in a permanent adherence. For the comparison between the studied cases, two types of percentage were used. The first one is defined as the ratio between the number of particles that hit (and then could stick to) the blade and the total number of injected particles. The second one is defined as the ratio η_{side} between the impacting particles on the PS $\eta_{SIDE,PS}$ or SS $\eta_{SIDE,SS}$ compared to the total number of impacting particles on the blade. In order to provide a useful value of the impacting particles, η_{hit} are also reported for the PS $\eta_{hit,PS}$ and SS $\eta_{hit,SS}$ and refers to the percentage of impacting particles on the PS or SS compared to the total number of injected particles, respectively. All values in Table 2 are reported but an extensive analysis can be found in Ref. [6].

Analyses of the Particle–Blade Interaction

Accretion Rate (AR). The first analysis of the particle–blade interaction refers to the quantity defined as AR which allows the identification of contaminant deposition intensity in terms of $\text{kg}/\text{m}^2 \text{ s}$. The AR, defined on a blade surface, allows the evaluation of the combined effects between the trajectories of the particles and the contaminant total mass flow rate. The AR values are obtained for each run of the respective cases. As mentioned above, each case was repeated for three different runs in order to avoid the problem caused by the statistical resolution of particle tracking. In Table 3, the values of the peak AR^* and the values obtained by a weight-area average AR for all of the executed runs are reported. From the values of Table 3, it is possible to note that the values obtained for the three runs of each case have the same order of magnitude, confirming the independence of the results from the statistical dispersion. With this evidence, it is possible to define an average value \overline{AR} of the AR for the three runs in each case. With these values, the amount of contaminants that affected the blade surface during the operation can be evaluated. In fact, it is possible to calculate the contaminant mass M_c on the blade surface as

where A_b is the blade surface and t is the operating time. Table 4 summarizes the mass of contaminants on the blade surface for different operation times.

The blade contamination is very noticeable with a very high contaminant mass on the blade surface (up to 76 kg) even after one operation week. These values of the mass deposits are due to the numerical model wall condition on the blade surfaces (trap) that imposes an ideal adhesion for each particle that hits the blade. Hence, the absolute values are not representative of the particle deposition because the sticking phenomena change during the particle deposition due to the different characteristics between the blade surface and the deposited particle layer. Rather, they are representative of the total amount of contaminants which hit the blade.

Impact Velocity

The first analysis is related to the particle impact velocity v_i . The modules of the particle impact velocity are reported in Fig. 1. The velocity values refer to the vector sum of the three velocity components u along the coordinate axes x , y , and z at the impact point on the blade surface. In Fig. 1, the most representative strips are reported: second, sixth, and tenth (12%, 47%, and 83% of the span blade, respectively) divided into PS and SS. Each dot on the graph corresponds to the impacting particle on the blade. From Fig. 1, it can be noticed that

- the impact velocity increases with the height of the blade and this phenomenon is due to the peripheral velocity;
- the lowest impact velocity can be found on the leading edge (LE) and on the trailing edge (TE) of the SS;
- the highest impact velocity can be found on SS, in particular, on the first part of the airfoil chord;
- the effects of flow separation (due to the shock wave) can be clearly seen on the SS. This phenomenon causes the drop of

Table 3 AR values ($\text{kg}/\text{m}^2 \text{ s}$)

d_p (μm)	First run		Second run		Third run		Average
	AR^*	\overline{AR}	AR^*	\overline{AR}	AR^*	\overline{AR}	\overline{AR}
0.25	4.1×10^{-3}	1.1×10^{-5}	1.3×10^{-2}	3.3×10^{-5}	2.1×10^{-2}	5.6×10^{-5}	3.3×10^{-5}
0.50	4.7×10^{-2}	1.1×10^{-4}	1.4×10^{-1}	3.1×10^{-4}	2.4×10^{-1}	5.2×10^{-4}	3.1×10^{-4}
1.00	3.0×10^0	1.0×10^{-2}	4.5×10^0	1.5×10^{-2}	7.6×10^0	2.5×10^{-2}	1.7×10^{-2}
1.50	5.6×10^{-1}	3.1×10^{-3}	4.4×10^{-1}	2.4×10^{-3}	7.4×10^{-1}	4.0×10^{-3}	3.2×10^{-3}
2.00	1.8×10^{-1}	8.3×10^{-4}	5.2×10^{-1}	2.5×10^{-3}	8.5×10^{-1}	4.1×10^{-3}	2.5×10^{-3}

