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# Measurement of the ratio of branching fractions

$$\mathcal{B}(B_c^+ \rightarrow B_s^0 \pi^+)/\mathcal{B}(B_c^+ \rightarrow J/\psi \pi^+)$$



## The LHCb collaboration

E-mail: [qi.shi@cern.ch](mailto:qi.shi@cern.ch)

**ABSTRACT:** The ratio of branching fractions of  $B_c^+ \rightarrow B_s^0 \pi^+$  and  $B_c^+ \rightarrow J/\psi \pi^+$  decays is measured with proton-proton collision data of a centre-of-mass energy of 13 TeV. The data were collected with the LHCb experiment during 2016–2018, corresponding to an integrated luminosity of  $5.4 \text{ fb}^{-1}$ . The  $B_s^0$  mesons are reconstructed via the decays  $B_s^0 \rightarrow J/\psi \phi$  and  $B_s^0 \rightarrow D_s^- \pi^+$ . The ratio of branching fractions is measured to be  $\mathcal{B}(B_c^+ \rightarrow B_s^0 \pi^+)/\mathcal{B}(B_c^+ \rightarrow J/\psi \pi^+) = 91 \pm 10 \pm 8 \pm 3$  where the first uncertainty is statistical, the second is systematic and the third is due to the knowledge of the branching fractions of the intermediate state decays.

**KEYWORDS:** B Physics, Branching fraction, Flavour Physics, Hadron-Hadron Scattering

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## 1 Introduction

The  $B_c^{(*)+}$  family of mesons is the only one formed by two different heavy flavour quarks ( $\bar{b}c$ ). Both the  $\bar{b}$  quark and the  $c$  quark can each decay with the other as a spectator, leading to final states such as  $J/\psi\pi^+$  and  $B_s^0\pi^+$ . In addition, the  $\bar{b}$  quark and the  $c$  quark can annihilate via a  $W^+$  boson, allowing pure leptonic final states such as  $\tau^+\nu_\tau$ .

The  $B_c^+ \rightarrow B_s^0 \pi^+$  decay<sup>1</sup> was first observed in 2013 by the LHCb collaboration [1], and can be used to tag the initial flavour of the  $B_s^0$  meson with the charge of the accompanying  $\pi^+$  meson. The  $B_c^+ \rightarrow B_s^0 \pi^+$  decay is also the first observed case of one  $B$  meson decaying weakly into another  $B$  meson. This property makes the decay ideal for testing theoretical models. The branching fraction of the  $B_c^+ \rightarrow B_s^0 \pi^+$  decay should be large as it is a Cabibbo-favoured decay. There are several predictions for the branching fraction of the  $B_c^+ \rightarrow B_s^0 \pi^+$  decay based on QCD sum rules or quark-potential models, which range between 2.5% and 16.4% [2–9]. A precise measurement of the branching fraction of the  $B_c^+ \rightarrow B_s^0 \pi^+$  decay will improve the understanding of the  $B_c^+$  theory models.

The  $B_c^+ \rightarrow \tau^+\nu_\tau$  decay is highly sensitive to new physics effects [10–12] but experimental accessibility to this mode is limited. Due to the large branching fraction of the  $B_c^+ \rightarrow B_s^0 \pi^+$  decay, its improved measurement contributes to a more stringent limit on the  $B_c^+ \rightarrow \tau^+\nu_\tau$  decay via the  $B_c^+$  total decay width, depending on the theoretical model assumed [11, 13, 14].

In this paper, the ratio between the branching fractions of the  $B_c^+ \rightarrow B_s^0 \pi^+$  and  $B_c^+ \rightarrow J/\psi\pi^+$  decays is measured, using proton-proton ( $pp$ ) collision data collected with the LHCb experiment between 2016 and 2018 at a centre-of-mass energy of 13 TeV, corresponding

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<sup>1</sup>Inclusion of charge conjugate processes is implied throughout this paper.

| Decay                             | Branching fraction [%] |
|-----------------------------------|------------------------|
| $D_s^- \rightarrow K^+ K^- \pi^-$ | $5.38 \pm 0.10$        |
| $B_s^0 \rightarrow D_s^- \pi^+$   | $0.298 \pm 0.014$      |
| $B_s^0 \rightarrow J/\psi \phi$   | $0.104 \pm 0.004$      |
| $\phi \rightarrow K^+ K^-$        | $49.1 \pm 0.5$         |
| $J/\psi \rightarrow \mu^+ \mu^-$  | $5.961 \pm 0.033$      |

**Table 1.** Branching fractions of intermediate-state decays [15].

to an integrated luminosity of  $5.4 \text{ fb}^{-1}$ . The ratio of branching fractions is measured separately using the  $B_s^0 \rightarrow J/\psi(\rightarrow \mu^+ \mu^-) \phi(\rightarrow K^+ K^-)$  and  $B_s^0 \rightarrow D_s^- (\rightarrow K^+ K^- \pi^-) \pi^+$  decays. These two measurements are then combined. The ratio of branching fractions is evaluated as

$$\mathcal{R}_X \equiv \frac{\mathcal{B}(B_c^+ \rightarrow B_s^0 \pi^+)}{\mathcal{B}(B_c^+ \rightarrow J/\psi \pi^+)} = \frac{N_{B_c^+ \rightarrow B_s^0 \pi^+}}{N_{B_c^+ \rightarrow J/\psi \pi^+}} \cdot \frac{\epsilon_{B_c^+ \rightarrow J/\psi \pi^+}}{\epsilon_{B_c^+ \rightarrow B_s^0 \pi^+}} \cdot \frac{1}{\mathcal{R}_{\text{int}}^X}, \quad (1.1)$$

where  $N$  is the signal yield,  $\epsilon$  is the total efficiency to reconstruct these decays, and  $\mathcal{R}_{\text{int}}^X$  is the ratio of branching fractions of the corresponding intermediate state decays, as shown in table 1. The  $X$  represents the final state of the  $B_s^0$  decay,  $J/\psi \phi$  or  $D_s^- \pi^+$ . The superscripts or subscripts  $J/\psi \phi$  and  $D_s^- \pi^+$  are used to indicate the  $B_c^+ \rightarrow B_s^0 (\rightarrow J/\psi \phi) \pi^+$  and  $B_c^+ \rightarrow B_s^0 (\rightarrow D_s^- \pi^+) \pi^+$  decay modes. The  $\mathcal{R}_{\text{int}}^X$  are defined as

$$\mathcal{R}_{\text{int}}^{J/\psi \phi} \equiv \mathcal{B}(B_s^0 \rightarrow J/\psi \phi) \cdot \mathcal{B}(\phi \rightarrow K^+ K^-), \quad (1.2)$$

$$\mathcal{R}_{\text{int}}^{D_s^- \pi^+} \equiv \mathcal{B}(B_s^0 \rightarrow D_s^- \pi^+) \cdot \mathcal{B}(D_s^- \rightarrow K^+ K^- \pi^-) / \mathcal{B}(J/\psi \rightarrow \mu^+ \mu^-) \quad (1.3)$$

with corresponding values of  $\mathcal{R}_{\text{int}}^{J/\psi \phi} = (5.11 \pm 0.20) \times 10^{-4}$  and  $\mathcal{R}_{\text{int}}^{D_s^- \pi^+} = (2.69 \pm 0.14) \times 10^{-3}$ .

