

A new 3D experimentally consistent XFEM to simulate delamination in FRP-reinforced concrete

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Abstract

FRP-based reinforcements are increasingly used for the structural rehabilitation of buildings damaged by seismic loadings. Single-lap shear tests on FRP-reinforced concrete blocks are often performed to assess the maximum transferable load before delamination. In this paper, an effective 3D regularized extended finite element model (XFEM) is proposed to study single-lap shear tests. The bending of the FRP plate, the influence of FRP plate width, and the two-way delamination for variable bonding length are assessed. Based on a suitable choice of the level sets associated with the XFEM enrichment, the proposed model can be used for design purposes besides, or in alternative to, experimental tests.

Keywords: A: 3-Dimensional reinforcement, B: delamination, C: Finite element analysis (FEA), C: Computational modelling, XFEM

1. INTRODUCTION

Fiber Reinforced Polymers (FRP) reinforcements are increasingly used for the post-seismic structural rehabilitation and to enhance the ultimate strength of concrete structural elements. Single-lap shear tests on FRP-reinforced concrete blocks are often performed to assess the ultimate load before the complete delamination of the FRP plate [1] occurs. Delamination is a complex three-dimensional (3D) process, and a challenging issue in computational mechanics [2, 3]. To study single-lap shear tests, we present an effective eXtended

18 Finite Element Model (XFEM) based on an experimentally consistent level set system. Key
19 aspects, such as the peeling and the bending of the FRP plate and the influence of the FRP
20 plate length and width, are investigated.

21 Based on experimental results, it is known that the delamination process involves a thin
22 portion of material localized in the concrete layer underlying the FRP plate. Within finite
23 element (FE) approaches, the concrete block is commonly modeled as a two-dimensional
24 (2D) body in plane stress state, and the FRP-concrete interface as a one-dimensional in-
25 terface subjected to distributed tangential stress by means of bond stress-slip laws [4, 5].
26 However, even if the reinforcement is loaded mainly in shear, the out-of-plane displacements
27 observed during the tests induce tensile and compression stresses orthogonal to the bond-
28 ing plane [6]. Hence, one-dimensional bond stress-slip relationships have been modified to
29 take into account failure mechanisms relying on mixed fracture modes [7, 8]. Furthermore,
30 experimental tests [9, 10] detected the occurrence of edge regions with high shear strains
31 associated with the stress transfer from FRP to concrete, and highlighted the key role of
32 the peel displacement in the delamination onset and propagation. In particular, a two-way
33 debonding mechanism was observed [6]. For very short bonding lengths, debonding starts
34 at the free end of the FRP plate and propagates towards the loaded end of the plate. For
35 sufficiently long bonding lengths, debonding starts at the loaded end and propagates towards
36 the free end until a critical bonding length is reached, at which delamination at the free end
37 starts and propagates towards the loaded end. Furthermore, the maximum transferable load
38 is influenced by the FRP plate bending, and the ratio between the widths of the plate and
39 the concrete block.

40 Width-effects associated with shear-strain-edge-regions cannot be taken into account
41 assuming a plane stress state. For this reason, 3D FE codes [12] have been used, most of them
42 assuming the concrete-FRP surface as a zero thickness interface. Based on this assumptions,
43 3D FE simulations highlighted the dependence of the concrete stress state on the bending
44 rigidity of the FRP plate and the edge stiffening effect on the shear stress components. To
45 this purpose, the analysis in reference [11] assumed elastic materials, and a perfect bond
46 between the adhesive and the FRP plate. An elastic-damage interface model governing the
47 interlaminar stresses acting in the sliding direction was adopted in a more recent study [12].
48 Furthermore, based on a bond stress-slip law, Neto et al. [13] have introduced an effective
49 bond width that comprehends of both the FRP plate and the contiguous concrete with

50 non-vanishing shear stress.

51 Delamination in composite laminates has been effectively simulated through the XFEM [14],
52 a partition of unity FE method proposed by Belytschko and coworkers [15]. This method
53 incorporates features of the expected solution at the nodal level. Unlike standard FE mod-
54 els, the XFEM makes it possible to adopt meshes, that are independent of the geometry
55 of cracks and interfaces. In the XFEM, interfaces are usually introduced in an implicit
56 way as level set functions [16]. Applications of the XFEM are really wide. For instance,
57 the XFEM was exploited to simulate fracture of composite laminates [2], and the dam-
58 age progression in laminated overheight compact tension specimens using representations of
59 individual cracks [17], or by means of an XFEM formulation based on a discrete damage
60 zone model [18]. The phantom-node method, a variation to the XFEM, has been used for
61 a mesh-independent 3D representation of matrix cracks as straight discontinuities in the
62 displacement field [19]. While the use of the XFEM for the 3D modelling of delamination
63 in composite laminates is quite established, 3D XFEM simulations of pull-out tests of FRP
64 plates bonded to concrete specimens seem not to have been presented in the literature.

65 We propose here an original 3D regularized version of the XFEM method with a global-
66 local level set system that is able to take into account edge effects. This regularized XFEM
67 approach was previously developed by the Authors for 2D plane-stress simulation of de-
68 lamination [20] to study, for instance, the transition from strain localization to crack [21],
69 and inclusions with imperfect interfaces [22]. In particular, the delamination strength was
70 predicted with a sufficient accuracy exploiting just the nominal values of the Young moduli
71 and the Poisson coefficients of concrete, glue, and reinforcement. The robust continuous-
72 discontinuous transition ensured mesh-size independent, energy-consistent structural results,
73 and avoided the sudden loss of stiffness that frequently occurs with other continuous-
74 discontinuous procedures [23]. While complex interface damage laws have been proposed in
75 the literature [24, 25], in the regularized XFEM approach to delamination [20], the damage
76 evolution was associated with both shearing and normal opening through a simple Rankine-
77 type law. Very reliable results and an excellent agreement with the experimental results
78 were reached.

79 With respect to the previous 2D model [20], the 3D regularized XFEM formulation
80 proposed herein is characterized by a global-local system of level sets taking into account
81 the edge effects and a fully 3D mixed shearing-peeling debonding mechanism. The main

82 aim of this contribution is not only to fit the experimental results, that are often subjected
 83 to several uncertainties, but also to provide a reliable technical tool that can be used for
 84 design purposes besides, or in alternative to, experimental tests.

85 After a basic introduction to the computational model in Sec. 2, the results obtained
 86 for the data sets [6] and [26] are shown in Sec. 3. Emphasis is put on the bending of the
 87 FRP plate, the presence of the peeling displacements at both the loaded and the free ends
 88 of the FRP plate, and the two-way delamination depending on the bonding length. The 3D
 89 behavior of the delamination process is assessed in Sec. 4, where the width dependence of
 90 the maximum transferable load is also displayed.

91 **2. Introduction to the computational model**

92 To model correctly the delamination process, an effective computational formulation
 93 has to: i) tackle the possible mesh-size dependency induced by the adoption of a softening
 94 constitutive law for concrete; ii) ensure a smooth continuous-discontinuous transition where
 95 delamination occurs; iii) reproduce the 3D debonding of the FRP plate taking into account
 96 the edge effects. In previous studies, we have proved that the regularized XFEM approach
 97 satisfies the mesh-independency requirement [22, 21], and ensures a smooth continuous-
 98 discontinuous transition [21]. Hereinafter, only the aspects related to the 3D debonding of
 99 the FRP plate taking into account the edge effects will be assessed. Sec. 2.1 focuses on
 100 the general aspects of the regularized XFEM approach, such as the regularized kinematics,
 101 and a smooth mechanically-consistent continuous-discontinuous-transition procedure. 3D
 102 modelling of delamination of FRP plates from concrete blocks requires a specific strategy
 103 different from that adopted to study 2D delamination, as discussed in Sec. 2.2.

104 *2.1. Kinematics, constitutive laws and continuous-discontinuous computational procedure in* 105 *the regularized XFEM*

106 Let the displacement field \mathbf{u} be discontinuous across the delamination surface $\mathcal{S} \in \mathbb{R}^3$ of
 107 normal \mathbf{n}_S . Within a single element of nodal degrees of freedom \mathbf{U}^e and \mathbf{A}^e interpolated by
 108 the usual FE interpolation functions \mathbf{N}^e , the regularized XFEM displacement \mathbf{u}^e of element
 109 e is [15]

$$110 \quad \mathbf{u}^e = \mathbf{N}^e \mathbf{U}^e + \mathcal{H}_\rho \mathbf{N}^e \mathbf{A}^e, \quad (1)$$

111 where \mathcal{H}_ρ is a regularized Heaviside function that approximates the Heaviside function for
 112 vanishing regularization length ρ . The vector \mathbf{A}^e collects the jump components along x , y
 113 and z for the finite element e . \mathcal{H}_ρ is assumed a function of the distance from the *global* level
 114 set plane. By compatibility, the strain field is

$$115 \quad \boldsymbol{\varepsilon}^e = \mathbf{B}^e \mathbf{U} + \mathcal{H}_\rho \mathbf{B}^e \mathbf{A}^e + \|\nabla \mathcal{H}_\rho\| (\mathbf{N}^e \mathbf{A}^e \otimes \mathbf{n}^e), \quad (2)$$

116 where $\mathbf{B}^e = \nabla \mathbf{N}^e$ is the standard FE compatibility matrix. Usually, \mathbf{n}^e denotes the vector
 117 normal to the surface across which the displacement field exhibits a discontinuity [15]. In
 118 our approach, we start from the experimental evidence [6] that the crack orientation in the
 119 concrete below the FRP plate is variable in the space, namely it is different at the boundaries
 120 of the detachment surface and within the detachment surface (Fig. 2). Therefore, \mathbf{n}^e is
 121 assumed to change from element to element according to a *local level set* attached to each
 122 Gauss point, as described in Sec. 2.2.