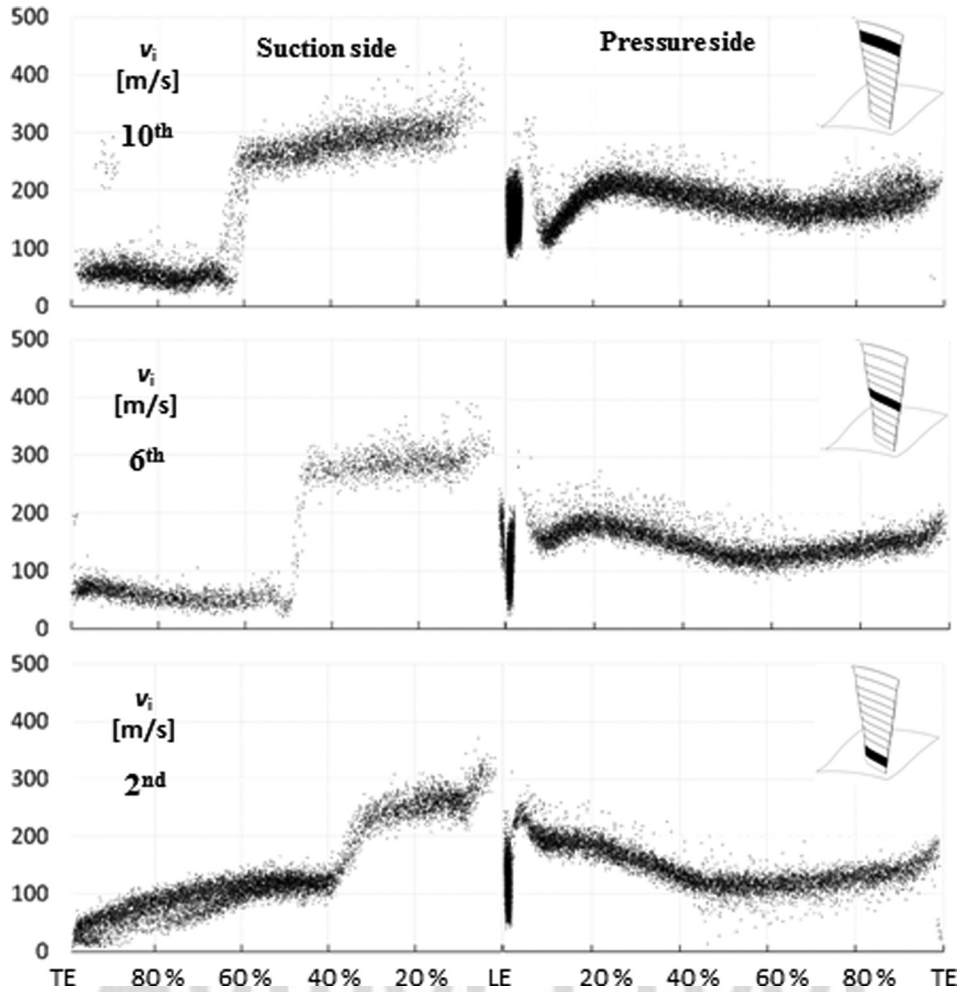


Fig. 1 Impact velocity v_i , second, sixth, tenth strip, case 1

394 the particle impact velocity and at about 50% of the chord we
 395 find the lowest impact velocity;
 396 – on the PS, the velocity trend is very similar for all the strips.
 397 On the first part of the chord, the particles reach the peak of
 398 impact velocity, while at about the 60% of the airfoil chord
 399 the impact velocity reaches a minimum.

400 The analysis of Fig. 1 shows that the particle impact velocity is
 401 very different on the same side of blade. This difference is due to
 402 the shape of the blade (e.g., the blade height) and the fluid
 403 dynamic phenomena (e.g., flow separation). Another fluid
 404 dynamic phenomenon that influenced the particle impact velocity
 405 at the top of the blade is the tip leakage vortex due to the blade tip
 406 gap. The effect on the particle impact location of this particular
 407 phenomenon is clearly investigated in Ref. [6] and in this second
 408 part, its effect on particle velocity impact can be seen. As is shown
 409 for the SS in Fig. 1, the rear part of the airfoil chord is impacted
 410 by particles with a very low impact velocity while for the 11th
 411 strip, reported in Fig. 2, this is not quite the case. The rear part of
 412 the airfoil chord of the 11th strip is impacted at the same time by
 413 particles with very low and very high impact velocity. The particles
 414 with the highest impact velocity are the particles dragged by
 415 the tip leakage vortex from the PS to the SS.

416 In this specific case, the wall condition imposed on the blade
 417 (trap) determines a smaller amount of particles that are dragged
 418 from the PS to the SS. Under real conditions, some particles
 419 bounce off the PS and could reach the other side of the blade
 420 through the tip gap. This effect plays a key role in the erosion
 421 problem not considered in this work due to the small particle
 422 sizes. In fact, the erosion phenomena require a particle diameter

larger than $10 \mu\text{m}$ as reported by Hamed et al. [13], Ghenaiet [16],
 and Kurz and Brun [21].

Impact Angle. As can be seen from the previous analyses, the
 particle impact velocity changes from the hub to the shroud, from
 the PS to the SS and along the airfoil chord. However, the impact
 velocity v_i is not the only parameter needed to determine particle
 adhesion on the blade surface. As mentioned above, particle adhesion
 is due to a combination of a number of effects, but the most
 important parameters are the normal v_n and tangential v_t velocity

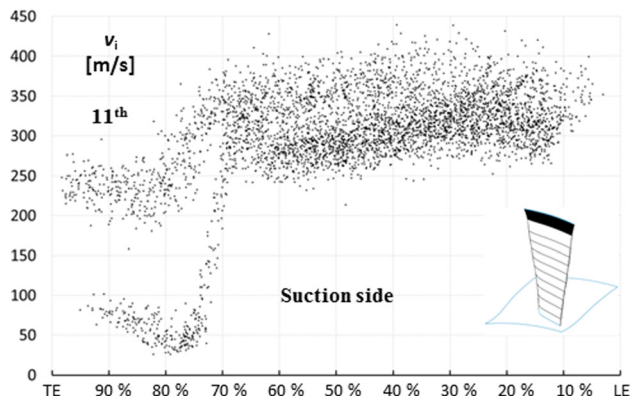


Fig. 2 Impact velocity v_i , 11th strip SS, case 1

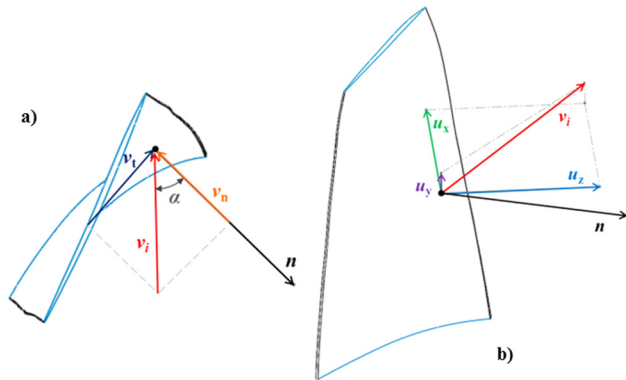


Fig. 3 Particle vectors velocity: (a) normal and tangential impact velocity and (b) impact velocity

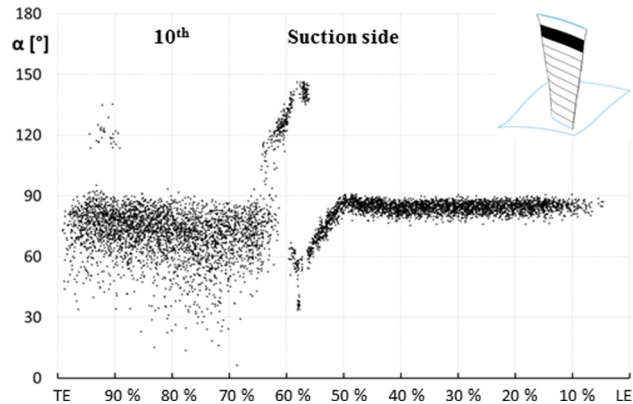


Fig. 5 Impact angle α , tenth strip, SS, case 1

432 components. In this paragraph, we analyze the particle impact
433 angle in order to better understand the particle kinematic impact.

434 The impact velocity was obtained by a vector sum of the three
435 velocity components u_x , u_y , and u_z along the axes x , y , and z ,
436 respectively (Fig. 3(b)). The impact velocity was decomposed
437 with respect to the normal (v_n) and tangential (v_t) direction
438 (Fig. 3(a)). Thus, the impact angle α is the angle between the surface
439 normal vector n and the impact velocity vector v_i .

440 In Figs. 4 and 5, the particle impact angle for the PS of the sixth
441 strip and for the SS of the tenth strip (case 1) are reported. We
442 find in Figs. 4 and 5 that in some instances the impact angle is
443 higher than 90 deg. This is due to (i) the surface local curvature
444 (e.g., at the LE and on the TE) and (ii) surface reconstruction
445 approximation during the particle impact postprocess. A deviation
446 can arise from the fact that the surface is reconstructed by interpo-
447 lating points on the mesh elements in the neighborhood of the
448 point of impact. The approximation introduced by this procedure
449 is considered acceptable by the authors, allowing for a confidence
450 band of ± 5 deg for all the results shown in this paper.