## 2 Detector and simulation

The LHCb detector [16, 17] is a single-arm forward spectrometer covering the pseudorapidity range  $2 < \eta < 5$ , designed for the study of particles containing  $b$  or  $c$  quarks. The detector includes a high-precision tracking system consisting of a silicon-strip vertex detector surrounding the  $pp$  interaction region [18], a large-area silicon-strip detector located upstream of a dipole magnet with a bending power of about 4 Tm, and three stations of silicon-strip detectors and straw drift tubes placed downstream of the magnet [19]. The tracking system provides a measurement of the momentum,  $p$ , of charged particles with a relative uncertainty that varies from 0.5% at low momentum to 1.0% at  $200 \text{ GeV}/c$ . The minimum distance of a track to a primary  $pp$  collision vertex (PV), the impact parameter (IP), is measured with a resolution of  $(15 + 29/p_T) \mu\text{m}$ , where  $p_T$  is the component of the momentum transverse to the beam, in  $\text{GeV}/c$ . Different types of charged hadrons are distinguished using information from two ring-imaging Cherenkov detectors [20]. Photons, electrons and hadrons are identified by a calorimeter system consisting of scintillating-pad and preshower

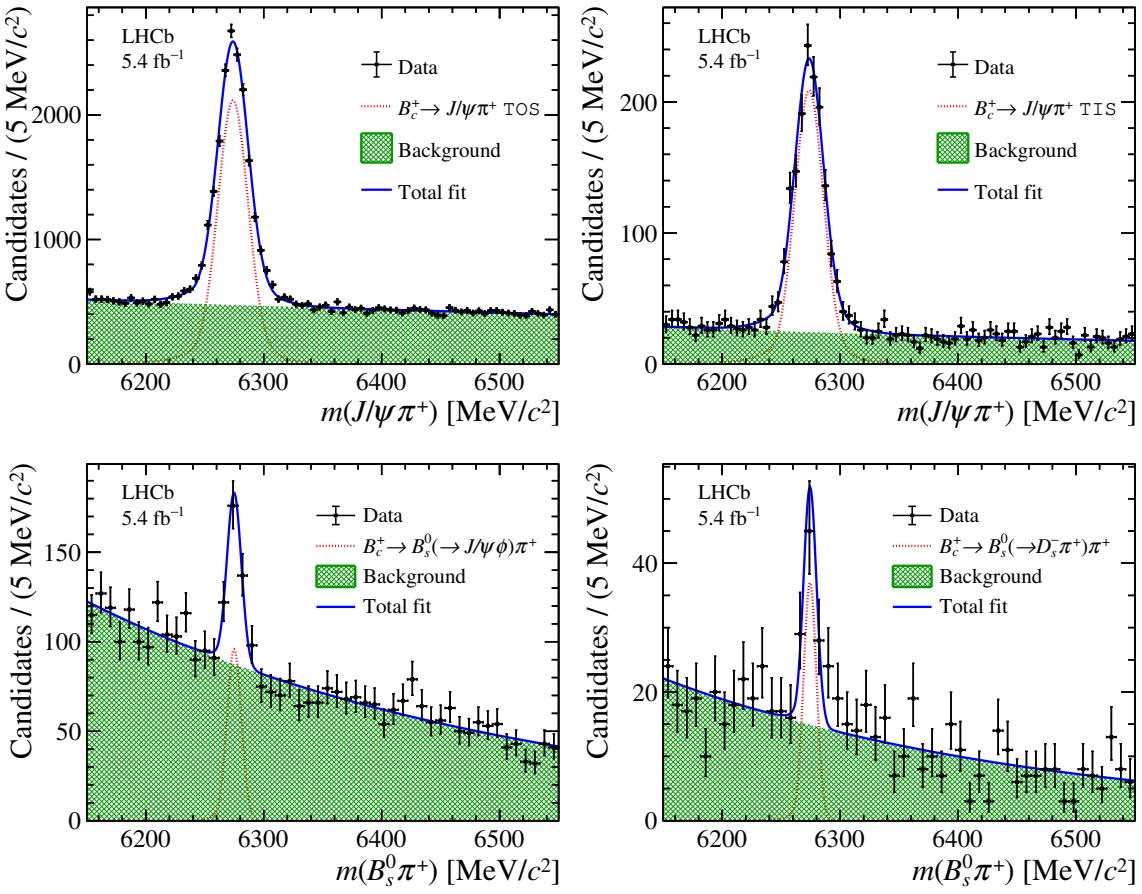
detectors, an electromagnetic and a hadronic calorimeter. Muons are identified by a system composed of alternating layers of iron and multiwire proportional chambers [21]. The online event selection is performed by a trigger [22], which consists of a hardware stage, based on information from the calorimeter and muon systems, followed by a software stage, which applies a full event reconstruction.

Simulation is required to model the effects of the detector acceptance and the imposed selection requirements. In the simulation,  $pp$  collisions are generated using PYTHIA 8 [23, 24] with a specific LHCb configuration [25]. A dedicated generator BCVEGPY [26] is used to simulate the production of  $B_c^+$  mesons. Decays of unstable particles are described by EVTGEN [27], in which final-state radiation is generated using PHOTOS [28]. The interaction of the generated particles with the detector, and its response, are implemented using the GEANT4 toolkit [29, 30] as described in ref. [31].

### 3 Event selection

In the  $B_c^+ \rightarrow B_s^0\pi^+$  decay mode, the  $B_s^0$  meson is reconstructed using two decay modes,  $B_s^0 \rightarrow J/\psi\phi$  and  $B_s^0 \rightarrow D_s^-\pi^+$ , with the intermediate states selected based on the  $D_s^- \rightarrow K^+K^-\pi^-$ ,  $J/\psi \rightarrow \mu^+\mu^-$  and  $\phi \rightarrow K^+K^-$  decays. The reconstructed  $B_s^0$  candidates are selected using the boosted decision tree (BDT) [32–34] classifier used in refs. [1, 35], with the same working points. In the  $B_c^+ \rightarrow J/\psi\pi^+$  decay channel, the  $J/\psi$  mesons are reconstructed using a pair of oppositely charged muons. The selected  $B_s^0$  and  $J/\psi$  candidates are combined with a track identified as a pion to reconstruct the  $B_c^+$  candidates. These candidates and all intermediate states are required to have good vertex-fit quality. All final state particles (pions, kaons and muons) are required to have a good track-fit quality and high transverse momentum. Particle identification (PID) information is used to suppress misidentified tracks. The reconstructed masses of intermediate state particles are required to be within three times the expected mass resolution of their known masses [15]. A second boosted decision tree classifier taken from refs. [35, 36] is used to further distinguish signal  $B_c^+$  mesons from combinatorial background. The vertex fit quality, transverse momentum and topological information are used in all BDT classifiers.