FRP and adhesive are modelled as linear elastic materials. At each Gauss point, we
 introduce the damage variables D and D_c for the concrete and the zone where debonding
 occurs, respectively, and compute the associated stresses thorough

$$\boldsymbol{\sigma}^e = (1 - D) \mathbf{E} \mathbf{B}^e \mathbf{U} + \mathcal{H}_\rho (1 - D) \mathbf{E} \mathbf{B}^e \mathbf{A}^e, \quad (3a)$$

$$\boldsymbol{\sigma}_c^e = (1 - D_c) \mathbf{E}_c \|\nabla \mathcal{H}_\rho\| (\mathbf{N}^e \mathbf{A}^e \otimes \mathbf{n}^e). \quad (3b)$$

123 Note that the stress $\boldsymbol{\sigma}_c^e$ (3b) is computed taking into account the *local level set system*
 124 described in Sec. 2.2. In particular, the concrete damage is governed by an exponential
 125 Rankine elasto-damaging law until D has reached the critical value D^{cr} [22]. As soon as
 126 $D \geq D^{cr}$, the evolution of D is dropped, i.e. the concrete can only elastically unload: a
 127 regularization zone replaces the discontinuity. In particular, the damage evolution can be
 128 D_c affecting the stress \mathbf{t} is

$$129 \quad D = \min\{D^{cr}, f(r)\}, \quad D_c = \max\{D^{cr}, f(r)\}, \quad f(r) = 1 - \frac{r_0}{r_c} \exp\left(-2H \frac{r_c - r_0}{r_0}\right), \quad (4)$$

130 where $r_c \geq r_0$, with $r_0 = f_t$ for tensile damage, and $r_0 = f_c$ for compressive damage. For a
 131 monotonic damage process in a one-dimensional bar the stress-strain law obeying Eq. (4) is
 132 shown in Fig. 1 for $H = 0.008$ MPa (blue continuous line) and $H = 0.005$ MPa (red dotted

133 line) with fixed $f_t = 3.21$ MPa.

134 The main differences between the proposed regularized XFEM procedure and standard
 135 applications of the XFEM are synthesized in Tab. (1). In particular, the standard elemental
 136 stiffness \mathbf{K}_{st}^e is

$$137 \quad \mathbf{K}_{st}^e = \begin{bmatrix} \mathbf{B}^{et} (1 - D) \mathbf{E} B_u & \mathbf{B}^{et} (1 - D) \mathbf{E} \nabla(\mathcal{H}_\rho \mathbf{N}^e) \\ \nabla(\mathcal{H}_\rho \mathbf{N}^e)^t (1 - D) \mathbf{E} \mathbf{B}^e & \nabla(\mathcal{H}_\rho \mathbf{N}^e)^t (1 - D) \mathbf{E} \nabla(\mathcal{H}_\rho \mathbf{N}^e) \end{bmatrix}, \quad (5)$$

138 with $\nabla(\mathcal{H}_\rho \mathbf{N}^e) = \mathcal{H}_\rho \mathbf{B}^e + \nabla \mathcal{H}_\rho \otimes \mathbf{N}^e$. The adopted stiffness matrix is

$$139 \quad \mathbf{K}_{reg}^e = \begin{bmatrix} \mathbf{B}^{et} (1 - D) \mathbf{E} B_u & \mathbf{B}^{et} \mathcal{H}_\rho (1 - D) \mathbf{E} \mathbf{B}^e \\ \mathbf{B}^{et} \mathcal{H}_\rho (1 - D) \mathbf{E} \mathbf{B}^e & \mathbf{B}^{et} \mathcal{H}_\rho^2 (1 - D) \mathbf{E} \mathbf{B}^e + \|\nabla \mathcal{H}_\rho\| \bar{\mathbf{N}}^e{}^t (1 - D_c) \mathbf{E}_c \bar{\mathbf{N}}^e \end{bmatrix}, \quad (6)$$

140 where the operator $\bar{\mathbf{N}}^e$ is such that $\nabla \mathcal{H}_\rho \otimes \mathbf{N}^e A \approx \|\nabla \mathcal{H}_\rho\| \bar{\mathbf{N}}^e \mathbf{A}_e$. The differences between
 141 the stiffness matrices (5) and (6) stem from the adopted variational formulation, which has
 142 been thoroughly described in [28, 27].

143 2.2. Definition of the fracture process zone based on global-local level set system

144 In the studied pull out tests, a concrete layer like that shown in Fig. 2 remains attached
 145 to the delaminated FRP plate [6], being thicker at the ends of the bonded zone, and in cor-
 146 respondence with large concrete aggregates. In this case, a shear process zone with cracks
 147 orthogonal to the maximum principal stress, thus inclined with respect to the adhesive sur-
 148 face, is often assumed [29]. In the previous applications of the regularized XFEM approach,
 149 a unique global level set was defined representing, for instance, the implicit surface of the
 150 inclusion in the study of particulate composites [28], or the crack path in strain localization
 151 problems [22, 21]. The normal vector \mathbf{n} coincided with the geometric normal to the inter-
 152 face surrounding the inclusion. In the present study, a different level set system is adopted,
 153 called *global-local level set* system. This is the main original aspect of the adopted approach
 154 with respect to the previous applications of the regularized XFEM approach. In this study,
 155 delamination is assumed to take place in a plane parallel to the adhesive layer, as usual
 156 in the literature [5, 3]. Such a plane has been located at 1 mm underneath the concrete
 157 surface. This corresponds to confine the delamination process in the finite element layer
 158 placed immediately below the glue. The hypothesis agrees with the experimental evidences
 159 in the central part of the plate (Fig. 2), where a layer of concrete 1 to 3 mm thick is usually

160 detached [6]. Of course, the concrete bulbs at the ends of the bonded zone are not caught.

161 The *global-local level set* procedure relies in the fact that we assume: a *global* level set
162 function denoting the surface of the delamination that is parallel to the adhesive surface
163 and located below the concrete surface, and a *local* level set system at each Gauss point
164 of the XFEM enriched finite elements, displayed in Fig. 3. The global level set surface is
165 defined a priori parallel to the adhesive layer, as usual in two-dimensional FE analysis and
166 in numerical models based on one-dimensional interface laws.

167 As for the *local* level set system, we have adopted the following path of reasoning. In
168 principle, the \mathbf{n}^e vector in the constitutive equation (3b) is orthogonal to the crack direction
169 and governed by the maximum tensile principal stress. This means that the computation of
170 the principal stresses has to be performed at each of the Gauss points for each equilibrium
171 iteration per each enriched finite elements. Therefore, hundreds of equilibrium iteration are
172 sometimes required to get convergence, the resulting computational procedure would turn
173 out being rather CPU demanding. Instead of proceeding in this way, in this study, we have
174 adopted the simplified *local* level set shown in Fig. 3. Each edge is associated with a local set
175 whose normal vectors are shown in red and blue. The blue normals are oriented at $(0, 0, -1)$
176 for $x > 0, y > 0, z > 0$. The red normals are oriented outwards the edges along the direction
177 $(1/\sqrt{3}, 1/\sqrt{3}, -1/\sqrt{3})$ for $x > 0, y > 0, z > 0$. The local level set adopted inside the enriched
178 layer is associated with a normal field, plotted in green in the same figure, oriented towards
179 the symmetry plane $y - z$ at $(-1/\sqrt{3}, 1/\sqrt{3}, -1/\sqrt{3})$ for $x > 0, y > 0, z > 0$, \mathbf{n} being
180 specular with respect to the $y - z$ plane in the domain $x < 0, y > 0, z > 0$. This assumption
181 ensures that the delamination process is captured consistently. Noteworthy, the adopted
182 *global-local level set* system is imposed a priori, but the final directions of the jump vectors
183 \mathbf{A}^e at the nodes are computed by the nonlinear iterative-incremental procedure and will not
184 generally coincide with that of the superimposed normals. Moreover, the enriched layer of
185 elements has the same width as the FRP plate, and is located 1 mm below the adhesive
186 layer.

187 **3. FRP plate bending, peeling and twofold delamination onset**

188 This section is devoted to the study of the influence of the FRP plate bending and
189 illustrates the occurrence of peeling at both the ends of the FRP plate depending on the
190 bonding length. The values of the Young modulus and the Poisson coefficient adopted for

191 the FRP plate, the glue, and the concrete used in the model are collected in Tab. 2 as
192 indicated in the references of the experimental works.

193 The regularized XFEM code used in this paper is a development of a Fortran 3D code
194 for elastic materials and standard XFEM [28].

195 After an introduction to the geometry and material properties in Sec. 3.1, Secs. 3.2
196 and 3.3 analyze the strain in the concrete substrate, the peeling. Moreover, the forthcoming
197 sections investigate the influence of the bonding length and the width of the FRP plate on
198 both the structural response and the delamination activation.

199 *3.1. Experimental tests geometry*

200 For a comparison with experimental results, we will consider the experimental campaign
201 of Chajes et al. [26], often studied in the literature, and the tests by Carrara et al. [6]. It
202 is remarkable that, while for the former data set the structural response has been recorded
203 up to the maximum load, for the latter data set, the whole load-displacement curves with
204 softening post-peak branches have been recorded, owing to a stable loading control procedure
205 developed for this purpose. All the data of the geometry and the material parameters of
206 the original tests are collected in Tab. 2. In the geometry adopted by Chajes et al. [26]
207 shown in Fig. 4a, the glue layer starts at the front of the specimen, namely close to the
208 loaded end of the FRP plate. On the contrary, in the geometry of Carrara et al. [6], shown
209 in Fig. 4b, the FRP plate is bonded at a certain distance from the front of the specimen
210 to avoid the detachment of a concrete wedge when pulling the FRP plate. Moreover, the
211 softening modulus H has been set equal to 0.005 MPa for the Chajes et al. [26] test, and
212 equal to 0.008 MPa for the Carrara et al. [6] test.

213 Based on symmetry reasons, only one half of the specimen has been meshed. Considering
214 that the concrete block stiffness plays an important role on the structural response of the
215 specimen, we have first run some simulations where the real dimensions of the concrete
216 block were modeled. Then, to reduce the computational burden, we have run the same
217 simulations but with a width of the meshed geometry of the concrete block equal to 80 mm.
218 These simulations gave the same results of the case with the entire width of the concrete
219 block. Also the height of the concrete block has been taken equal to 35 mm to reduce the
220 computational effort. These results seem to be confirmed also in [31]. The mesh adopted
221 for Chajes et al. [26] in the $L_b = 101.6$ mm case shown in Fig. 5 is made of 250548 elements
222 with 243930 total degrees of freedom.