451 Figures 4 and 5 illustrate the following observations:

- 452 – the impact angle at the LE (Fig. 4(a)) assumes different val-
453 ues from 30 deg to 120 deg;
- 454 – on the PS (Fig. 4(a)), the particle impact angle is very close
455 to 90 deg (i.e., the particles are tangential to the blade surface)
456 almost everywhere on the airfoil. A particular area can be
457 noticed in the middle of the chord where the particle impact
458 angle reaches 120 deg. This fact is consistent with Fig. 5(b)
459 where the representation of the PS curvature is reported. The
460 blue zone refers to a lower curvature, while the red zone
461 refers to higher curvature. The local variation of the impact
462 angle (gray box) corresponds to the local variation of the sur-
463 face curvature (gray circle). Thus, it is clearly shown that the

464 local curvature of the airfoil (e.g., dimples, surface damage,
465 etc.) changes the particle impact angle in a significant way
466 and, more generally, the local shape of the blade changes the
467 particle deposition. A different impact angle can determine
468 whether the particle sticks or slips and thus, the actual shape
469 of the blade surface would determine the magnitude and the
470 rate of the fouling. These findings represent a useful guide for
471 blade surface treatment and control during the manufacturing
472 and maintenance process. The same phenomenon can be
473 noticed for all the strips;

- 474 – for the SS, there is also a variation of the particle impact
475 angle in the middle of the chord due to the airfoil curvature.
476 However, it is less noticeable than on the PS;
- 477 – on the SS, the particle impact angle is lower than the PS and
478 this implies that the particle hits the surface with a value of
479 normal velocity higher than the tangential velocity. This is
480 noticeable in the last part of the chord where the flow is sepa-
481 rated from the blade.

482 Areas characterized at the same time by very high tangential
483 velocity and very low normal velocity (impact angle close to
484 90 deg) should not be subject to particle deposition because in this
485 case the particles tend to slip on the blade surface. However, in
486 the other areas with a lower impact angle, the normal velocity pro-
487 motes particle sticking (e.g., in the case of ductile particles). Simi-
488 lar evaluations can be made for cases in which the blade surface is
489 contaminated by water, oil, or grease and in the case of viscous
490 particles (e.g., oils and grease) that should stick to the blade sur-
491 face more easily because of the high normal velocity. A study on
492 particle sticking by liquid materials can be found in Ref. [22] in
493 which the authors show the coefficient of restitution trend for mol-
494 ten granulate. For this material, the results show a lower bounce
495 capability at high impact velocities.

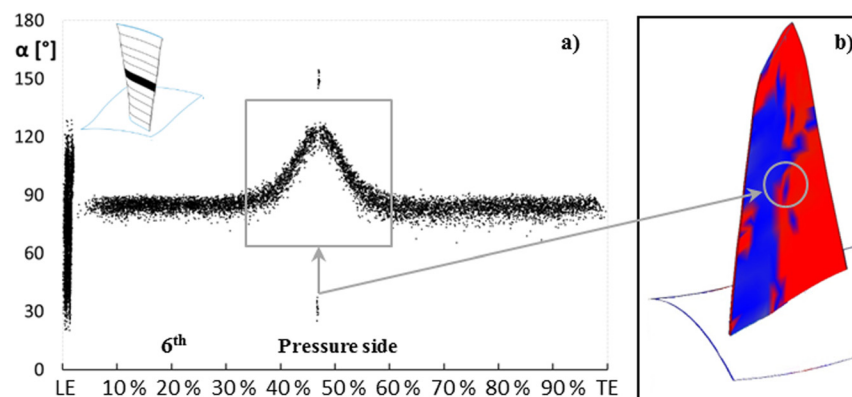


Fig. 4 (a) Impact angle α , sixth strip, PS, case 1 and (b) contour plot of the PS curvature

496 However, these preliminary considerations must also be related
 497 to the material properties (particles and surface) and to the surface
 498 roughness (high or low). The particle slip is related to the surface
 499 roughness. Greater roughness causes more slip resistance
 500 and therefore easier adhesion or easier breaking away of
 501 particles [23].

502 As shown in the first paragraphs, the study of particle adhesion
 503 on a surface comprises a large number of aspects and probabilistic
 504 analyses are often used due to the unique nature of each contact.
 505 In this paper, the authors provide a quantitative analysis of particle
 506 adhesion by using the experimental results found in Ref. [11] in
 507 which particle velocity and materials are among the most similar
 508 to the particles causing fouling phenomena.

509 **SP.** The SP analysis is closely related to the experimental
 510 results provided by Poppe et al. [11]. The SP defined in Ref. [11]
 511 was calculated for each normal impact velocity v_n by sliding aver-
 512 aging in groups. The groups consisted of 11 collision events for
 513 the smallest and largest velocities, respectively, and of up to 71
 514 collision events for the intermediate velocities, thus accounting
 515 for the uneven velocity distribution of the impacts. The upper and
 516 lower standard deviations (1σ) for the SP are reported in Figs. 6
 517 and 7 (black continuous lines). The trends refer to irregular grains
 518 of silicon carbide ($\rho_p \approx 3000 \text{ kg/m}^3$, $E \approx 410 \text{ GPa}$, and hardness
 519 $\approx 2800 \text{ kg/mm}^2$) with an average diameter equal to $0.37 \mu\text{m}$ and
 520 $0.64 \mu\text{m}$ that impacts a dry, polished silica surface. In Figs. 6 and

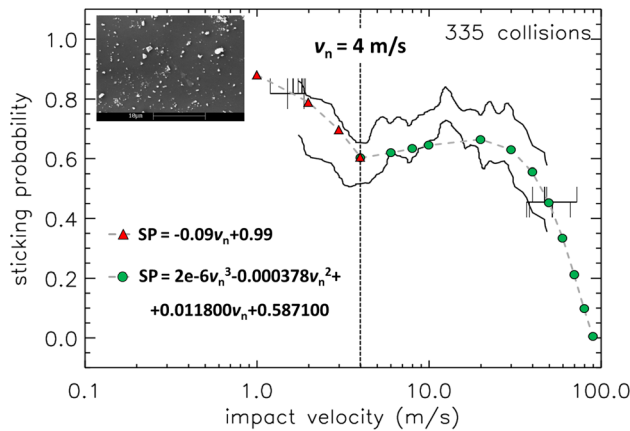


Fig. 6 SP versus normal impact velocity v_n of silicon carbide particles, $0.37 \mu\text{m}$ on silica target [11], and trend of adopted equations superimposed