In the offline selection, trigger signatures are associated with reconstructed particles. Selection requirements can therefore be made on the trigger selection itself and on whether the event was selected at trigger level because of the signal candidate (denoted TOS), or because of the decay products of the  $b$ - or  $c$ -hadrons produced together with the signal  $B_c^+$  meson (denoted TIS). For the  $B_c^+ \rightarrow B_s^0(\rightarrow J/\psi\phi)\pi^+$  decay, the reconstructed  $J/\psi$  mesons are required to pass the hardware trigger that selects a muon or a pair of muons with high  $p_T$  (the muon TOS trigger). For the  $B_c^+ \rightarrow B_s^0(\rightarrow D_s^-\pi^+)\pi^+$  decay, hadronic triggers are avoided due to the complexity of simulating the detector response of many-body final states overlapping in the calorimeter. Instead, events containing  $B_c^+ \rightarrow B_s^0\pi^+$  decays are required to pass the trigger in case there is a muon or pair of muons in the rest of the event, not related to the  $B_c^+$  daughter particles (the muon TIS trigger). This requirement gives better control of trigger efficiencies, even though it reduces the signal yields. In the  $\mathcal{R}_{J/\psi\phi}$  case, the TOS trigger requirement is applied to the  $B_c^+ \rightarrow B_s^0\pi^+$  and  $B_c^+ \rightarrow J/\psi\pi^+$  decays, and in the  $\mathcal{R}_{D_s^-\pi^+}$  case, the TIS trigger requirement is applied.



**Figure 1.** Invariant-mass distributions of  $B_c^+$  candidates: (top left)  $B_c^+ \rightarrow J/\psi\pi^+$  TOS, (top right)  $B_c^+ \rightarrow J/\psi\pi^+$  TIS, (bottom left)  $B_c^+ \rightarrow B_s^0(\rightarrow J/\psi\phi)\pi^+$ , and (bottom right)  $B_c^+ \rightarrow B_s^0(\rightarrow D_s^-\pi^+)\pi^+$ . The black dots represent the data, the blue line represents the results of the total fit, and the red dashed line and green shaded area show the signal and combinatorial background components, respectively.

#### 4 Signal yields and efficiencies

The signal yields are determined from unbinned maximum-likelihood fits to the  $B_c^+$  candidate invariant-mass distributions. The invariant mass spectrum of  $B_c^+$ -meson candidates is obtained from a fit to the whole decay chain with the masses of  $B_s^0$ ,  $J/\psi$ , and  $D_s^-$  mesons constrained to their known values [15], as described in ref. [37]. For the  $B_c^+ \rightarrow J/\psi\pi^+$  decay, the signal component is modelled using the sum of two double-sided Crystal Ball (DSCB) functions [38] with a common peak position. For the  $B_c^+ \rightarrow B_s^0\pi^+$  decay, one DSCB function is sufficient to model the signal shape. The tail parameters of all DSCB functions are determined from simulation, as detailed in ref. [35]. In both cases the combinatorial background is modelled with an exponential function. The fit results are shown in figure 1, and the signal yields are summarised in table 2.

The total efficiency is calculated as the product of the geometrical detector acceptance and of the efficiencies related to particle reconstruction, event selection, particle identification

| Decay channel  | Signal yields   |
|--|-----------------|
| $B_c^+ \rightarrow J/\psi\pi^+$ (TOS)                  | $14641 \pm 195$ |
| $B_c^+ \rightarrow J/\psi\pi^+$ (TIS)                  | $1376 \pm 45$   |
| $B_c^+ \rightarrow B_s^0(\rightarrow J/\psi\phi)\pi^+$ | $186 \pm 25$    |
| $B_c^+ \rightarrow B_s^0(\rightarrow D_s^-\pi^+)\pi^+$ | $65 \pm 12$     |

**Table 2.** Yields of the normalisation and signal decay modes.

| Decay channel  | Efficiency [%]        |
|--|-----------------------|
| $B_c^+ \rightarrow J/\psi\pi^+$ (TOS)                  | $1.6141 \pm 0.0028$   |
| $B_c^+ \rightarrow J/\psi\pi^+$ (TIS)                  | $0.15846 \pm 0.00063$ |
| $B_c^+ \rightarrow B_s^0(\rightarrow J/\psi\phi)\pi^+$ | $0.4526 \pm 0.0013$   |
| $B_c^+ \rightarrow B_s^0(\rightarrow D_s^-\pi^+)\pi^+$ | $0.02801 \pm 0.00028$ |

**Table 3.** Total efficiencies of the  $B_c^+ \rightarrow J/\psi\pi^+$  and  $B_c^+ \rightarrow B_s^0\pi^+$  decay modes. The uncertainties are due to the finite size of the simulated samples.

and trigger decision. The efficiencies are determined using simulation calibrated with data. Since the number of final-state particles is different between the signal and normalisation modes (five and three, respectively), accurate modelling of the track reconstruction efficiency is important. For every event, a track-by-track correction to the track reconstruction efficiency of all final states is applied using a calibration sample of  $J/\psi \rightarrow \mu^+\mu^-$  decays [39]. The PID information used in the event selection is corrected with high-yield calibration samples ( $D^{*+} \rightarrow D^0\pi^+$  with  $D^0 \rightarrow K^-\pi^+$  decays for hadrons, and the  $J/\psi \rightarrow \mu^+\mu^-$  decay for muons) [40]. The total efficiencies are listed in table 3.

## 5 Systematic uncertainties

The measurement is affected by the systematic uncertainty in the determination of signal yields and efficiencies, as summarised in table 4.

The systematic uncertainty due to the signal lineshape modelling is studied using pseudoexperiments. The signal candidates in the fully simulated samples are mixed with background events that are randomly generated with the exponential shape and fraction determined from the fits to the data. The same fit model used for the data fit is applied to these samples. The difference between the fitted value of the signal yield ratio  $N_{B_c^+ \rightarrow B_s^0\pi^+}/N_{B_c^+ \rightarrow J/\psi\pi^+}$  and the generated ratio in simulated samples is taken as the systematic uncertainty, 3.0% for  $\mathcal{R}_{J/\psi\phi}$  and 4.4% for  $\mathcal{R}_{D_s^-\pi^+}$ . The systematic uncertainty of the background lineshape is estimated by using a first-order polynomial function as an alternative shape, and the difference with the default result is taken as the systematic uncertainty, which is 5.5% for  $\mathcal{R}_{J/\psi\phi}$  and 4.5% for  $\mathcal{R}_{D_s^-\pi^+}$ .

|   | $\mathcal{R}_{J/\psi\phi}$ [%] | $\mathcal{R}_{D_s^-\pi^+}$ [%] |
|---|--------------------------------|--------------------------------|
| Signal model                                | 3.0                            | 4.4                            |
| Background model                            | 5.5                            | 4.5                            |
| $B_c^+ \rightarrow B_s^{*0}\pi^+$ component | 3.2                            | 3.1                            |
| Data-simulation agreement                   | 2.3                            | 4.8                            |
| Tracking efficiency                         | 3.2                            | 6.0                            |
| Simulation sample size                      | 0.3                            | 1.1                            |
| Trigger                                     | 1.0                            | 1.0                            |
| Total                                       | 8.1                            | 10.5                           |

**Table 4.** Relative systematic uncertainties in the ratio of the branching fractions where the symbol  $\mathcal{R}_{J/\psi\phi}$  indicates the ratio of branching fractions measured with the  $B_c^+ \rightarrow B_s^0(\rightarrow J/\psi\phi)\pi^+$  decay channel and  $\mathcal{R}_{D_s^-\pi^+}$  that using the  $B_c^+ \rightarrow B_s^0(\rightarrow D_s^-\pi^+)\pi^+$  decay.