223 *3.2. Structural response, delamination onset and propagation*

224 As the complete data set of Carrara et al. [6] test is available up to the final delamination
225 stage is available, the corresponding delamination process has been assessed for various
226 bonding lengths.

227 To compare the computed structural responses with the experimental ones, we have
228 plotted the load vs both the z displacement u_1 at the loaded end and the z displacement
229 u_2 at the free end of the FRP plate. Different bonding lengths L_b have been considered,
230 from the short length $L_b = 30$ mm shown in Fig. 6a, to the longer lengths $L_b = 90$ mm and
231 $L_b = 120$ mm, displayed in Figs. 6b and 6c, respectively. The snap-back in the structural
232 response plotted in terms of the front-displacement increases as the bonding length increases,
233 as shown in the case of $L_b = 90$ mm in Fig. 6b, and in the case of $L_b = 120$ mm in Fig. 6c.

234 The maximum transferable loads computed with the present analysis are shown for
235 variable bonding length in Fig. 7a. In addition to the previous load-displacement simulations,
236 the structural responses up to the maximum transferable load have been determined for the
237 Chajes et al. [26] geometry too. For this latter case, the maximum transferable loads vs the
238 bonding length are shown in Fig. 7b.

239 The different types of structural behavior for the Carrara et al. [6] tests shown in Fig. 6
240 are associated with different positions of the delamination onset. Because the detachment is
241 strictly related to the evolution of the damage variable D_c , it is thus convenient to consider
242 the evolution of D_c , which is the damage on the concrete surface below the glue. Figs. 8
243 display the sequence of the contour plots of D_c for $L_b = 30$ mm and $L_b = 90$ mm at the pre-
244 peak, the peak and the post-peak loads, from top to bottom. In the case of short bonding
245 length, $L_b = 30$ mm, damage starts at the free end. On the contrary, for the $L_b = 90$ mm
246 case displayed in Fig. 8b, the damage evolution is different, as shown in Fig. 9 for the pre-
247 peak, peak and post-peak stages. In a first stage, the damage develops at both the loaded
248 and the free end of the FRP plate (a, b), it subsequently propagates from the loaded end
249 towards the free end (c) up to a certain value of the actual bonded length, approximately
250 equal to 30 mm (d), then it propagates from the free to the loaded end (e) up to complete
251 delamination (f). An animation of the damage propagation for the $L_b = 90$ mm case is also
252 shown in the supplementary file.

253 For all the investigated bonding lengths, the bending of the FRP plate influences the
254 damage onset and its evolution. Such an influence is shown in Fig. 10 for the Chajes et

255 al. [26] test with bonding length $L_b = 50.8$ mm. The same figure shows that, for this
 256 bonding length, delamination proceeds from the free end to the loaded end of the FRP
 257 plate.

258 To complete the picture of the 3D delamination process, the tangential stress τ_{yz} and the
 259 peel stress σ_{yy} have been evaluated at the Gauss points on the concrete surface below the
 260 glue for the Carrara et al. [6] test. Figs. 11 display the spatial frames of σ_{yy} and τ_{yz} stress
 261 components, on the left and on the right column, respectively, detected at the pre-peak, the
 262 peak and the post-peak for the $L_b = 90$ mm case. In particular, in a first stage, the peel
 263 stress σ_{yy} is activated at the loaded end (a) and subsequently at the free end (b, c, d). In
 264 the final stages (e, f), peeling at the loaded end is predominant. The shearing stress τ_{yz} first
 265 propagates from the loaded end to the free end (a,b), then its peak moves from the free to
 266 the loaded end (c, d, e). Moreover, the shear stress τ_{yz} evaluated at the end of the FRP plate
 267 reaches values of 15 MPa, which are slightly high with respect to the experimental ones. The
 268 profile of τ_{yz} cannot be directly compared to the "shear stress" of one-dimensional shear-
 269 stress-slip laws [5], which are average values obtained from the variation of two subsequent
 270 strains measured by strain gauges on the plate surface. For instance, an example of shear-
 271 stress-slip law equivalent to the experimental one was deduced from the results obtained by
 272 means of the regularized XFEM approach in the 2D case [20].

273 3.3. Strain evolution and deformability

274 The strain evolution for various load levels have been compared with the available ex-
 275 perimental results for both the experimental data sets. For this purpose, we have ob-
 276 tained from our 3D results a shear strain equivalent to that usually obtained in experi-
 277 ments from the displacements recorded at the strain gauges. In particular, the expression
 278 $\varepsilon = (u_{z,i+1} - u_{z,i})/\Delta z_i$, has been exploited, where z is the longitudinal axis, and $u_{z,i+1}$
 279 and $u_{z,i}$ denote the displacements recorded on the FRP plate at discrete positions z_i and
 280 $z_{i+1} = z_i + \Delta z_i$. In Fig. 12, the evolution of the strain along the z -axis is shown for the
 281 Chajes et al. [26] tests at the maximum transferrable load. The post-peak strain profiles
 282 have not been reported, because the experimental post-peak data are not available. Profiles
 283 of the same colors correspond to the same load level, while the markers indicate the experi-
 284 mental results. For the Carrara et al. [6] experimental data, the strain profiles corresponding
 285 to both the pre-peak and the post-peak branches are available. In Fig. 13 the strain profiles

286 computed for the two bonding lengths $L_b = 90$ mm and $L_b = 120$ mm are compared with
287 the homologous experimental profiles.

288 A good agreement between numerical values and tests is confirmed for both the experi-
289 mental campaigns.

290 4. 3D aspects: Width and edge effects

291 To highlight the 3D aspects of delamination, this section investigates the profiles of
292 several relevant stress components along the width and the length of the FRP-plate. In
293 Sec. 4.1, the influence of the FRP plate width on the structural response is assessed. For
294 this purpose, the Chajes et al. [26] tests are taken into account. In Sec. 4.2, the main results
295 obtained for the cases investigated, are discussed.

296 4.1. Influence of the FRP plate width

297 The delamination analysis has been performed for variable FRP plate widths b . For
298 $b/B > 0.5$, the width of the concrete is insufficient to allow a full transmission of the stresses
299 from the FRP plate to the concrete substrate [10]. Therefore, the concrete width of the
300 specimen analyzed has been set equal to 80 mm. In particular, the peel stress σ_{yy} and the
301 shearing stress τ_{yz} have been plotted at the Gauss points of the finite elements within the
302 bonded concrete. They are displayed in Fig. 14 for two widths of the FRP plate, namely
303 $b = 15$ mm and $b = 45$ mm for the same bonding length $L_b = 101.6$ mm. The profiles have
304 been detected at the peak of the transferable load. Fig. 14 display no appreciable edge effect
305 for both the values of b .

306 The profiles of the shearing and the axial strain components ε_{xz} and ε_{zz} , respectively,
307 have been plotted across the width at different locations along the FRP plate length for
308 $b=15$ mm, $b=45$ mm, and $L_b = 101.6$ mm. Figs. 15 show the values detected during the
309 elastic stage, before that the delamination process starts, while Figs. 16 display the evolution
310 of these specific strain components during the delamination process. The profiles correspond
311 to $z = 0.5$ mm (cyan dotted line), $z = 18.1$ mm (pink dashed line), $z = 38.1$ mm (yellow
312 dash-dotted line), $z = 58.1$ mm (green dotted line), $z = 78.1$ mm (red dashed line) and
313 $z = 95.1$ mm (blu continuous line). During the elastic stage, the profiles of ε_{xz} and ε_{zz}
314 extend over a region significantly larger than b (Fig. 15). This confirms that the modelled
315 concrete support must be sufficiently large compared to the FRP plate width to allow a

316 full diffusion of the shearing stress components in the concrete surrounding the FRP plate,
 317 as observed by Subramaniam et al. [10]. While ε_{zz} at the center of the FRP plate are
 318 almost constant, ε_{xz} is antisymmetric with respect to the symmetry plane. Both of them
 319 display high variations over an edge region approximately 20 mm wide. When the load
 320 reaches 90% of the peak load, the damage has been activated at the edges of the concrete
 321 substrate underlying the FRP plate. Correspondingly, the ε_{xz} and ε_{zz} shown in Fig. 16
 322 display narrower edge regions of width approximately equal to 10 mm, in agreement with
 323 the experimental results reported in [6].

324 Furthermore, the dependence of the strain profiles on the bonding length has been as-
 325 sessed. Figs. 17 and 18 display the strain components during the elastic stage and at the
 326 peak, respectively, for the bonding lengths $L_b = 152.4$ mm and $L_b = 202.3$ mm. These
 327 bonding lengths are close to the asymptotic value predicted by Fig. 7b. The strain profiles
 328 have been evaluated at a distance from the front of the concrete block of 25.4 mm (blue line),
 329 50.8 mm (red line), and 76.2 mm (green line). The results show that the strain diffusion in
 330 the concrete surrounding the FRP is substantially independently of L_b .

331 The widths of the edge regions computed through the present model are smaller than
 332 those experimentally detected by Subramaniam et al. [10], who measured edge regions ap-
 333 proximately 20 mm wide for both ε_{xz} and ε_{zz} . Analogously to Subramaniam et al. [10], the
 334 width of the edge regions measured in this study is independent of b . These widths have
 335 not been imposed, and have been observed after post-processing the results. Indeed, the
 336 enriched layer of elements where the debonding is simulated has a width equal to that of
 337 the FRP plate.