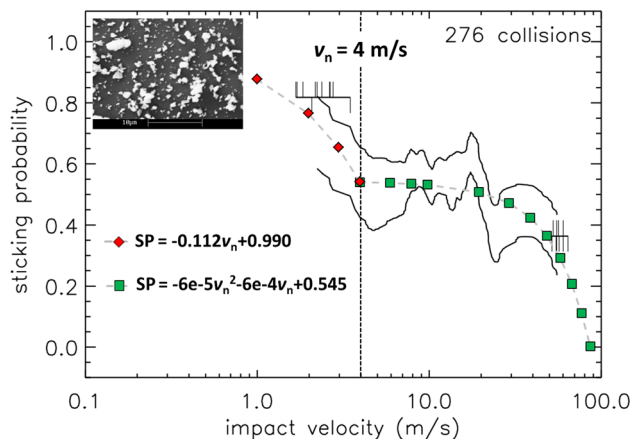


Fig. 7 SP versus normal impact velocity v_n of silicon carbide particles, $0.64 \mu\text{m}$ on silica target [11], and trend of adopted equations superimposed

7, SEM images of a silicon carbide sample are also reported. The
 SP among the 11 slowest and the 11 fastest collisions is separately
 given as a constant value in the corresponding velocity interval
 (reported with a linear segment in Figs. 6 and 7). The capture
 velocity is the velocity where the 1σ limits of the SP are 0.5. Such
 a definition results in a physically meaningful quantity only if the
 SP behaves similarly to a step function. More details on materials
 and experimental results can be found in Ref. [11].

529 With the experimental SP trend reported in Figs. 6 and 7, it is
 530 possible to define representative trends for the correlation between
 531 the normal impact velocity v_n and the SP. For the smaller silicon
 532 carbide particles ($0.37 \mu\text{m}$) reported in Fig. 6, the trend can be re-
 533 presented by two equations. The first one refers to the lower normal
 534 impact velocity ($<4 \text{ m/s}$)

$$S_p = -0.09 v_n + 0.99 \quad (2)$$

535 and the second one refers to normal impact velocity in the range
 536 of $4 \text{ m/s} - 90 \text{ m/s}$

$$S_p = 2 \times 10^{-6} v_n^3 - 0.000378 v_n^2 + 0.011800 v_n + 0.587100 \quad (3)$$

537 The results of Eqs. (2) and (3) are reported in Fig. 6, with the
 538 experimental results obtained by Poppe et al. [11] superimposed.
 539 As can be noticed from Eq. (2), in the case of the normal impact
 540 velocity equal to 0 m/s , the SP is equal to 0.99 .

541 In the same way, for the larger silicon carbide ($0.64 \mu\text{m}$)
 542 reported in Fig. 7, the trend can be represented by two equations.
 543 The first one refers to the lower normal impact velocity ($<4 \text{ m/s}$)

$$S_p = -0.112 v_n + 0.990 \quad (4)$$

544 and the second one refers to normal impact velocity in the range
 545 of $4 \text{ m/s} - 90 \text{ m/s}$

$$S_p = -6 \times 10^{-5} v_n^2 - 6e-4 v_n + 0.545 \quad (5)$$

546 Again, the results of Eqs. (4) and (5) are reported in Fig. 7, with
 547 the experimental results obtained by Poppe et al. [11] superim-
 548 posed. The threshold normal velocity (equal to 4 m/s) and the
 549 degree of the polynomials were chosen in order to better describe
 550 the experimental trend results.

551 With the definition of the SP (Eqs. (3) and (5)), for cases 1 and
 552 2, the $SP = 0.5$ is in correspondence to a normal impact velocity
 553 v_n equal to 48.35 m/s . However, for case 3, the $SP = 0.5$ is in cor-
 554 respondence to a normal impact velocity v_n equal to 22.85 m/s .
 555 Thus, the smaller particles have a wider range of normal impact
 556 velocity for which particle impact with the blade surface becomes
 557 (with a high probability) a permanent adhesion.

558 Equations (2)–(5) are used to calculate the SP for each particle
 559 stuck to the blade surface by using the normal impact velocity.
 560 The particle characteristics used in Ref. [11] are quite different
 561 compared to the classic particle characteristics involved in fouling
 562 phenomena. In particular, the silicon carbide particles [11] have a
 563 very high level of hardness and this implies that the rebound prop-
 564 erties could be different from those found in the real fouling
 565 applications.

566 In Fig. 8, the SP for the sixth strip (case 1) is reported. Each dot
 567 on the graph represents a particle that hit the blade surface with a
 568 normal impact velocity less than 90 m/s . Only the particles with a
 569 normal velocity component toward the surface are taken into
 570 account. This procedure allows the identification of the dangerous
 571 particle (that will be able to stick) with respect to fouling phenom-
 572 enon only. Figure 8 illustrates that

- the SS is completely covered by particles that have a SP of about 0.7 ;
- the PS shows an area, in the middle of the airfoil chord, in which the particles have a SP equal to zero. This effect is due

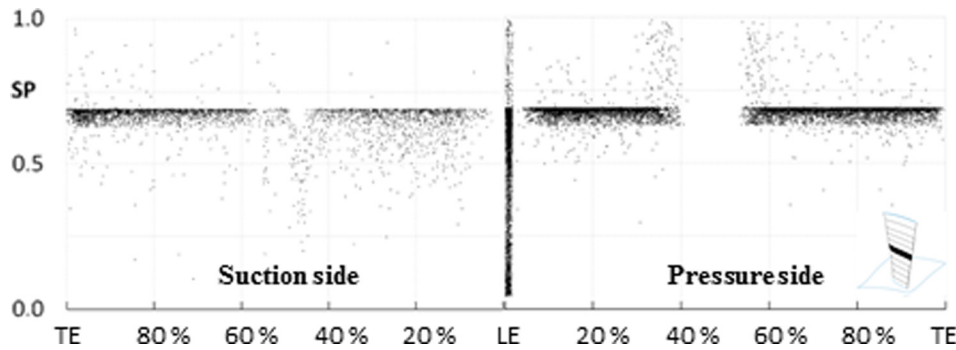


Fig. 8 SP, sixth strip, case 1

577 to the airfoil shape as highlighted in Fig. 4. For the other
 578 regions in the PS, the SP is comparable with the SP on the
 579 SS, and in some cases reaches the unit;
 580 – on the LE, there are very dispersed values of SP, probably
 581 due to the wide range of the impact angle as can be seen in
 582 Fig. 4.

583 The other strips show similar features. As mentioned above, the
 584 SP defined in Ref. [11] only considers the normal impact velocity.
 585 However, in this application, particular attention must be paid to
 586 the tangential impact velocity. In fact, as can be seen in Fig. 9 for
 587 the sixth strip, the magnitude of the tangential impact velocity is
 588 not negligible. The tangential impact velocity can reach 250 m/s
 589 or 400 m/s in the PS and SS, respectively. These very high values
 590 may diminish the SP and transform the adhesion-impact in to the
 591 slip-impact. Conversely, it can be noted that in the separation
 592 zone on SS, where the SP is equal to 0.7, the tangential impact
 593 velocity is much smaller, thus limiting the possibility of slip
 594 between the particle and blade surface. Regarding this aspect,
 595 some field data can be found in Ref. [24], where the authors high-
 596 lighted the higher deposition rate where the shear stress between
 597 air and blade surface is lower. This confirms the results obtained
 598 in this work by linking together the SP data and the impact
 599 dynamic characteristics.