The  $B_c^+$  mesons can also decay to a  $B_s^{*0}\pi^+$  state, followed by the  $B_s^{*0} \rightarrow B_s^0\gamma$  decay. This final state, with an unobserved photon, mimics the signal but with the reconstructed  $m(B_s^0\pi^+)$  shifted to lower values by the mass difference between  $B_s^{*0}$  and  $B_s^0$  mesons to a good approximation. In this case, the  $B_c^+$  mass resolution is almost unaffected [1, 41]. To estimate the impact of this potential partially reconstructed decay, an alternative fit to the  $m(B_s^0\pi^+)$  mass distribution is performed including an additional  $B_s^{*0}\pi^+$  component. The fit model for  $B_c^+ \rightarrow B_s^{*0}\pi^+$  is a DSCB function with the same tail parameters as in the default fit. The mass resolution is assumed to be the same as the signal and the mass shift is fixed to the known mass difference between  $B_s^0$  and  $B_s^{*0}$  mesons,  $48.5_{-1.5}^{+1.8}$  MeV/ $c^2$  [15]. The difference in the signal yield ratio is taken as the relevant systematic uncertainty, which is 3.2% for  $\mathcal{R}_{J/\psi\phi}$  and 3.1% for  $\mathcal{R}_{D_s^-\pi^+}$ .

As some efficiencies are determined from simulated samples, any discrepancy between data and simulated events can introduce a bias. Distributions of kinematic variables in simulation and in background-subtracted data using the *sPlot* technique [42] are compared. The differences of the efficiencies between the unchanged simulation and after the alignment between data and simulation are 2.3% for  $\mathcal{R}_{J/\psi\phi}$  and 4.8% for  $\mathcal{R}_{D_s^-\pi^+}$ , respectively, which are taken as systematic uncertainties.

The tracking efficiency, which is determined on data calibration samples, has a systematic uncertainty of 1% per track. Due to the hadronic interactions and the uncertainty of the LHCb material-budget in simulation, there is 1.1% additional uncertainty for kaons and 1.4% additional uncertainty for pions. The uncertainties for the same particles in numerator and denominator largely cancel in the ratio. Therefore, the systematic uncertainty is found to be 3.2% for the  $\mathcal{R}_{J/\psi\phi}$  measurement and 6.0% for  $\mathcal{R}_{D_s^-\pi^+}$ .

The efficiencies determined from simulated events have uncertainties due to the limited size of the samples. The relative uncertainty in the total efficiency ratio  $\epsilon(B_c^+ \rightarrow J/\psi\pi^+)/\epsilon(B_c^+ \rightarrow B_s^0\pi^+)$  is considered as the systematic uncertainty, the value is 0.3% for  $\mathcal{R}_{J/\psi\phi}$  and 1.1% for  $\mathcal{R}_{D_s^-\pi^+}$ . The imperfect simulation of the trigger system can also cause a bias in the trigger selection efficiency. An uncertainty of 1.0% is assigned to the trigger

selection following ref. [43]. Other potential sources of systematic uncertainty, such as those associated with the PID calibration, are found to be negligible. A summary of the systematic uncertainties is given in table 4.

## 6 Results

The ratio of the total branching fractions, including the branching fractions of the intermediate states, is determined as

$$\mathcal{R}_{J/\psi\phi}^{\text{tot}} \equiv \frac{N_{B_c^+ \rightarrow B_s^0(\rightarrow J/\psi\phi)\pi^+}}{N_{B_c^+ \rightarrow J/\psi\pi^+}} \cdot \frac{\epsilon_{B_c^+ \rightarrow J/\psi\pi^+}}{\epsilon_{B_c^+ \rightarrow B_s^0(\rightarrow J/\psi\phi)\pi^+}}, \quad (6.1)$$

$$\mathcal{R}_{D_s^-\pi^+}^{\text{tot}} \equiv \frac{N_{B_c^+ \rightarrow B_s^0(\rightarrow D_s^-\pi^+)\pi^+}}{N_{B_c^+ \rightarrow J/\psi\pi^+}} \cdot \frac{\epsilon_{B_c^+ \rightarrow J/\psi\pi^+}}{\epsilon_{B_c^+ \rightarrow B_s^0(\rightarrow D_s^-\pi^+)\pi^+}}, \quad (6.2)$$

and are measured to be

$$\begin{aligned} \mathcal{R}_{J/\psi\phi}^{\text{tot}} &= 0.0453 \pm 0.0061 \text{ (stat)} \pm 0.0037 \text{ (syst)}, \\ \mathcal{R}_{D_s^-\pi^+}^{\text{tot}} &= 0.267 \pm 0.050 \text{ (stat)} \pm 0.028 \text{ (syst)}, \end{aligned}$$

in each decay mode. Accounting for the branching fractions of the intermediate state decays, as shown in table 1, the ratio between the branching fractions of the  $B_c^+ \rightarrow B_s^0\pi^+$  and  $B_c^+ \rightarrow J/\psi\pi^+$  decays is

$$\begin{aligned} \mathcal{R}_{J/\psi\phi} &= 89 \pm 12 \text{ (stat)} \pm 7 \text{ (syst)} \pm 4 \text{ (\mathcal{B})}, \\ \mathcal{R}_{D_s^-\pi^+} &= 99 \pm 19 \text{ (stat)} \pm 10 \text{ (syst)} \pm 5 \text{ (\mathcal{B})}, \end{aligned}$$

for the two decay modes considered. The third uncertainty is due to the knowledge of the branching fractions of intermediate state decays (denoted  $\mathcal{B}$  hereafter).

The ratio of the branching fractions measured using the two decay modes can then be combined using the BLUE method [44]. The systematic uncertainties due to the signal lineshape, background lineshape,  $B_c^+ \rightarrow B_s^{*0}\pi^+$  component, data-simulation agreement, tracking efficiency and trigger decision are considered to be 100% correlated between the two decay modes, while the uncertainty due to limited simulation sample size is uncorrelated. Uncertainties in this ratio arise from knowledge of the branching fractions of intermediate state decays, which is the uncertainty of  $\mathcal{R}_{\text{int}}$ . The uncertainties from  $\mathcal{R}_{\text{int}}^{J/\psi\phi}$  and  $\mathcal{R}_{\text{int}}^{D_s^-\pi^+}$  are considered uncorrelated. The combined result is

$$\mathcal{R} = \frac{\mathcal{B}(B_c^+ \rightarrow B_s^0\pi^+)}{\mathcal{B}(B_c^+ \rightarrow J/\psi\pi^+)} = 91 \pm 10 \text{ (stat)} \pm 8 \text{ (syst)} \pm 3 \text{ (\mathcal{B})}.$$

This is the first direct measurement of the ratio of branching fractions  $\mathcal{B}(B_c^+ \rightarrow B_s^0\pi^+)/\mathcal{B}(B_c^+ \rightarrow J/\psi\pi^+)$ .

## 7 Discussion of $\mathcal{B}(B_c^+ \rightarrow B_s^0 \pi^+)$

The branching fraction of the  $B_c^+ \rightarrow B_s^0 \pi^+$  decay can be extracted from this result with the knowledge of  $\mathcal{B}(B_c^+ \rightarrow J/\psi \pi^+)$ . The branching fraction of the  $B_c^+ \rightarrow J/\psi \pi^+$  decay is predicted to be  $(0.291^{+0.043}_{-0.050})\%$  [45] with a  $B_c^+$  lifetime of 0.453 ps. Using this input and the  $B_c^+$  lifetime of 0.510 ps taken from ref. [15], the branching fraction of the  $B_c^+ \rightarrow B_s^0 \pi^+$  decay is found to be

$$\mathcal{B}(B_c^+ \rightarrow B_s^0 \pi^+) = (30.0 \pm 3.3 \text{ (stat)} \pm 2.6 \text{ (syst)} \pm 1.0 (\mathcal{B})_{-5.5}^{+4.6} (B_c^+ \rightarrow J/\psi \pi^+)) \%, \quad (7.1)$$

where the fourth uncertainty is the uncertainty from the prediction of  $\mathcal{B}(B_c^+ \rightarrow J/\psi \pi^+)$ .