338 In Fig. 19a, the dependence of the maximum stress $\sigma_u = P/(bt_f)$ transferable through
 339 the FRP plate width is displayed. The same figure displays the results obtained with the
 340 CNR-design formula [32]

$$341 \quad P_{max} = b_f \sqrt{2 E_f t_f \Gamma_F} \quad (7)$$

342 where Γ_F is the specific fracture energy formulated as

$$343 \quad \Gamma_F = k_b k_G \sqrt{f_{cm} f_{ctm}} \quad (8)$$

344 where $k_b = \sqrt{\frac{2-b_f/b}{1+b_f/b}} \geq 1$ for $b_f/b \geq 0.25$, and $k_b = 1.1832$ for $b_f/b < 0.25$, and k_g is usually
 345 taken in a range going from 0.063 for preformed composites to 0.077 for in situ impreg-

346 nated composites. The results of our simulations are perfectly fitted by the CNR rule (7)
347 using $k_b = 1.1832$ and $k_g = 0.065$, and the nominal properties of the concrete and the FRP
348 that are, $f_{cm}=43$ MPa, $f_{ctm}=3.21$ MPa and $t_f=1$ mm, $L_b=101.6$ mm, $E_f=108380$ MPa.

349 4.2. Discussion

350 The agreement with the CNR rule [32] shown in Fig. 19a confirms that the developed
351 computational model is a reliable design tool. The computed trend of the maximum trans-
352 missible stress shown in Fig. 19a is decreasing with increasing FRP plate width b and tends
353 to an asymptote for large widths. On the contrary, based on their experimental tests, where
354 the concrete widths B equal to 52 mm and to 125 mm were taken, Subramaniam et al. [9, 10]
355 concluded that, below the critical value 0.5 of the ratio b/B , the maximum transmissible
356 stress increases for increasing width of the concrete support. On the other hand, Fig. 19b
357 shows that the maximum transferable load increases with the width b . We have investigated
358 also larger widths B , but the same decreasing trend has been found. A more extensive
359 numerical campaign is necessary to understand whether larger ratios b/B , namely larger
360 concrete supports, may correspond to an increasing ultimate stress for increasing b . Fur-
361 thermore, the occurrence of peel stress can potentially affect the results for relatively small
362 bonding lengths, such as in the current case. To summarize, the results presented in Sec. 4
363 have shown that:

- 364 • Where the FRP plate is still attached to the concrete substrate, the profiles of the
365 strain components ε_{xz} and ε_{zz} along the length of the FRP plate are not uniform
366 across the width, and display high gradients localized at edge regions of the FRP plate
367 comprehensive of both the FRP plate edges and the surrounding concrete.
- 368 • Where delamination has been activated, the edge regions, intended as the regions with
369 high gradients, corresponding to the ε_{xz} and ε_{zz} profiles are localized in a narrower
370 zone.
- 371 • The edge regions corresponding to the ε_{xz} and ε_{zz} profiles have a width independent
372 of b and L_b .
- 373 • For the investigated geometries, the maximum transferable load increases for increasing
374 width b of the FRP plate.
- 375 • For the assumed b/B ratios, the maximum load increases while the maximum stress
376 transmissible through the FRP plate slightly decreases with the width in agreement

377 with the CNR-rules [32].

378 5. CONCLUSIONS

379 A regularized XFEM approach with an experimentally consistent level set system has
380 been proposed for single-lap shear tests. It has been shown that: i) the bending of the FRP
381 plate plays a remarkable role on the debonding of the FRP plate; ii) an edge strengthening
382 effect due to the shear strain localization occurs along the edges of the FRP plate; iii) the
383 common design rules prescribing the variation of the nominal maximum stress with the
384 bonded width of the FRP plates have been confirmed; iv) a two-way delamination can be
385 observed.

386 As for the question whether a 3D or a 2D analysis should be preferred, the obtained
387 results have shown that, while the strain profiles along the FRP plate length, the peeling
388 and the different delamination onset locations can be detected indifferently through 3D and
389 2D analyses, the shear and axial strain on the debonded concrete surface display an edge
390 effect that can be captured only by means of a 3D analysis. Moreover, the dependence of
391 the maximum transferable load on the bonding width can be assessed only through a 3D
392 analysis.

393 Finally, the proposed regularized XFEM approach fits, and can be used as an alternative
394 to, experimental tests.

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	XFEM	Regularized XFEM
discontinuity	strong	weak
cohesive zone	traction-separation law	regularized stress-strain law
\mathbf{K}^e	\mathbf{K}_{st}^e (5)	\mathbf{K}_{reg}^e (6)
quadrature	sub-element quadrature	Gauss quadrature
internal length	no	yes

Table 1: Differences between the proposed approach and Belytschko’s et al. [14] original approach; matrix \mathbf{B}_e collects all the elemental contributions and \mathbf{K}^e is the elemental stiffness

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	Chajes et al.[26]			Carrara et al. [6]		
	FRP	glue	concrete	FRP	glue	concrete
E [MPa]	108380	1585	33640	168500	3517.3	28700
ν	0.248	0.315	0.2	0.248	0.315	0.2
f_t [MPa]	—	—	3.21	—	—	2.85
thickness [mm]	1.0	1.0	—	1.3	1.3	—
H [mm]	152.4	152.4	152.4	90	90	90
B [mm]	150	150	150	152.4	152.4	152.4
L [mm]	300	300	300	228.6	228.6	228.6
b [mm]	25.4	25.4	25.4	30	30	30

Table 2: Material and geometry parameters

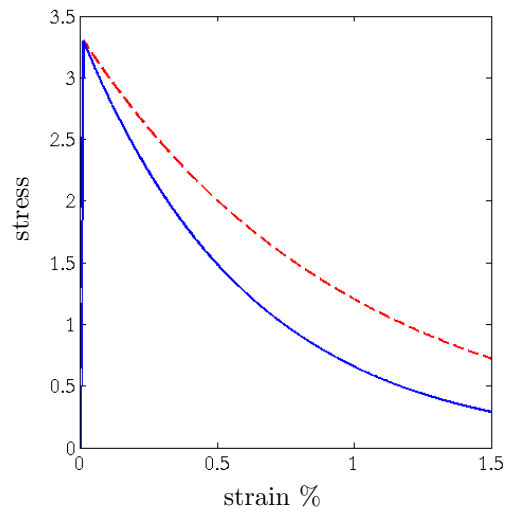


Figure 1:

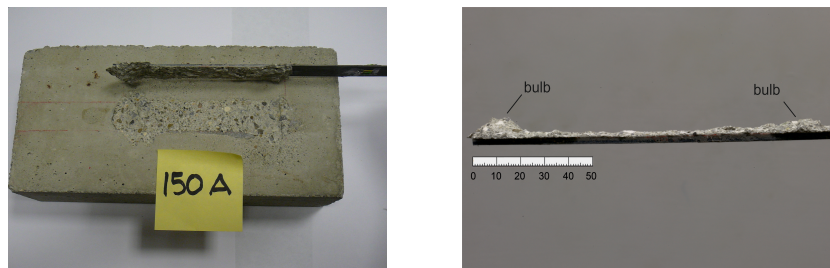


Figure 2:

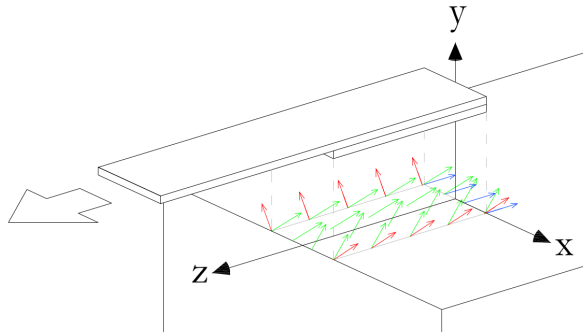


Figure 3:

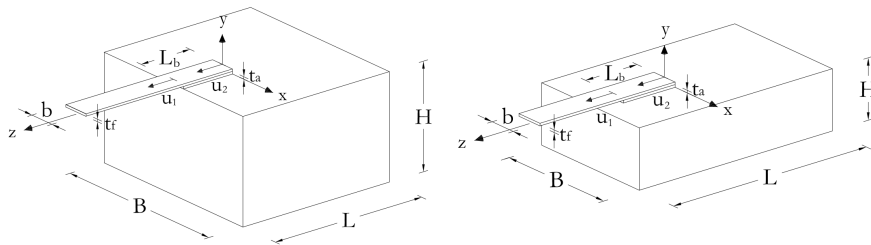


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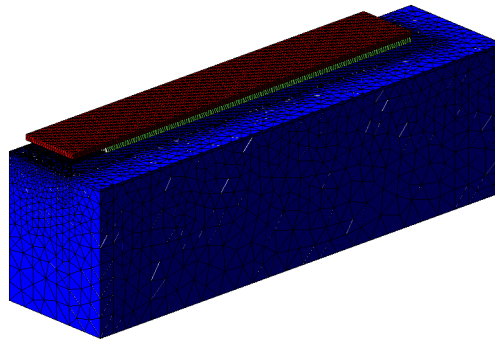


Figure 5:

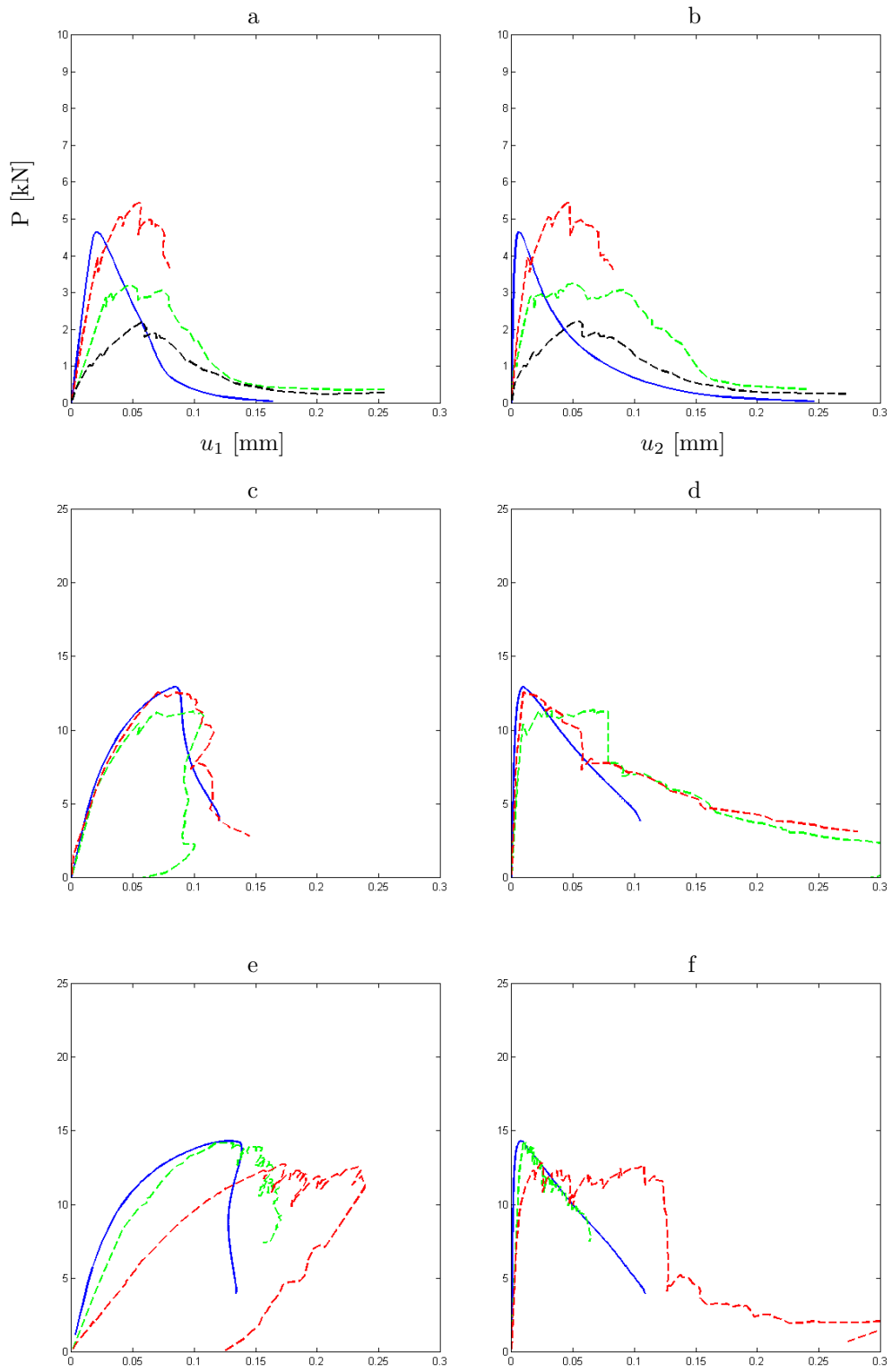


Figure 6:

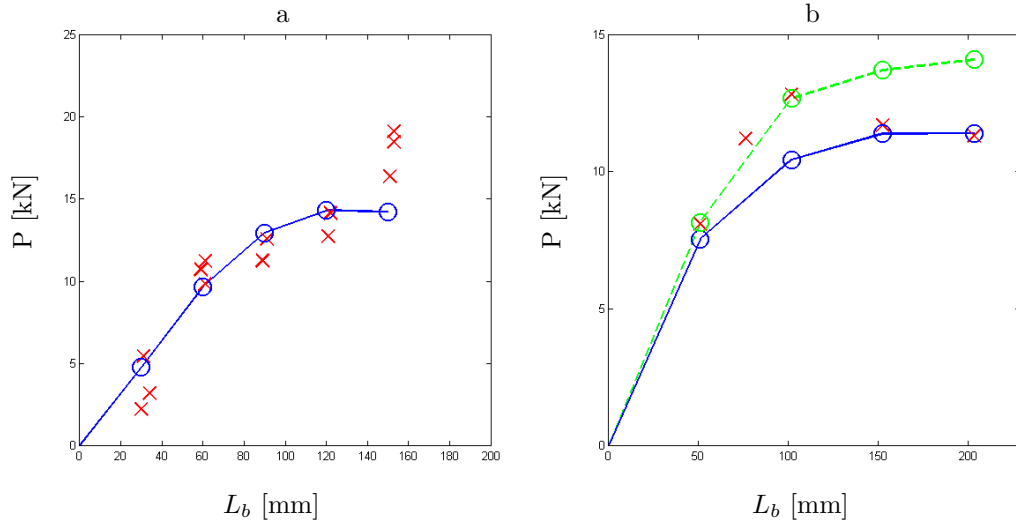


Figure 7:

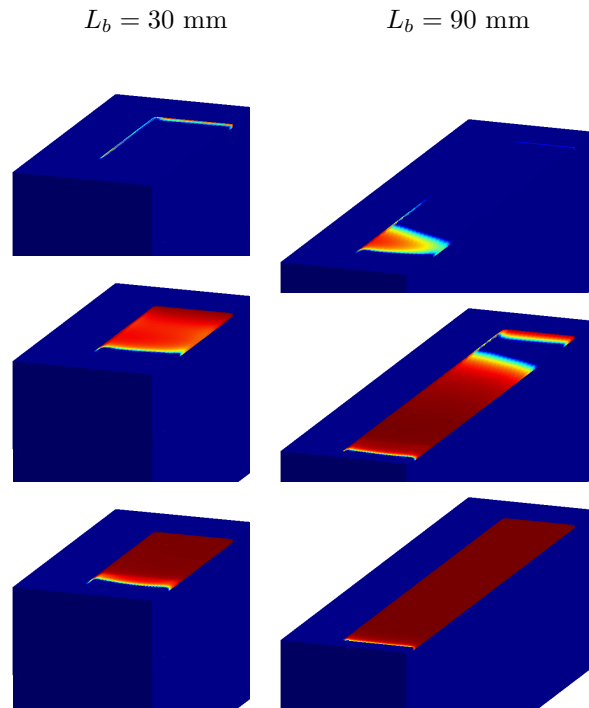


Figure 8:

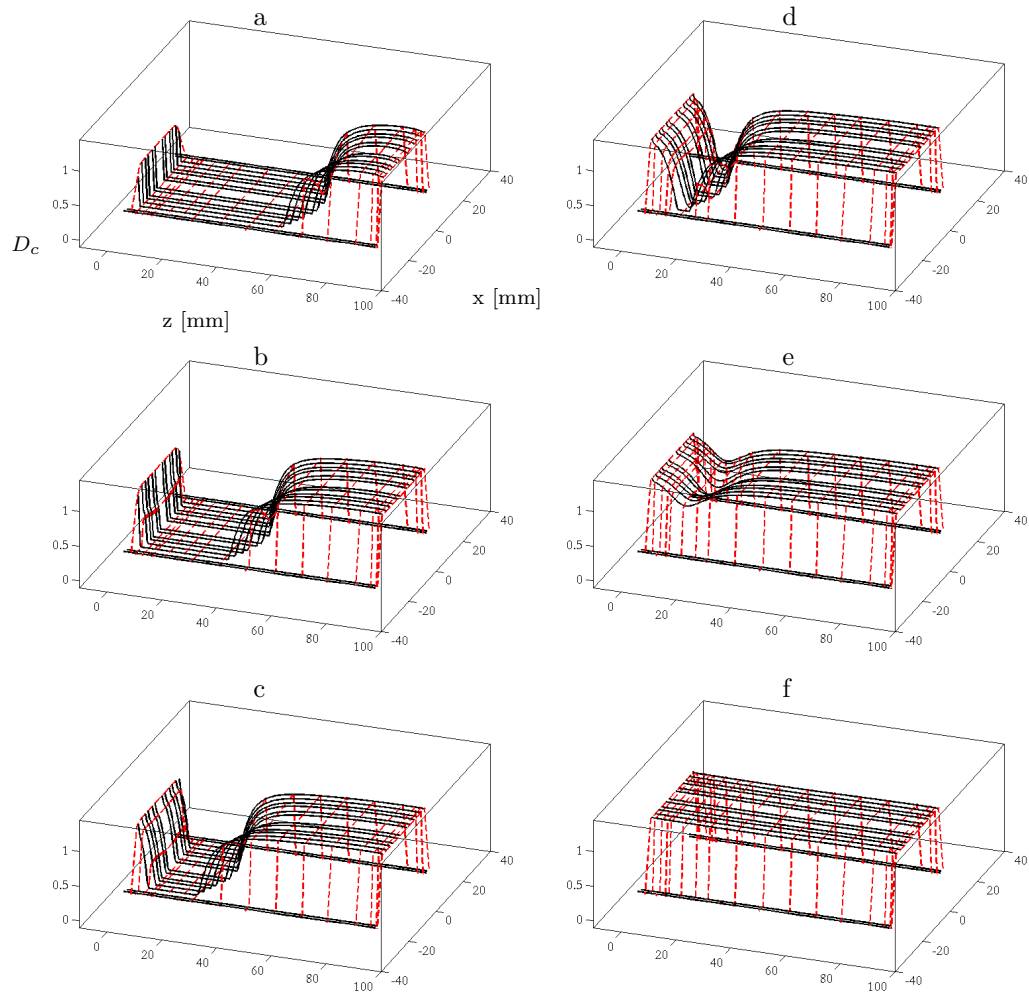


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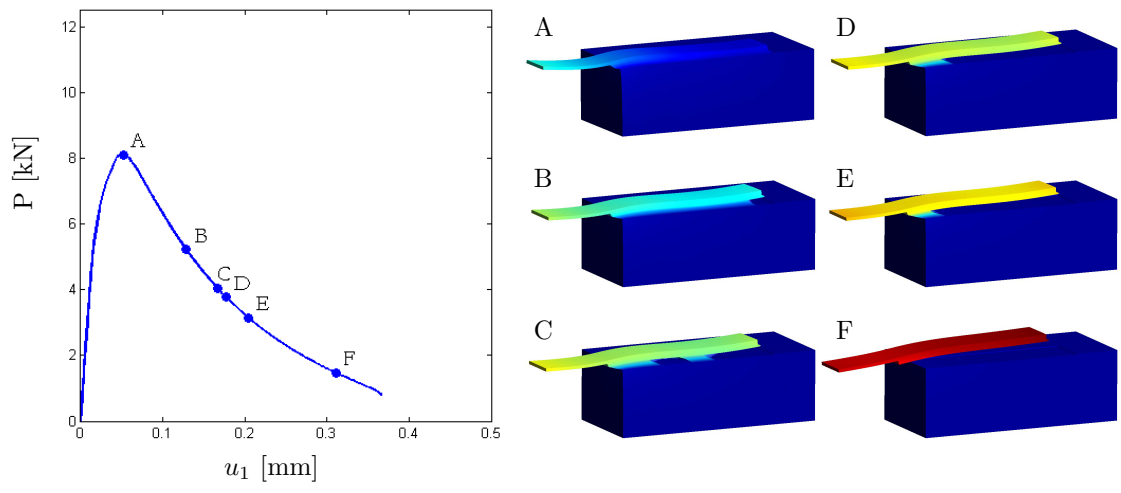