600 Unfortunately, specific studies on the interaction between normal
 601 impact velocity, tangential impact velocity, and surface
 602 roughness are not available in literature. In addition, specific stud-
 603 ies on the variation of the SP due to the presence of a third mate-
 604 rial at the interface between surface and particle are not available
 605 in literature. Poppe et al. [11] pointed out that the presence of
 606 hydrophobic silane coating did not change the collisional behavior
 607 with respect to another test in which the surface was only cleaned
 608 with alcohol and subsequently dried with pressurized air.

609 Generally, in the actual compressors, the presence of a third
 610 substance (such as oil and grease) on the blade surface could
 611 increase the SP of the particle, but, at the moment, there are no
 612 specific studies that allow the quantification of this effect.

613 In Table 5, all the impact characteristics are reported for cases
 614 1, 2, and 3 that are considered by the authors the most interesting
 615 cases from a fouling point of view.

616 The particles are subdivided by using normal impact velocity
 617 criteria. In particular, the following three categories are defined:

- the particles that move away from the surface (called harmless);
- the particles that have a normal impact velocity less than 90 m/s and for which it can be possible to define an SP by using Eqs. (2)–(5);
- the particles that have an impact normal velocity higher than 90 m/s and for which the SP is assumed equal to zero.

625 Special attention must be paid to the last category, character-
 626 ized by an impact normal velocity higher than 90 m/s and an SP
 627 equal to zero. These particles possess high kinetic energy that
 628 decreases by an order of magnitude during the first impact as
 629 reported in Ref. [11]. This phenomenon implies that these par-
 630 ticles will not be able to stick during the first contact but instead,
 631 it will most likely be during the second one. In fact, the decrease
 632 in kinetic energy is strongly related to the decrease in velocity
 633 and, consequently, an increase of SP. A similar effect can also be
 634 found in turbomachinery applications. In Ref. [15], the authors
 635 have described the poor erosion capacity shown by the particles
 636 during the second contact with the blade caused by a low level of
 637 kinetic energy that corresponds to that observed in Ref. [11]. If
 638 this phenomenon is not important from an erosion point of view
 639 (due to the particle diameter lower than 10 μm), for the fouling
 640 problems, a low particle kinetic energy during impact with the
 641 blade leads to a high SP.

642 Table 5 shows for all categories listed above: (i) the absolute
 643 number of particles N that have impacted on that side (PS or SS)
 644 and on that band (1st–11th), (ii) the ratio n_{SIDE} between the abso-
 645 lute number N and the total number of particles that impacted on
 646 that side of the blade, and (iii) the ratio n_{hit} between the abso-
 647 lute number N and the total number of injected particles. Thus, the
 648 ratio n defines the kinematic characteristics distribution on one