The branching fraction of the  $B_c^+ \rightarrow J/\psi \pi^+$  decay can also be determined from the theoretical prediction of  $\mathcal{B}(B_c^+ \rightarrow J/\psi \mu^+ \nu_\mu)$  and the measurement of  $\mathcal{B}(B_c^+ \rightarrow J/\psi \pi^+)/\mathcal{B}(B_c^+ \rightarrow J/\psi \mu^+ \nu_\mu) = 0.0469 \pm 0.0028 \pm 0.0046$  [46]. The average of several theoretical predictions for  $\mathcal{B}(B_c^+ \rightarrow J/\psi \mu^+ \nu_\mu)$ , as considered in ref. [47], is  $\mathcal{B}(B_c^+ \rightarrow J/\psi \mu^+ \nu_\mu) = (1.95 \pm 0.46)\%$ . Using these inputs, the value of  $\mathcal{B}(B_c^+ \rightarrow J/\psi \pi^+)$  is found to be  $(0.091 \pm 0.024)\%$ . The branching fraction of the  $B_c^+ \rightarrow B_s^0 \pi^+$  decay is determined to be

$$\mathcal{B}(B_c^+ \rightarrow B_s^0 \pi^+) = (8.3 \pm 0.9 \text{ (stat)} \pm 0.7 \text{ (syst)} \pm 0.3 (\mathcal{B}) \pm 2.2 (B_c^+ \rightarrow J/\psi \pi^+)) \%. \quad (7.2)$$

The fourth uncertainty in eq. (7.2) is due to the uncertainty on the value of  $\mathcal{B}(B_c^+ \rightarrow J/\psi \pi^+)$ .

Equations (7.1) and (7.2) show a discrepancy in  $\mathcal{B}(B_c^+ \rightarrow B_s^0 \pi^+)$  which is larger than the quoted uncertainties, due to different predictions for the value of  $\mathcal{B}(B_c^+ \rightarrow J/\psi \pi^+)$ . Until this discrepancy is resolved, caution is advised when converting the measured ratio into a measurement of  $\mathcal{B}(B_c^+ \rightarrow B_s^0 \pi^+)$ . However, irrespective of this discrepancy, either value of the branching fraction  $\mathcal{B}(B_c^+ \rightarrow J/\psi \pi^+)$  considered here results in the  $B_c^+ \rightarrow B_s^0 \pi^+$  decay having the largest branching fraction of all  $B_c^+$  decays measured to date.