Figure 10:

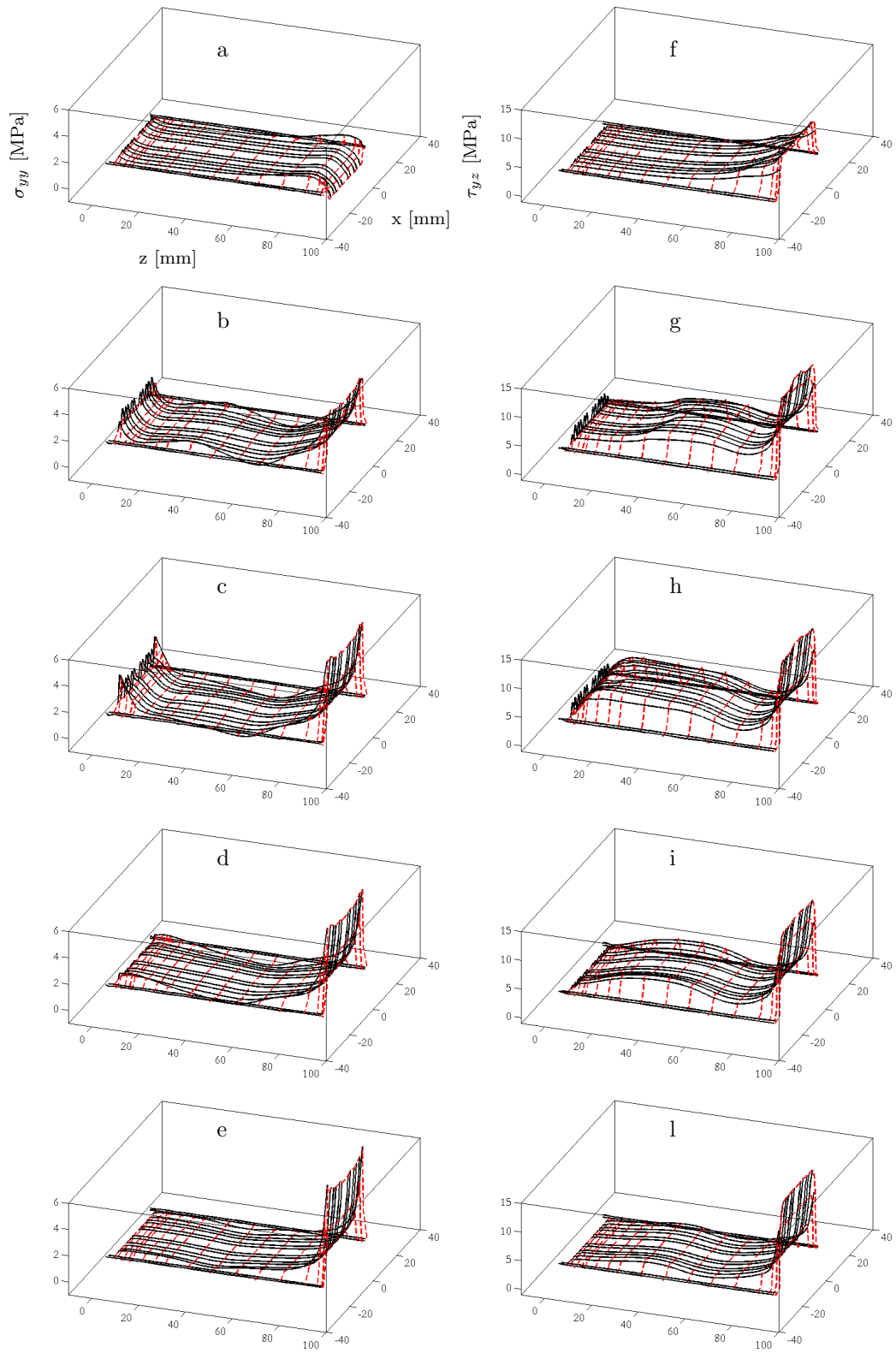


Figure 11:

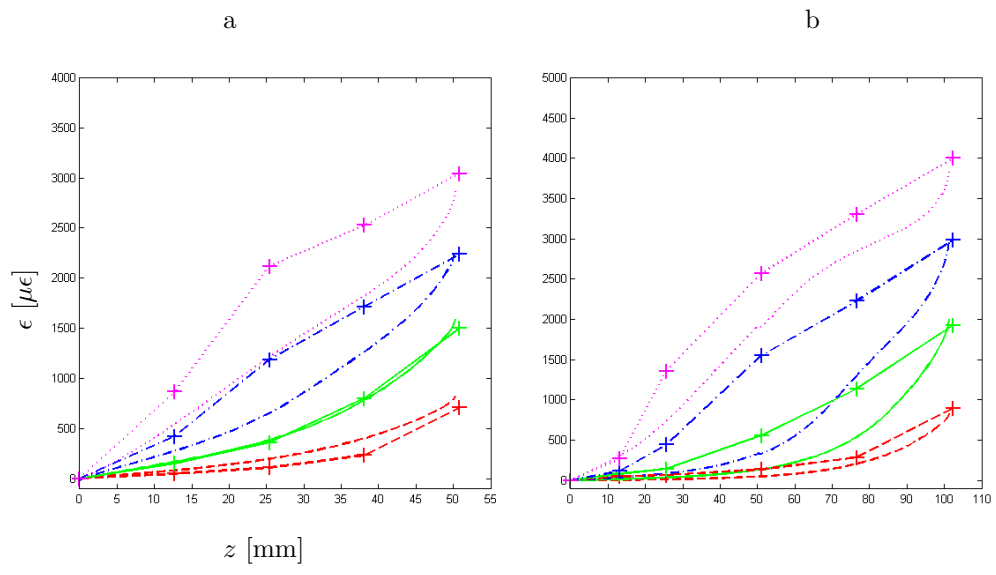


Figure 12:

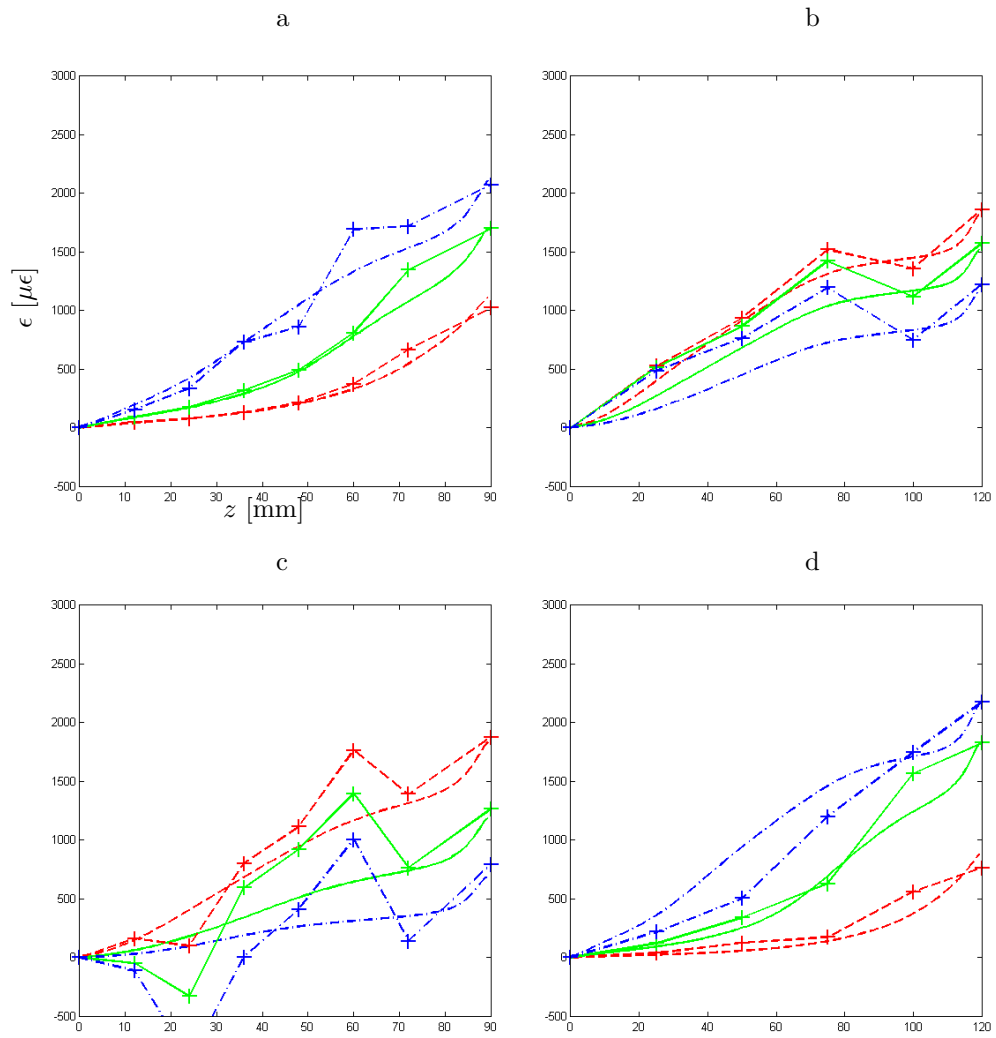


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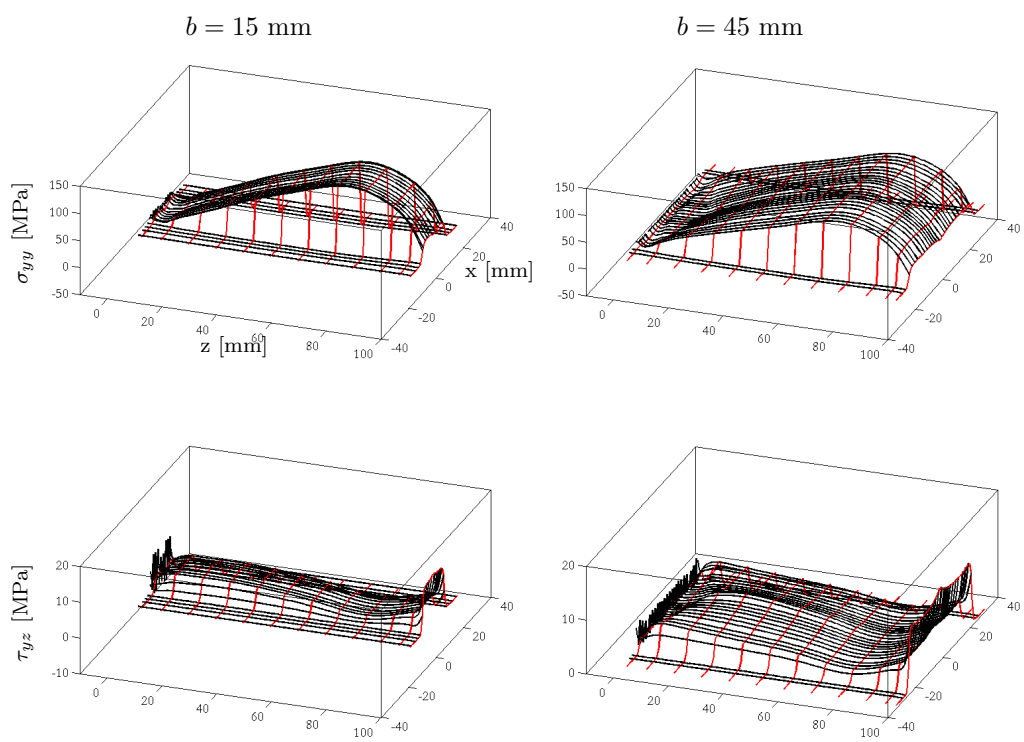


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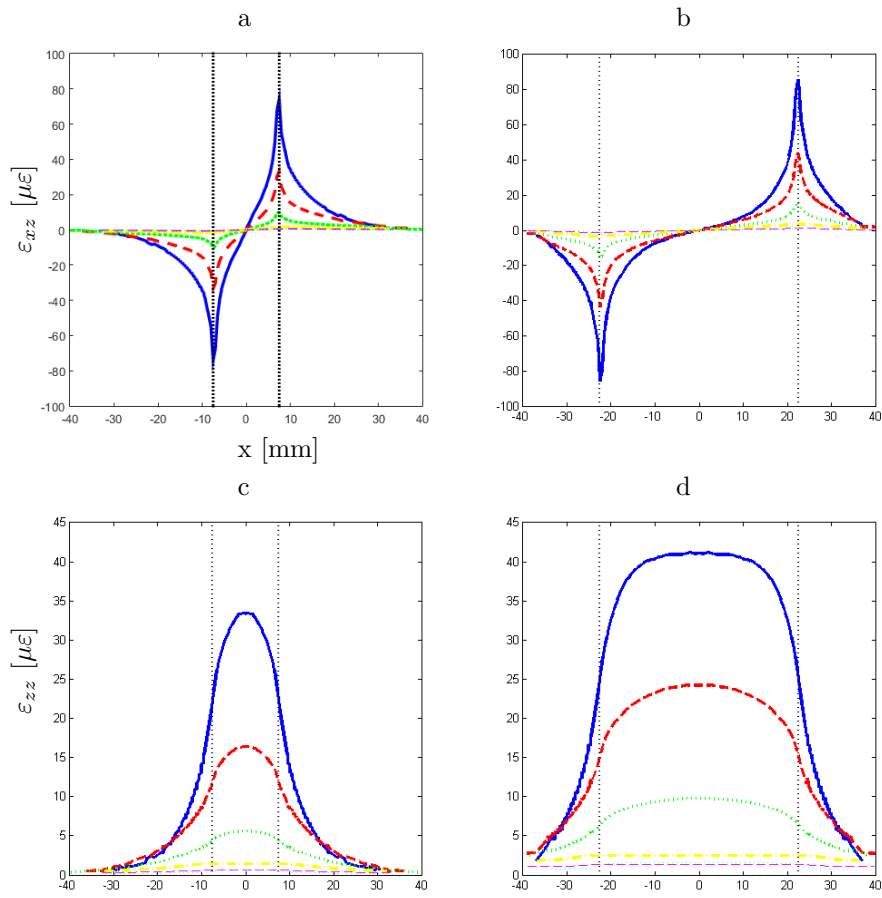


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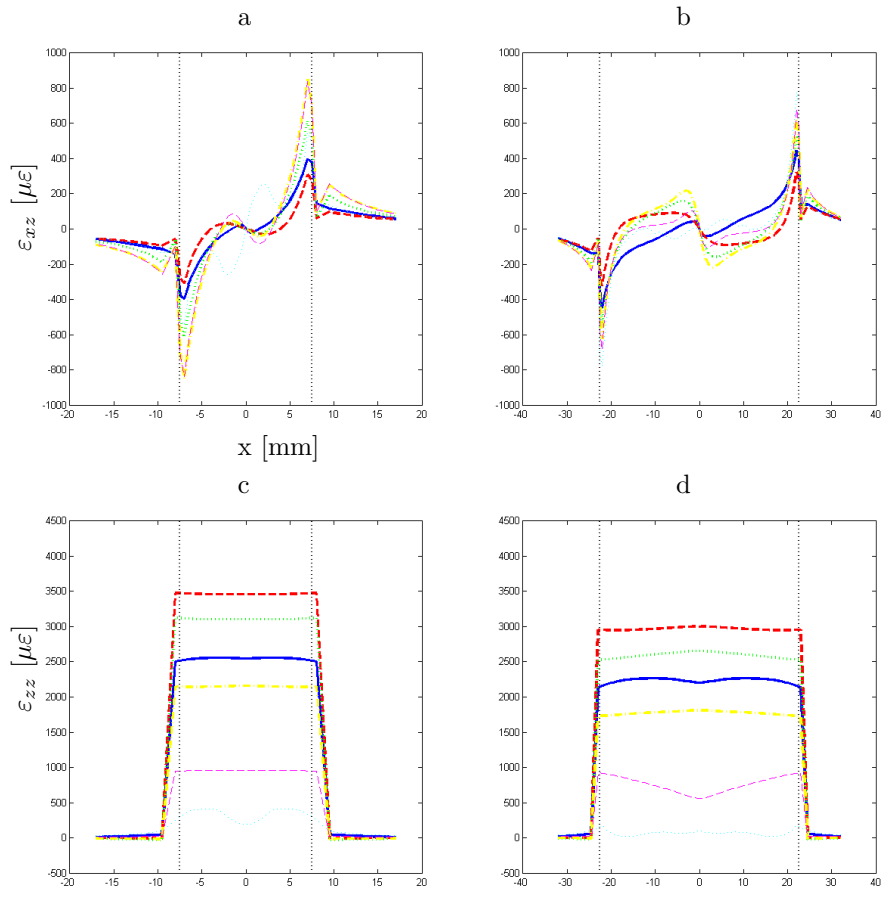


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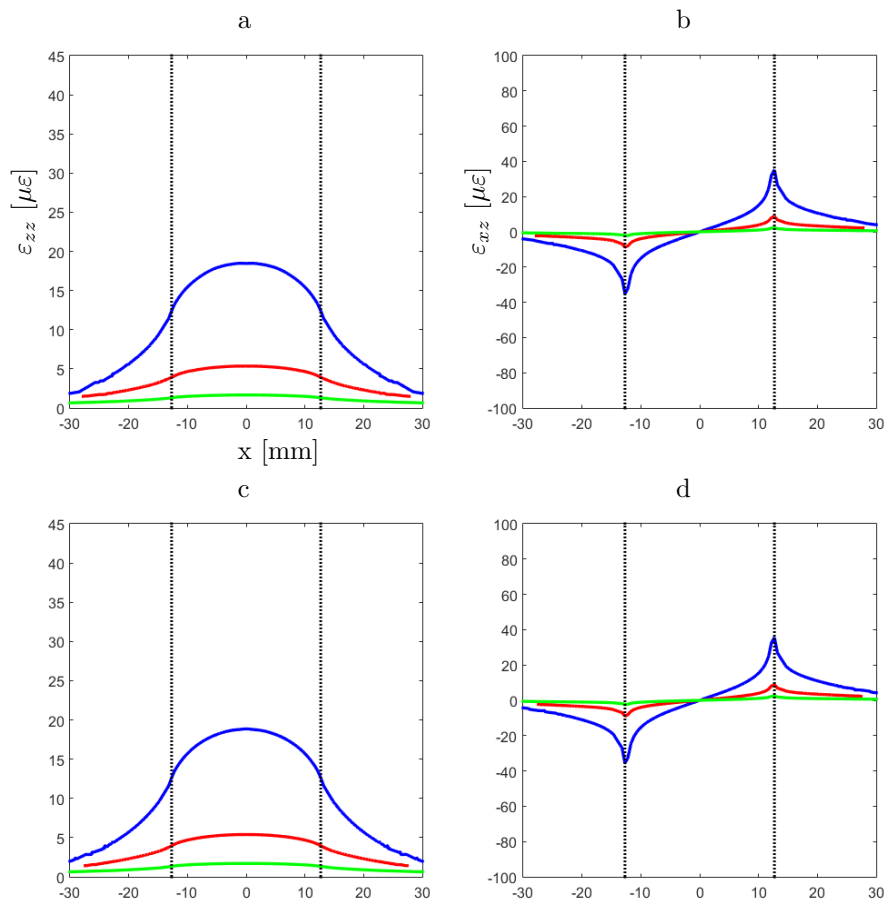