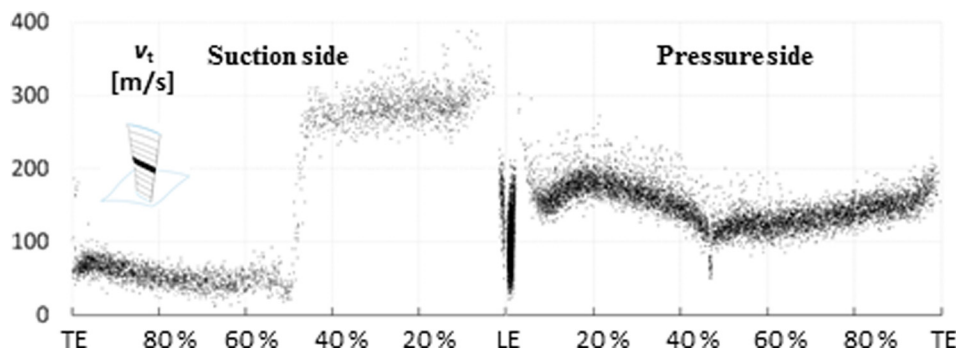


Fig. 9 Tangential velocity v_t , sixth strip, case 1

Table 5 Particle-blade interaction

		Case 1 ($d_p = 0.25 \mu\text{m}$)						Case 2 ($d_p = 0.50 \mu\text{m}$)						Case 3 ($d_p = 1.00 \mu\text{m}$)					
		PS			SS			PS			SS			PS			SS		
		N	n_{SIDE} (%)	n_{hit} (%)	N	n_{SIDE} (%)	n_{hit} (%)	N	n_{SIDE} (%)	n_{hit} (%)	N	n_{SIDE} (%)	n_{hit} (%)	N	n_{SIDE} (%)	n_{hit} (%)	N	n_{SIDE} (%)	n_{hit} (%)
11th	Harmless	8004	5.69	0.09	1333	2.18	0.02	17,091	7.32	0.20	423	1.15	0.00	29,258	8.00	0.35	118	0.81	0.00
	$0 < v_n \leq 90 \text{ m/s}$	4392	3.12	0.05	2678	4.38	0.03	3600	1.54	0.04	1074	2.92	0.01	7259	1.99	0.09	96	0.66	0.00
	$v_n > 90 \text{ m/s}$	283	0.20	0.00	328	0.54	0.00	290	0.12	0.00	102	0.28	0.00	357	0.10	0.00	13	0.09	0.00
	$\text{SP} \geq 0.5$	4045	2.88	0.05	1717	2.81	0.02	3354	1.44	0.04	828	2.25	0.01	4337	1.19	0.05	44	0.30	0.00
Tenth	Harmless	5606	3.99	0.07	335	0.55	0.00	11,646	4.99	0.14	157	0.43	0.00	13,168	3.60	0.16	84	0.58	0.00
	$0 < v_n \leq 90 \text{ m/s}$	15,102	10.74	0.18	6385	10.44	0.08	21,913	9.39	0.26	3278	8.92	0.04	33,297	9.11	0.39	78	0.54	0.00
	$v_n > 90 \text{ m/s}$	669	0.48	0.01	191	0.31	0.00	83	0.04	0.00	97	0.26	0.00	146	0.04	0.00	1	0.01	0.00
	$\text{SP} \geq 0.5$	11,943	8.49	0.14	5951	9.73	0.07	18,183	7.79	0.21	3065	8.34	0.04	22,125	6.05	0.26	64	0.44	0.00
Ninth	Harmless	2249	1.60	0.03	71	0.12	0.00	1851	0.79	0.02	50	0.14	0.00	1620	0.44	0.02	4	0.03	0.00
	$0 < v_n \leq 90 \text{ m/s}$	11,079	7.88	0.13	5011	8.20	0.06	18,876	8.09	0.22	2849	7.75	0.03	25,338	6.93	0.30	175	1.20	0.00
	$v_n > 90 \text{ m/s}$	1216	0.86	0.01	113	0.18	0.00	1784	0.76	0.02	52	0.14	0.00	3755	1.03	0.04	4	0.03	0.00
	$\text{SP} \geq 0.5$	9049	6.43	0.11	4729	7.73	0.06	12,442	5.33	0.15	2634	7.17	0.03	14,819	4.05	0.17	103	0.71	0.00
Eighth	Harmless	2159	1.53	0.03	385	0.63	0.00	1996	0.85	0.02	90	0.24	0.00	2544	0.70	0.03	0	0.00	0.00
	$0 < v_n \leq 90 \text{ m/s}$	13,206	9.39	0.16	6724	11.00	0.08	24,456	10.48	0.29	4011	10.91	0.05	35,757	9.78	0.42	207	1.42	0.00
	$v_n > 90 \text{ m/s}$	1403	1.00	0.02	35	0.06	0.00	3462	1.48	0.04	11	0.03	0.00	9268	2.54	0.11	0	0.00	0.00
	$\text{SP} \geq 0.5$	11,711	8.33	0.14	6321	10.34	0.07	13,522	5.79	0.16	3704	10.08	0.04	18,874	5.16	0.22	126	0.87	0.00
Seventh	Harmless	2974	2.11	0.04	319	0.52	0.00	3203	1.37	0.04	63	0.17	0.00	5041	1.38	0.06	8	0.05	0.00
	$0 < v_n \leq 90 \text{ m/s}$	13,162	9.36	0.16	4969	8.13	0.06	24,158	10.35	0.28	2159	5.87	0.03	43,160	11.81	0.51	65	0.45	0.00
	$v_n > 90 \text{ m/s}$	1032	0.73	0.01	1	0.00	0.00	3440	1.47	0.04	5	0.01	0.00	9057	2.48	0.11	1	0.01	0.00
	$\text{SP} \geq 0.5$	12,071	8.58	0.14	4749	7.77	0.06	17,038	7.30	0.20	2001	5.44	0.02	26,929	7.37	0.32	36	0.25	0.00
Sixth	Harmless	2785	1.98	0.03	381	0.62	0.00	2087	0.89	0.02	15	0.04	0.00	3275	0.90	0.04	0	0.00	0.00
	$0 < v_n \leq 90 \text{ m/s}$	9623	6.84	0.11	3184	5.21	0.04	11,885	5.09	0.14	845	2.30	0.01	18,351	5.02	0.22	1	0.01	0.00
	$v_n > 90 \text{ m/s}$	539	0.38	0.01	1	0.00	0.00	3000	1.29	0.04	1	0.00	0.00	1207	0.33	0.01	0	0.00	0.00
	$\text{SP} \geq 0.5$	8711	6.19	0.10	3063	5.01	0.04	8282	3.55	0.10	810	2.20	0.01	14,195	3.88	0.17	1	0.01	0.00
Fifth	Harmless	1967	1.40	0.02	65	0.11	0.00	665	0.28	0.01	21	0.06	0.00	219	0.06	0.00	0	0.00	0.00
	$0 < v_n \leq 90 \text{ m/s}$	4646	3.30	0.05	4785	7.83	0.06	4649	1.99	0.05	1718	4.67	0.02	1131	0.31	0.01	7	0.05	0.00
	$v_n > 90 \text{ m/s}$	350	0.25	0.00	0	0.00	0.00	1504	0.64	0.02	2	0.01	0.00	999	0.27	0.01	0	0.00	0.00
	$\text{SP} \geq 0.5$	4127	2.93	0.05	4624	7.56	0.05	2192	0.94	0.03	1642	4.47	0.02	968	0.26	0.01	5	0.03	0.00
Fourth	Harmless	5202	3.70	0.06	115	0.19	0.00	5787	2.48	0.07	33	0.09	0.00	1823	0.50	0.02	0	0.00	0.00
	$0 < v_n \leq 90 \text{ m/s}$	8021	5.70	0.09	6057	9.91	0.07	15,357	6.58	0.18	1938	5.27	0.02	17,564	4.81	0.21	222	1.52	0.00
	$v_n > 90 \text{ m/s}$	254	0.18	0.00	12	0.02	0.00	553	0.24	0.01	22	0.06	0.00	723	0.20	0.01	0	0.00	0.00
	$\text{SP} \geq 0.5$	7411	5.27	0.09	5393	8.82	0.06	11,720	5.02	0.14	1722	4.68	0.02	3321	0.91	0.04	22	0.15	0.00
Third	Harmless	4840	3.44	0.06	311	0.51	0.00	7866	3.37	0.09	576	1.57	0.01	6359	1.74	0.08	46	0.32	0.00
	$0 < v_n \leq 90 \text{ m/s}$	4205	2.99	0.05	7074	11.57	0.08	8848	3.79	0.10	2664	7.25	0.03	24,561	6.72	0.29	1416	9.73	0.02
	$v_n > 90 \text{ m/s}$	526	0.37	0.01	123	0.20	0.00	4149	1.78	0.05	39	0.11	0.00	3112	0.85	0.04	0	0.00	0.00
	$\text{SP} \geq 0.5$	3731	2.65	0.04	5940	9.71	0.07	6281	2.69	0.07	2510	6.83	0.03	10,785	2.95	0.13	1326	9.11	0.02
Second	Harmless	6703	4.77	0.08	1372	2.24	0.02	8021	3.44	0.09	5601	15.24	0.07	18,922	5.18	0.22	1774	12.19	0.02
	$0 < v_n \leq 90 \text{ m/s}$	3350	2.38	0.04	6497	10.63	0.08	7879	3.37	0.09	5557	15.12	0.07	13,813	3.78	0.16	2322	15.95	0.03
	$v_n > 90 \text{ m/s}$	669	0.48	0.01	656	1.07	0.01	3524	1.51	0.04	1409	3.83	0.02	10,017	2.74	0.12	5217	35.84	0.06
	$\text{SP} \geq 0.5$	2130	1.51	0.03	3632	5.94	0.04	4166	1.78	0.05	2133	5.80	0.03	1567	0.43	0.02	692	4.75	0.01