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## The LHCb collaboration

- R. Aaij <sup>32</sup>, A.S.W. Abdelmotteleb <sup>50</sup>, C. Abellan Beteta<sup>44</sup>, F. Abudinén <sup>50</sup>, T. Ackernley <sup>54</sup>,  
 B. Adeva <sup>40</sup>, M. Adinolfi <sup>48</sup>, P. Adlarson <sup>77</sup>, H. Afsharnia<sup>9</sup>, C. Agapopoulou <sup>13</sup>,  
 C.A. Aidala <sup>78</sup>, Z. Ajaltouni<sup>9</sup>, S. Akar <sup>59</sup>, K. Akiba <sup>32</sup>, P. Albicocco <sup>23</sup>, J. Albrecht <sup>15</sup>,  
 F. Alessio <sup>42</sup>, M. Alexander <sup>53</sup>, A. Alfonso Albero <sup>39</sup>, Z. Aliouche <sup>56</sup>, P. Alvarez Cartelle <sup>49</sup>,  
 R. Amalric <sup>13</sup>, S. Amato <sup>2</sup>, J.L. Amey <sup>48</sup>, Y. Amhis <sup>11,42</sup>, L. An <sup>42</sup>, L. Anderlini <sup>22</sup>,  
 M. Andersson <sup>44</sup>, A. Andreianov <sup>38</sup>, M. Andreotti <sup>21</sup>, D. Andreou <sup>62</sup>, D. Ao <sup>6</sup>,  
 F. Archilli <sup>17</sup>, A. Artamonov <sup>38</sup>, M. Artuso <sup>62</sup>, E. Aslanides <sup>10</sup>, M. Atzeni <sup>44</sup>,  
 B. Audurier <sup>12</sup>, I.B Bachiller Perea <sup>8</sup>, S. Bachmann <sup>17</sup>, M. Bachmayer <sup>43</sup>, J.J. Back <sup>50</sup>,  
 A. Bailly-reyre<sup>13</sup>, P. Baladron Rodriguez <sup>40</sup>, V. Balagura <sup>12</sup>, W. Baldini <sup>21,42</sup>,  
 J. Baptista de Souza Leite <sup>1</sup>, M. Barbetti <sup>22,j</sup>, R.J. Barlow <sup>56</sup>, S. Barsuk <sup>11</sup>, W. Barter <sup>52</sup>,  
 M. Bartolini <sup>49</sup>, F. Baryshnikov <sup>38</sup>, J.M. Basels <sup>14</sup>, G. Bassi <sup>29,p</sup>, B. Batsukh <sup>4</sup>,  
 A. Battig <sup>15</sup>, A. Bay <sup>43</sup>, A. Beck <sup>50</sup>, M. Becker <sup>15</sup>, F. Bedeschchi <sup>29</sup>, I.B. Bediaga <sup>1</sup>,  
 A. Beiter <sup>62</sup>, V. Belavin<sup>38</sup>, S. Belin <sup>40</sup>, V. Bellee <sup>44</sup>, K. Belous <sup>38</sup>, I. Belov <sup>38</sup>, I. Belyaev <sup>38</sup>,  
 G. Benane <sup>10</sup>, G. Bencivenni <sup>23</sup>, E. Ben-Haim <sup>13</sup>, A. Berezhnoy <sup>38</sup>, R. Bernet <sup>44</sup>,  
 S. Bernet Andres <sup>76</sup>, D. Berninghoff<sup>17</sup>, H.C. Bernstein<sup>62</sup>, C. Bertella <sup>56</sup>, A. Bertolin <sup>28</sup>,  
 C. Betancourt <sup>44</sup>, F. Betti <sup>42</sup>, Ia. Bezshyiko <sup>44</sup>, S. Bhasin <sup>48</sup>, J. Bhom <sup>35</sup>, L. Bian <sup>68</sup>,  
 M.S. Bieker <sup>15</sup>, N.V. Biesuz <sup>21</sup>, S. Bifani <sup>47</sup>, P. Billoir <sup>13</sup>, A. Biolchini <sup>32</sup>, M. Birch <sup>55</sup>,  
 F.C.R. Bishop <sup>49</sup>, A. Bitadze <sup>56</sup>, A. Bizzeti <sup>10</sup>, M.P. Blago <sup>49</sup>, T. Blake <sup>50</sup>, F. Blanc <sup>43</sup>,  
 J.E. Blank <sup>15</sup>, S. Blusk <sup>62</sup>, D. Bobulska <sup>53</sup>, J.A. Boelhauve <sup>15</sup>, O. Boente Garcia <sup>12</sup>,  
 T. Boettcher <sup>59</sup>, A. Boldyrev <sup>38</sup>, C.S. Bolognani <sup>74</sup>, R. Bolzonella <sup>21,i</sup>, N. Bondar <sup>38,42</sup>,  
 F. Borgato <sup>28</sup>, S. Borghi <sup>56</sup>, M. Borsato <sup>17</sup>, J.T. Borsuk <sup>35</sup>, S.A. Bouchiba <sup>43</sup>,  
 T.J.V. Bowcock <sup>54</sup>, A. Boyer <sup>42</sup>, C. Bozzi <sup>21</sup>, M.J. Bradley <sup>55</sup>, S. Braun <sup>60</sup>,  
 A. Brea Rodriguez <sup>40</sup>, J. Brodzicka <sup>35</sup>, A. Brossa Gonzalo <sup>40</sup>, J. Brown <sup>54</sup>, D. Brundu <sup>27</sup>,  
 A. Buonaaura <sup>44</sup>, L. Buonincontri <sup>28</sup>, A.T. Burke <sup>56</sup>, C. Burr <sup>42</sup>, A. Bursche <sup>66</sup>,  
 A. Butkevich <sup>38</sup>, J.S. Butter <sup>32</sup>, J. Buytaert <sup>42</sup>, W. Byczynski <sup>42</sup>, S. Cadeddu <sup>27</sup>, H. Cai <sup>68</sup>,  
 R. Calabrese <sup>21,i</sup>, L. Calefice <sup>15</sup>, S. Cali <sup>23</sup>, M. Calvi <sup>26,m</sup>, M. Calvo Gomez <sup>76</sup>,  
 P. Campana <sup>23</sup>, D.H. Campora Perez <sup>74</sup>, A.F. Campoverde Quezada <sup>6</sup>, S. Capelli <sup>26,m</sup>,  
 L. Capriotti <sup>20</sup>, A. Carbone <sup>20,g</sup>, G. Carboni <sup>31</sup>, R. Cardinale <sup>24,k</sup>, A. Cardini <sup>27</sup>,  
 P. Carniti <sup>26,m</sup>, L. Carus <sup>14</sup>, A. Casais Vidal <sup>40</sup>, R. Caspary <sup>17</sup>, G. Casse <sup>54</sup>, M. Cattaneo <sup>42</sup>,  
 G. Cavallero <sup>55</sup>, V. Cavallini <sup>21,i</sup>, S. Celani <sup>43</sup>, J. Cerasoli <sup>10</sup>, D. Cervenkov <sup>57</sup>,  
 A.J. Chadwick <sup>54</sup>, I. Chahrour <sup>78</sup>, M.G. Chapman <sup>48</sup>, M. Charles <sup>13</sup>, Ph. Charpentier <sup>42</sup>,  
 C.A. Chavez Barajas <sup>54</sup>, M. Chefdeville <sup>8</sup>, C. Chen <sup>3</sup>, S. Chen <sup>4</sup>, A. Chernov <sup>35</sup>,  
 S. Chernyshenko <sup>46</sup>, V. Chobanova <sup>40</sup>, S. Cholak <sup>43</sup>, M. Chrzaszcz <sup>35</sup>, A. Chubykin <sup>38</sup>,  
 V. Chulikov <sup>38</sup>, P. Ciambrone <sup>23</sup>, M.F. Cicala <sup>50</sup>, X. Cid Vidal <sup>40</sup>, G. Ciezarek <sup>42</sup>,  
 G. Ciullo <sup>i,21</sup>, P.E.L. Clarke <sup>52</sup>, M. Clemencic <sup>42</sup>, H.V. Cliff <sup>49</sup>, J. Closier <sup>42</sup>,  
 J.L. Cobbledick <sup>56</sup>, V. Coco <sup>42</sup>, J.A.B. Coelho <sup>11</sup>, J. Cogan <sup>10</sup>, E. Cogneras <sup>9</sup>,  
 L. Cojocariu <sup>37</sup>, P. Collins <sup>42</sup>, T. Colombo <sup>42</sup>, L. Congedo <sup>19</sup>, A. Contu <sup>27</sup>, N. Cooke <sup>47</sup>,  
 I. Corredoira <sup>40</sup>, G. Corti <sup>42</sup>, B. Couturier <sup>42</sup>, D.C. Craik <sup>44</sup>, M. Cruz Torres <sup>1,e</sup>,  
 R. Currie <sup>52</sup>, C.L. Da Silva <sup>61</sup>, S. Dadabaev <sup>38</sup>, L. Dai <sup>65</sup>, X. Dai <sup>5</sup>, E. Dall'Occo <sup>15</sup>,  
 J. Dalseno <sup>40</sup>, C. D'Ambrosio <sup>42</sup>, J. Daniel <sup>9</sup>, A. Danilina <sup>38</sup>, P. d'Argent <sup>19</sup>, J.E. Davies <sup>56</sup>,  
 A. Davis <sup>56</sup>, O. De Aguiar Francisco <sup>56</sup>, J. de Boer <sup>42</sup>, K. De Bruyn <sup>73</sup>, S. De Capua <sup>56</sup>,