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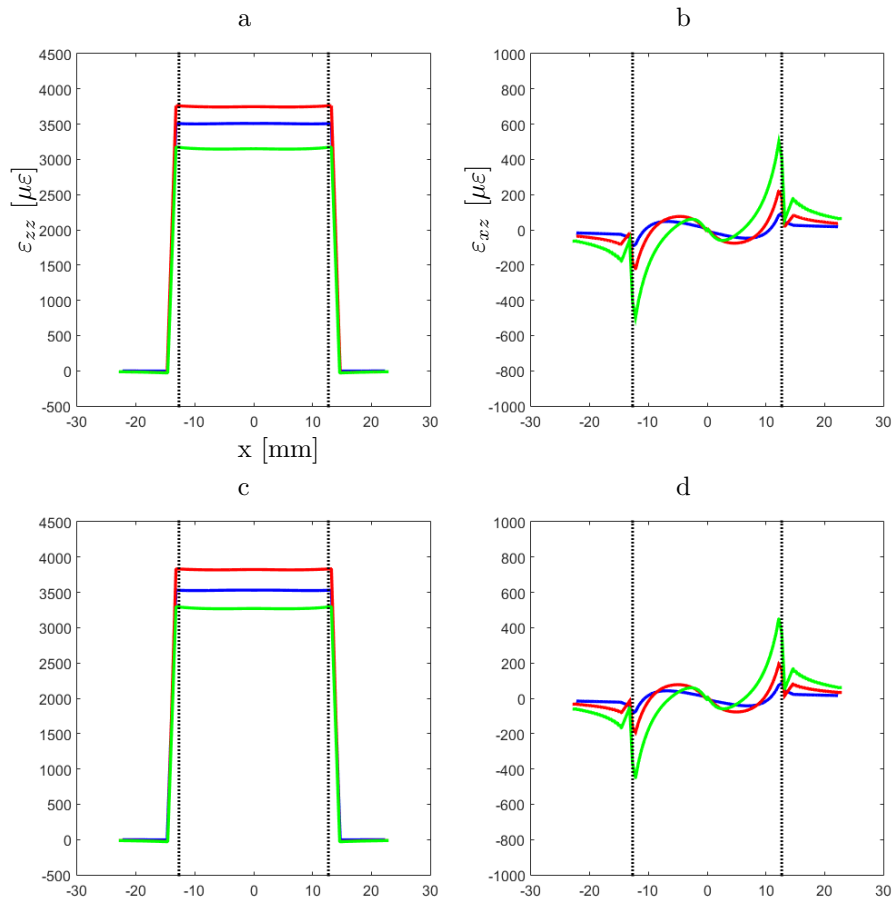


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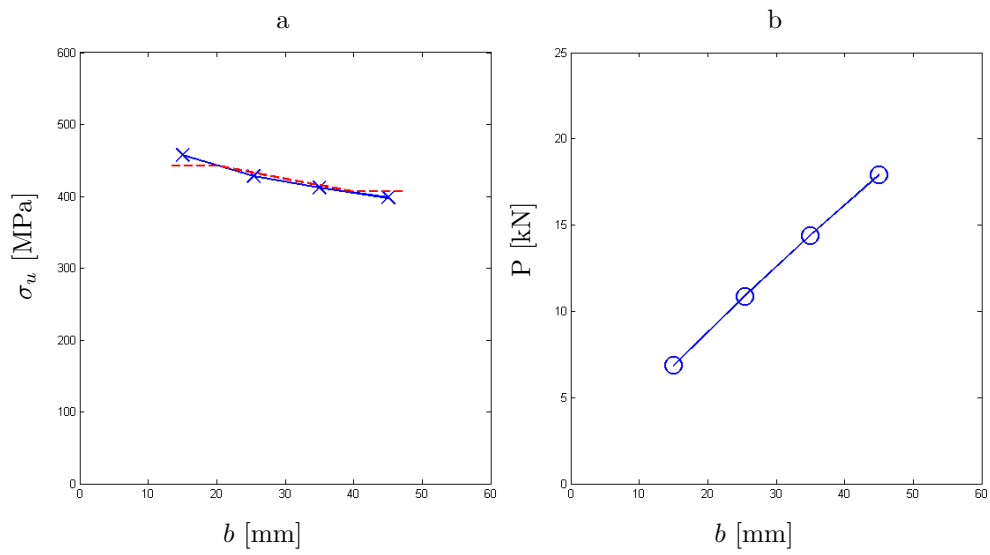


Figure 19:

Figures captions

- 480
- 481 Fig.1 Stress strain law corresponding to $H = 0.008$ (blue continuous line) and $H = 0.005$
482 (red dotted line) and $f_t = 3.21$ MPa
- 483 Fig.2 Photographs of the detached concrete layer after delamination for the tests of Carrara
484 et al. [6]
- 485 Fig.3 Qualitative picture of the vector \mathbf{n} associated with the *local* level set adopted in
486 Eq. (3b)
- 487 Fig.4 Geometry of the Chajes et al. [26] (a) and Carrara et al. [6] (b) specimens
- 488 Fig.5 Mesh of half of the Chajes et al. [26] specimen
- 489 Fig.6 Computed (continuous line) and experimental [6] (dashed lines) load-displacement
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491 loaded and the free end (see Fig. 4)
- 492 Fig.7 Computed (“o”) and experimental maximum loads (“x”) vs bonding length L_b [mm]
493 for the tests [6] a) and [26] b), where the dashed and the continuous lines refer to
494 $H = 0.005$ and $H = 0.008$, respectively
- 495 Fig.8 Damage evolution for the test [6] with $L_b = 30$ mm, on the left, and $L_b = 90$ mm, on
496 the right, evaluated (from top to bottom) at $P = 2.60$ kN, $P = 4.76$ kN, and $P = 3.75$
497 kN
- 498 Fig.9 Pre-peak (a,b,c), peak (d) and post-peak (e,f) damage profiles for $L_b = 90$ mm for
499 the test [6]
- 500 Fig.10 Computed load-displacement profile for $L_b = 50.8$ mm and corresponding deformed
501 mesh evolution at the loads A, B, C, D, E, F for the test [26]
- 502 Fig.11 Pre-peak (a), peak (b) and post-peak (c) profiles of σ_{yy} MPa (peeling) on the left and
503 τ_{yz} MPa (on the right) for the test [6] with $L_b = 90$ mm
- 504 Fig.12 1D equivalent axial strain along z for the test [26] obtained for $L_b = 50.8$ mm (a) and
505 $L_b = 101.6$ mm (b). The load levels are: $P = 1.94$ kN (red dashed line), $P = 4.06$ kN
506 (green continuous line), $P = 6.01$ kN (blue dash dotted line line), $P = 8.10$ kN (pink
507 dotted line) (a), and $P = 2.23$ kN (red dashed line), $P = 5.03$ kN (green continuous
508 line), $P = 7.71$ kN (blue dash dotted line line), $P = 10.29$ kN (pink dotted line) (b)
- 509 Fig.13 1D equivalent axial strain along z for the test [6] for $L_b = 90$ mm (c) and $L_b = 120$ mm
510 (d), at $P = 6.00$ kN (red dashed line), $P = 10.01$ kN (green continuous line), $P = 12.64$
511 kN (blue dash dotted line line) in (c), and $P = 5.01$ kN (red dashed line), $P = 12.00$ kN
512 (green continuous line), $P = 14.28$ kN (blue dash dotted line line) (d); $P = 11.15$ kN
513 (red dashed line), $P = 7.47$ kN (green continuous line), $P = 4.67$ kN (blue dash dotted
514 line line) in (e), and $P = 12.21$ kN (red dashed line), $P = 10.33$ kN (green continuous
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- 516 Fig.14 3D view of σ_{yy} and τ_{yz} for $b = 15$ mm and $b = 45$ mm with $L_b = 101.6$ mm evaluated
517 at the maximum load for the test [26]
- 518 Fig.15 Front view of ε_{xz} (a) and ε_{zz} (b) along x during the elastic stage for $b = 15$ mm (on
519 the left) and $b = 45$ mm (on the right) at $z = 0.5$ mm (cyan dotted line), $z = 18.1$ mm
520 (pink dashed line), $z = 38.1$ mm (yellow dash-dotted line), $z = 58.1$ mm (green dotted
521 line), $z = 78.1$ mm (red dashed line) and $z = 95.1$ mm (blu continuous line) for the
522 test [26]
- 523 Fig.16 Front view of ε_{xz} (a) and ε_{zz} (b) along x at the maximum load for $b = 15$ mm (on
524 the left) and $b = 45$ mm (on the right) in the test [26] (notation as in the previous
525 Figure)
- 526 Fig.17 Front view of ε_{zz} (a,c) and ε_{xz} (b,d) along x during the elastic stage for $L_b = 152.4$ mm
527 (on the left) and $L_b = 203.2$ mm (on the right) evaluated at a distance from the front
528 of 25.4 mm (blue line), 50.8 mm (red line), and 76.2 mm (green line) for the test [26]
- 529 Fig.18 Front view of ε_{xz} (a,c) and ε_{zz} (b,d) along x at the peak load for $L_b = 152.4$ mm (on
530 the left) and $L_b = 203.2$ mm (on the right) evaluated at a distance from the front of
531 the concrete block of 25.4 mm (blue line), 50.8 mm (red line), and 76.2 mm (green
532 line) for the test [26]
- 533 Fig.19 Computed nominal stress (a) and maximum load (b) vs bond width b and the CNR-
534 design-rule (7) with $\kappa_g = 0.065$ (dashed line) for the test [26]