Table 5. Continued

	Case 1 ($d_p = 0.25 \mu\text{m}$)						Case 2 ($d_p = 0.50 \mu\text{m}$)						Case 3 ($d_p = 1.00 \mu\text{m}$)					
	PS		SS		PS		SS		PS		SS		PS		SS			
	N	$n_{\text{SIDE}} (\%)$	$n_{\text{hit}} (\%)$	N	$n_{\text{SIDE}} (\%)$	$n_{\text{hit}} (\%)$	N	$n_{\text{SIDE}} (\%)$	$n_{\text{hit}} (\%)$	N	$n_{\text{SIDE}} (\%)$	$n_{\text{hit}} (\%)$	N	$n_{\text{SIDE}} (\%)$	$n_{\text{hit}} (\%)$	N	$n_{\text{SIDE}} (\%)$	$n_{\text{hit}} (\%)$
First	Harmless	3133	2.23	0.04	250	0.41	0.00	2.90	0.08	1339	3.64	0.02	13,741	3.76	0.16	120	0.82	0.00
	$0 < v_n \leq 90 \text{ m/s}$	1170	0.83	0.01	1385	2.27	0.02	0.96	0.03	560	1.52	0.01	4753	1.30	0.06	2578	17.71	0.03
	$v_n > 90 \text{ m/s}$	148	0.11	0.00	0	0.00	0.00	0.36	0.01	0	0.00	0.00	5924	1.62	0.07	1	0.01	0.00
	$\text{SP} \geq 0.5$	712	0.51	0.01	1345	2.20	0.02	0.30	0.01	508	1.38	0.01	566	0.15	0.01	2227	15.30	0.03
Side	Harmless	45,622	32.43	0.54	4937	8.07	0.06	28.69	0.79	8368	22.76	0.10	95,970	26.26	1.13	2154	14.80	0.03
	$0 < v_n \leq 90 \text{ m/s}$	87,956	62.53	1.04	54,749	89.54	0.65	61.62	1.70	26,653	72.50	0.31	22,4984	61.55	2.65	7167	49.23	0.08
	$v_n > 90 \text{ m/s}$	7089	5.04	0.08	1460	2.39	0.02	9.69	0.27	1740	4.73	0.02	44,565	12.19	0.53	5237	35.97	0.06
	$\text{SP} \geq 0.5$	75,641	53.77	0.89	47,464	77.62	0.56	41.92	1.15	21,557	58.64	0.25	11,8486	32.42	1.40	4646	31.91	0.05

side of the blade for all the particles upon the impact and the ratio n_{hit} shows a global overview, in line with the fouling susceptibility criteria that consists of the ratio between the number of stuck particles and the total number of particles injected in the flow path.

In Table 5, N , n_{SIDE} , and n_{hit} related to the particles characterized by an SP equal to or greater than 0.5 are also reported. Finally, the rows grouped by the name *side* contain the sum of the values reported for each strip. With this global overview, it is possible to highlight the different behavior of particle deposition on the blade surface, in particular:

- the percentages of the particles with $v_n > 90 \text{ m/s}$ are higher for the strips close to the hub in the SS, while in PS this value is higher for the strips close to the blade tip;
- the comparison between cases 1 and 2 shows a different value of the SP > 0.5 between the PS and SS along the blade span. In fact, for the strip close to the hub (second), the particle percentage on SS of the SP > 0.5 is higher than the PS, while at midspan (sixth and seventh), the PS percentage is higher than the SS. At blade tip (11th), the two values are quite similar.

With the spanwise subdivision of the results shown in Table 5, we can highlight the difference in terms of particle–blade interaction behavior between the SS and PS.

Thanks to the sum of the values N , n_{SIDE} , and n_{hit} for the two sides of the blade, further analysis regarding particle–blade interaction is possible.

Figure 10 shows two bar charts relative to the sum of values for the ratio n_{SIDE} , reported for each strip, and indicated with the name n_{SIDE} . On the SS, the percentage of particles with SP > 0.5 is greater than the PS for all cases even if, for case 3, the phenomenon is much less obvious. This result shows how on the SS there are some fluid dynamic conditions that make it more sensitive to particle sticking. On the SS, there are fewer particles than the PS but these particles have a higher sticking capacity. From the compressor performance point of view, the sensitivity to fouling of the SS appears to be greater than the PS [25], thus a greater particle tendency to stick to the SS is an important result and focuses attention not only on the quantity of ingested contaminants but also to the fluid dynamic phenomena that characterize the flow around the blade. On the SS, case 1 is the most severe from a fouling point of view. The particles arrive with a normal impact velocity that makes it extremely effective in sticking to the blade surface. The percentage of particles with an SP > 0.5 reaches almost 90%. On the PS, the differences of the particle impact kinematics are less evident between the cases and all the percentages are quite similar to each other even if case 3 uses a particle diameter four times higher than case 1. The PS, for all cases, shows a higher percentage for the *harmless* particle category. This effect is directly related to the fact that in PS the separation that afflicts the SS does not take place. For the PS and SS, it can be seen that the particle percentage of the $v_n > 90 \text{ m/s}$ category increases with the increase of the particle diameter and the SP decreases as the diameter increases. This phenomenon is the precursor of the erosive effects that are produced by the particles with a diameter greater than $10 \mu\text{m}$, as reported in Ref. [16]. In fact, the normal impact velocity increases with the increase of the particle diameter and, in the same way, the particles become less able to stick, although the impact is more dangerous for the blade surface. The final analysis is related to the particles that have the SP > 0.5 . In particular, in Fig. 11, the trend of the ratio $n_{\text{hit,SP}>0.5}$ (black continuous line) for the particles with SP > 0.5 superimposed with the trend of the n_{hit} (grey dotted line) is reported. The two trends refer to both sides of the blade (PS and SS).

As mentioned above, n_{hit} represents the fouling susceptibility and its values represent a key result for gas turbine operators.

As can be seen from Fig. 11, for the PS, the trend of $n_{\text{hit,PS,SP}>0.5}$ does not follow the trend of $n_{\text{hit,PS}}$ unlike the trends reported for the SS. For the PS, the number of stuck particles is quite independent of the total number of particles that hit the blade and the $n_{\text{hit,PS,SP}<0.5}$ remains almost the same for the three