- M. De Cian  $\textcolor{blue}{D}^{43}$ , U. De Freitas Carneiro Da Graca  $\textcolor{blue}{D}^1$ , E. De Lucia  $\textcolor{blue}{D}^{23}$ , J.M. De Miranda  $\textcolor{blue}{D}^1$ ,  
L. De Paula  $\textcolor{blue}{D}^2$ , M. De Serio  $\textcolor{blue}{D}^{19,f}$ , D. De Simone  $\textcolor{blue}{D}^{44}$ , P. De Simone  $\textcolor{blue}{D}^{23}$ , F. De Vellis  $\textcolor{blue}{D}^{15}$ ,  
J.A. de Vries  $\textcolor{blue}{D}^{74}$ , C.T. Dean  $\textcolor{blue}{D}^{61}$ , F. Debernardis  $\textcolor{blue}{D}^{19,f}$ , D. Decamp  $\textcolor{blue}{D}^8$ , V. Dedu  $\textcolor{blue}{D}^{10}$ ,  
L. Del Buono  $\textcolor{blue}{D}^{13}$ , B. Delaney  $\textcolor{blue}{D}^{58}$ , H.-P. Dembinski  $\textcolor{blue}{D}^{15}$ , V. Denysenko  $\textcolor{blue}{D}^{44}$ , O. Deschamps  $\textcolor{blue}{D}^9$ ,  
F. Dettori  $\textcolor{blue}{D}^{27,h}$ , B. Dey  $\textcolor{blue}{D}^{71}$ , P. Di Nezza  $\textcolor{blue}{D}^{23}$ , I. Diachkov  $\textcolor{blue}{D}^{38}$ , S. Didenko  $\textcolor{blue}{D}^{38}$ ,  
L. Dieste Maronas<sup>40</sup>, S. Ding  $\textcolor{blue}{D}^{62}$ , V. Dobishuk  $\textcolor{blue}{D}^{46}$ , A. Dolmatov<sup>38</sup>, C. Dong  $\textcolor{blue}{D}^3$ ,  
A.M. Donohoe  $\textcolor{blue}{D}^{18}$ , F. Dordei  $\textcolor{blue}{D}^{27}$ , A.C. dos Reis  $\textcolor{blue}{D}^1$ , L. Douglas<sup>53</sup>, A.G. Downes  $\textcolor{blue}{D}^8$ , P. Duda  $\textcolor{blue}{D}^{75}$ ,  
M.W. Dudek  $\textcolor{blue}{D}^{35}$ , L. Dufour  $\textcolor{blue}{D}^{42}$ , V. Duk  $\textcolor{blue}{D}^{72}$ , P. Durante  $\textcolor{blue}{D}^{42}$ , M.M. Duras  $\textcolor{blue}{D}^{75}$ , J.M. Durham  $\textcolor{blue}{D}^{61}$ ,  
D. Dutta  $\textcolor{blue}{D}^{56}$ , A. Dziurda  $\textcolor{blue}{D}^{35}$ , A. Dzyuba  $\textcolor{blue}{D}^{38}$ , S. Easo  $\textcolor{blue}{D}^{51}$ , U. Egede  $\textcolor{blue}{D}^{63}$ , V. Egorychev  $\textcolor{blue}{D}^{38}$ ,  
C. Eirea Orro<sup>40</sup>, S. Eisenhardt  $\textcolor{blue}{D}^{52}$ , E. Ejopu  $\textcolor{blue}{D}^{56}$ , S. Ek-In  $\textcolor{blue}{D}^{43}$ , L. Eklund  $\textcolor{blue}{D}^{77}$ , J. Ellbracht  $\textcolor{blue}{D}^{15}$ ,  
S. Ely  $\textcolor{blue}{D}^{55}$ , A. Ene  $\textcolor{blue}{D}^{37}$ , E. Epple  $\textcolor{blue}{D}^{59}$ , S. Escher  $\textcolor{blue}{D}^{14}$ , J. Eschle  $\textcolor{blue}{D}^{44}$ , S. Esen  $\textcolor{blue}{D}^{44}$ , T. Evans  $\textcolor{blue}{D}^{56}$ ,  
F. Fabiano  $\textcolor{blue}{D}^{27,h}$ , L.N. Falcao  $\textcolor{blue}{D}^1$ , Y. Fan  $\textcolor{blue}{D}^6$ , B. Fang  $\textcolor{blue}{D}^{11,68}$ , L. Fantini  $\textcolor{blue}{D}^{72,o}$ , M. Faria  $\textcolor{blue}{D}^{43}$ ,  
S. Farry  $\textcolor{blue}{D}^{54}$ , D. Fazzini  $\textcolor{blue}{D}^{26,m}$ , L.F. Felkowski  $\textcolor{blue}{D}^{75}$ , M. Feo  $\textcolor{blue}{D}^{42}$ , M. Fernandez Gomez  $\textcolor{blue}{D}^{40}$ ,  
A.D. Fernez  $\textcolor{blue}{D}^{60}$ , F. Ferrari  $\textcolor{blue}{D}^{20}$ , L. Ferreira Lopes  $\textcolor{blue}{D}^{43}$ , F. Ferreira Rodrigues  $\textcolor{blue}{D}^2$ ,  
S. Ferreres Sole  $\textcolor{blue}{D}^{32}$ , M. Ferrillo  $\textcolor{blue}{D}^{44}$ , M. Ferro-Luzzi  $\textcolor{blue}{D}^{42}$ , S. Filippov  $\textcolor{blue}{D}^{38}$ , R.A. Fini  $\textcolor{blue}{D}^{19}$ ,  
M. Fiorini  $\textcolor{blue}{D}^{21,i}$ , M. Firlej  $\textcolor{blue}{D}^{34}$ , K.M. Fischer  $\textcolor{blue}{D}^{57}$ , D.S. Fitzgerald  $\textcolor{blue}{D}^{78}$ , C. Fitzpatrick  $\textcolor{blue}{D}^{56}$ ,  
T. Fiutowski  $\textcolor{blue}{D}^{34}$ , F. Fleuret  $\textcolor{blue}{D}^{12}$ , M. Fontana  $\textcolor{blue}{D}^{13}$ , F. Fontanelli  $\textcolor{blue}{D}^{24,k}$ , R. Forty  $\textcolor{blue}{D}^{42}$ ,  
D. Foulds-Holt  $\textcolor{blue}{D}^{49}$ , V. Franco Lima  $\textcolor{blue}{D}^{54}$ , M. Franco Sevilla  $\textcolor{blue}{D}^{60}$ , M. Frank  $\textcolor{blue}{D}^{42}$ , E. Franzoso  $\textcolor{blue}{D}^{21,i}$ ,  
G. Frau  $\textcolor{blue}{D}^{17}$ , C. Frei  $\textcolor{blue}{D}^{42}$ , D.A. Friday  $\textcolor{blue}{D}^{53}$ , J. Fu  $\textcolor{blue}{D}^6$ , Q. Fuehring  $\textcolor{blue}{D}^{15}$ , T. Fulghesu  $\textcolor{blue}{D}^{13}$ ,  
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S. Gambetta  $\textcolor{blue}{D}^{52,42}$ , Y. Gan  $\textcolor{blue}{D}^3$ , M. Gandelman  $\textcolor{blue}{D}^2$ , P. Gandini  $\textcolor{blue}{D}^{25}$ , Y. Gao  $\textcolor{blue}{D}^7$ , Y. Gao  $\textcolor{blue}{D}^5$ ,  
M. Garau  $\textcolor{blue}{D}^{27,h}$ , L.M. Garcia Martin  $\textcolor{blue}{D}^{50}$ , P. Garcia Moreno  $\textcolor{blue}{D}^{39}$ , J. García Pardiñas  $\textcolor{blue}{D}^{26,m}$ ,  
B. Garcia Plana<sup>40</sup>, F.A. Garcia Rosales  $\textcolor{blue}{D}^{12}$ , L. Garrido  $\textcolor{blue}{D}^{39}$ , C. Gaspar  $\textcolor{blue}{D}^{42}$ , R.E. Geertsema  $\textcolor{blue}{D}^{32}$ ,  
D. Gerick<sup>17</sup>, L.L. Gerken  $\textcolor{blue}{D}^{15}$ , E. Gersabeck  $\textcolor{blue}{D}^{56}$ , M. Gersabeck  $\textcolor{blue}{D}^{56}$ , T. Gershon  $\textcolor{blue}{D}^{50}$ ,  
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E.L. Gkougkousis  $\textcolor{blue}{D}^{42}$ , V.V. Gligorov  $\textcolor{blue}{D}^{13,42}$ , C. Göbel  $\textcolor{blue}{D}^{64}$ , E. Golobardes  $\textcolor{blue}{D}^{76}$ , D. Golubkov  $\textcolor{blue}{D}^{38}$ ,  
A. Golutvin  $\textcolor{blue}{D}^{55,38}$ , A. Gomes  $\textcolor{blue}{D}^{1,a}$ , S. Gomez Fernandez  $\textcolor{blue}{D}^{39}$ , F. Goncalves Abrantes  $\textcolor{blue}{D}^{57}$ ,  
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J. Heuel  $\textcolor{blue}{D}^{14}$ , A. Hicheur  $\textcolor{blue}{D}^2$ , D. Hill  $\textcolor{blue}{D}^{43}$ , M. Hilton  $\textcolor{blue}{D}^{56}$ , S.E. Hollitt  $\textcolor{blue}{D}^{15}$ , J. Horswill  $\textcolor{blue}{D}^{56}$ ,  
R. Hou  $\textcolor{blue}{D}^7$ , Y. Hou  $\textcolor{blue}{D}^8$ , J. Hu<sup>17</sup>, J. Hu  $\textcolor{blue}{D}^{66}$ , W. Hu  $\textcolor{blue}{D}^5$ , X. Hu  $\textcolor{blue}{D}^3$ , W. Huang  $\textcolor{blue}{D}^6$ , X. Huang<sup>68</sup>,  
W. Hulsbergen  $\textcolor{blue}{D}^{32}$ , R.J. Hunter  $\textcolor{blue}{D}^{50}$ , M. Hushchyn  $\textcolor{blue}{D}^{38}$ , D. Hutchcroft  $\textcolor{blue}{D}^{54}$ , P. Ibis  $\textcolor{blue}{D}^{15}$ ,  
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K. Ivshin  $\textcolor{blue}{D}^{38}$ , R. Jacobsson  $\textcolor{blue}{D}^{42}$ , H. Jage  $\textcolor{blue}{D}^{14}$ , S.J. Jaimes Elles  $\textcolor{blue}{D}^{41}$ , S. Jakobsen  $\textcolor{blue}{D}^{42}$ , E. Jans  $\textcolor{blue}{D}^{32}$ ,