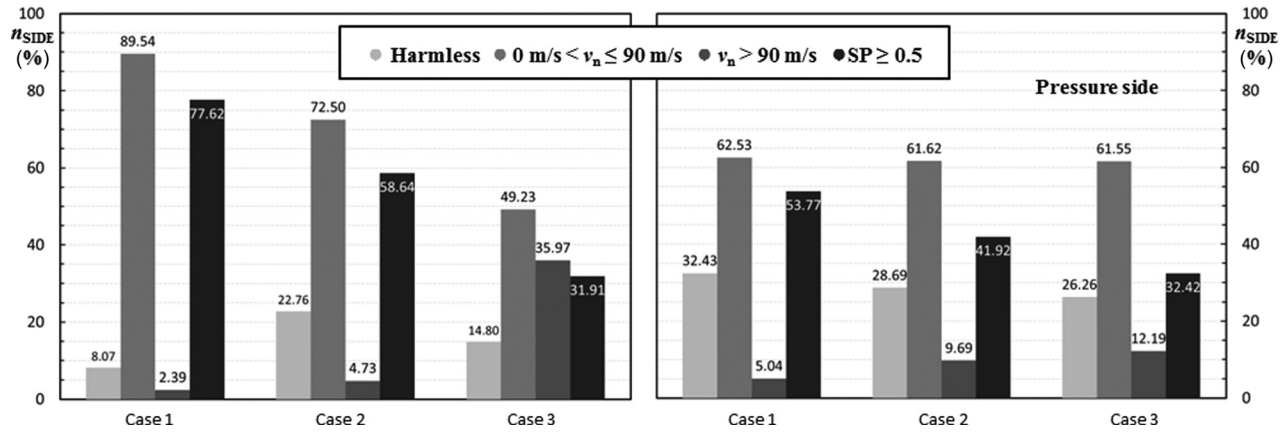


Fig. 10 Ratio n_{SIDE} for the SS and PS of cases 1, 2, and 3

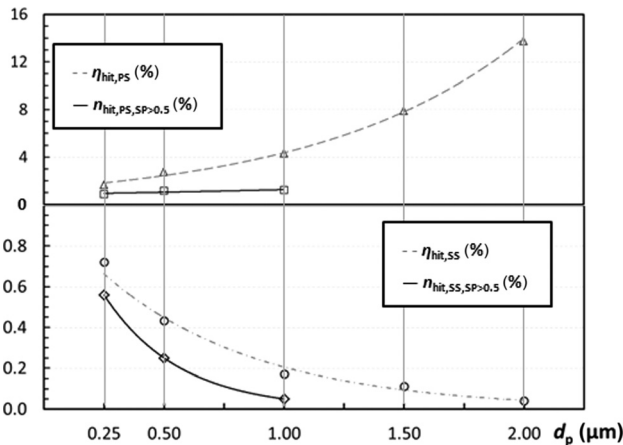


Fig. 11 Trends of the ratio $n_{hit,SP>0.5}$ and n_{hit} superimposed

718 considered cases. In this case, the higher particles produce more
 719 fouling effects due to their higher diameter and thus the higher
 720 mass. For the SS, the ratio $n_{hit,SS,SP>0.5}$ shows a very high percent-
 721 age of particles able to stick for the smallest diameters compared
 722 to the total number of particles that hit the SS.

723 Considering the different effect of the SS and PS deposits on
 724 axial compressor performance drop [25], some general guidelines
 725 for proper management of the power plant can be described as
 726 follows:

- 727 – as reported in Ref. [25], the deposits on the SS have the high-
 728 est influence on the axial compressor performance drop.
 729 Therefore, the filtration system must be designed to remove
 730 the smaller particles (up to $0.5 \mu\text{m}$) from the airflow stream
 731 because the bigger particles have a smaller SP and are not
 732 able to reach the SS due to their inertia;
- 733 – as reported in Ref. [25], the deposits on the PS have the low-
 734 est influence on the axial compressor performance drop.
 735 Therefore, the bigger particles that could stick on the PS do
 736 not determine a great performance drop and these deposits
 737 could be removed by proper periodic washing operations.

738 Conclusions

739 In this paper, an extended study on microparticle adhesion on
 740 the axial compressor blade surface was carried out. The micropar-
 741 ticles dragged by the airflow through the air filtration systems are
 742 responsible for compressor fouling if they come into contact with
 743 the compressor airfoils and stick there.

744 The kinematic characteristics for the impacting particles were
 745 obtained by a numerical model validated by data from literature
 746 and are used in order to describe for the first time how the sub

micro-sizes particles hit the blade surface. Special attention was
 given to the particle–blade interaction in terms of impact velocity
 and impact angle.

Thanks to experimental data reported in literature and the
 knowledge of the impact velocity component, it has been possible
 to highlight which blade areas are more affected by particle depo-
 sition and their sensitivity to particle diameters and fluid dynamic
 phenomena.

The key results can be summarized as follows:

- the particle impact velocity depends on several factors: shape
 of the blade (e.g., the blade height and local airfoil curvature),
 design characteristics (e.g., tip gap), and fluid dynamic phe-
 nomena (e.g., flow separation and tip leakage vortex);
- on the PS, the particle impact angle is very close to 90 deg in
 almost all of the airfoil extension;
- on the SS, the particle impact angle is lower than the PS due
 to the separation phenomena. This fact implies that particles
 hit the surface with a value of normal velocity higher than
 tangential velocity;
- thanks to the experimental results reported in literature regard-
 ing SP, it was possible to define two representative trends for
 the correlation between normal impact velocity and SP;
- on the SS, the smallest particles are the most numerous from
 a fouling point of view due to the high total number of par-
 ticles characterized by a SP greater than 0.5.

Regarding the management of gas turbine installations, the
 results of this study highlight the advantage of installing air filtra-
 tion systems that can remove small and very small particles from
 the air stream. This would allow the use of effective online wash-
 ing using larger droplets that typically would only hit and clean
 the PS of the blade.

The CFD numerical simulations link the design characteristic of
 the machine and the fluid dynamic phenomena. As shown in this
 work, these two items determine the particle deposition on the blade
 surface and thus the fouling phenomena. Future studies would have
 to analyze the behavior of a subsonic stage where the fluid dynamic
 phenomena are quite different compared to a transonic stage.

An increase in the knowledge of fouling through the use of
 numerical codes may therefore constitute a decisive element for
 better planning of maintenance of turbomachinery. In this sense,
 studies (experimental and numerical) dedicated to the interaction
 between the particles responsible for fouling (in terms of size and
 material) with blade surfaces are fundamental to allow for better
 simulations with numerical codes.

Nomenclature

- A = area
- d = diameter

794 E = Young module
 795 k = turbulent kinetic energy
 796 m = mass flow rate
 797 M = mass
 798 n = ratio
 799 N = total number (referred to particles)
 800 r = radius
 801 St = Stokes number
 802 t = operation time
 803 u = velocity component
 804 U = averaged velocity
 805 v = velocity
 806 X = particle concentration (blade)
 807 x, y, z = axis coordinate
 808 α = impact angle
 809 β = compression ratio
 810 ε = dissipation rate of turbulent kinetic energy
 811 η = efficiency
 812 μ = dynamic viscosity
 813 ρ = density
 814 σ = standard deviation

815 **Subscripts and Superscripts**

816 b = blade
 817 c = contaminant
 818 h = hydraulic
 819 hit = hit (referred to blade, side, and slice)
 820 i = impact
 821 n = normal direction
 822 p = particle
 823 $SIDE$ = side (referred to the blade division)
 824 $STRIP$ = strip (referred to spanwise division)
 825 t = tangential direction
 826 TT = total-to-total
 827 x, y, z = axis coordinate
 828 1 = inlet
 829 $-$ = average
 830 \sim = weighted-area average
 831 $*$ = peak

832 **Acronyms**

833 AR = accretion rate
 834 CFD = computational fluid dynamics
 835 DRW = discrete random walk
 836 FDS = flux-difference splitting
 837 LE = leading edge
 838 PS = pressure side
 839 SEM = scanning electron microscope
 840 SP = sticking probability
 841 SS = suction side
 842 TE = trailing edge

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