- B.K. Jashal<sup>41</sup>, A. Jawahery<sup>60</sup>, V. Jevtic<sup>15</sup>, E. Jiang<sup>60</sup>, X. Jiang<sup>4,6</sup>, Y. Jiang<sup>6</sup>, M. John<sup>57</sup>, D. Johnson<sup>58</sup>, C.R. Jones<sup>49</sup>, T.P. Jones<sup>50</sup>, B. Jost<sup>42</sup>, N. Jurik<sup>42</sup>, I. Juszcza<sup>35</sup>, S. Kandybei<sup>45</sup>, Y. Kang<sup>3</sup>, M. Karacson<sup>42</sup>, D. Karpenkov<sup>38</sup>, M. Karpov<sup>38</sup>, J.W. Kautz<sup>59</sup>, F. Keizer<sup>42</sup>, D.M. Keller<sup>62</sup>, M. Kenzie<sup>50</sup>, T. Ketel<sup>32</sup>, B. Khanji<sup>15</sup>, A. Kharisova<sup>38</sup>, S. Kholodenko<sup>38</sup>, G. Khreich<sup>11</sup>, T. Kirn<sup>14</sup>, V.S. Kirsebom<sup>43</sup>, O. Kitouni<sup>58</sup>, S. Klaver<sup>33</sup>, N. Kleijne<sup>29,p</sup>, K. Klimaszewski<sup>36</sup>, M.R. Kmiec<sup>36</sup>, S. Kolliiev<sup>46</sup>, L. Kolk<sup>15</sup>, A. Kondybayeva<sup>38</sup>, A. Konoplyannikov<sup>38</sup>, P. Kopciewicz<sup>34</sup>, R. Kopecna<sup>17</sup>, P. Koppenburg<sup>32</sup>, M. Korolev<sup>38</sup>, I. Kostiuk<sup>32,46</sup>, O. Kot<sup>46</sup>, S. Kotriakhova<sup>16</sup>, A. Kozachuk<sup>38</sup>, P. Kravchenko<sup>38</sup>, L. Kravchuk<sup>38</sup>, R.D. Krawczyk<sup>42</sup>, M. Kreps<sup>50</sup>, S. Kretzschmar<sup>14</sup>, P. Krokovny<sup>38</sup>, W. Krupa<sup>34</sup>, W. Krzemien<sup>36</sup>, J. Kubat<sup>17</sup>, S. Kubis<sup>75</sup>, W. Kucewicz<sup>35</sup>, M. Kucharczyk<sup>35</sup>, V. Kudryavtsev<sup>38</sup>, A. Kups<sup>77</sup>, D. Lacarrere<sup>42</sup>, G. Lafferty<sup>56</sup>, A. Lai<sup>27</sup>, A. Lampis<sup>27,h</sup>, D. Lancierini<sup>44</sup>, C. Landesa Gomez<sup>40</sup>, J.J. Lane<sup>56</sup>, R. Lane<sup>48</sup>, C. Langenbruch<sup>14</sup>, J. Langer<sup>15</sup>, O. Lantwin<sup>38</sup>, T. Latham<sup>50</sup>, F. Lazzari<sup>29,q</sup>, M. Lazzaroni<sup>25,l</sup>, R. Le Gac<sup>10</sup>, S.H. Lee<sup>78</sup>, R. Lefèvre<sup>9</sup>, A. Leflat<sup>38</sup>, S. Legotin<sup>38</sup>, P. Lenisa<sup>i,21</sup>, O. Leroy<sup>10</sup>, T. Lesiak<sup>35</sup>, B. Leverington<sup>17</sup>, A. Li<sup>3</sup>, H. Li<sup>66</sup>, K. Li<sup>7</sup>, P. Li<sup>17</sup>, P.-R. Li<sup>67</sup>, S. Li<sup>7</sup>, T. Li<sup>4</sup>, T. Li<sup>66</sup>, Y. Li<sup>4</sup>, Z. Li<sup>62</sup>, X. Liang<sup>62</sup>, C. Lin<sup>6</sup>, T. Lin<sup>51</sup>, R. Lindner<sup>42</sup>, V. Lisovskyi<sup>15</sup>, R. Litvinov<sup>27,h</sup>, G. Liu<sup>66</sup>, H. Liu<sup>6</sup>, Q. Liu<sup>6</sup>, S. Liu<sup>4,6</sup>, A. Lobo Salvia<sup>39</sup>, A. Loi<sup>27</sup>, R. Lollini<sup>72</sup>, J. Lomba Castro<sup>40</sup>, I. Longstaff<sup>53</sup>, J.H. Lopes<sup>2</sup>, A. Lopez Huertas<sup>39</sup>, S. López Soliño<sup>40</sup>, G.H. Lovell<sup>49</sup>, Y. Lu<sup>4,b</sup>, C. Lucarelli<sup>22,j</sup>, D. Lucchesi<sup>28,n</sup>, S. Luchuk<sup>38</sup>, M. Lucio Martinez<sup>74</sup>, V. Lukashenko<sup>32,46</sup>, Y. Luo<sup>3</sup>, A. Lupato<sup>56</sup>, E. Luppi<sup>21,i</sup>, A. Lusiani<sup>29,p</sup>, K. Lynch<sup>18</sup>, X.-R. Lyu<sup>6</sup>, R. Ma<sup>6</sup>, S. Maccolini<sup>15</sup>, F. Machefert<sup>11</sup>, F. Maciuc<sup>37</sup>, I. Mackay<sup>57</sup>, V. Macko<sup>43</sup>, L.R. Madhan Mohan<sup>48</sup>, A. Maevskiy<sup>38</sup>, D. Maisuzenko<sup>38</sup>, M.W. Majewski<sup>34</sup>, J.J. Malczewski<sup>35</sup>, S. Malde<sup>57</sup>, B. Malecki<sup>35,42</sup>, A. Malinin<sup>38</sup>, T. Maltsev<sup>38</sup>, G. Manca<sup>27,h</sup>, G. Mancinelli<sup>10</sup>, C. Mancuso<sup>11,25,l</sup>, R. Manera Escalero<sup>39</sup>, D. Manuzzi<sup>20</sup>, C.A. Manzari<sup>44</sup>, D. Marangotto<sup>25,l</sup>, J.F. Marchand<sup>8</sup>, U. Marconi<sup>20</sup>, S. Mariani<sup>22,j</sup>, C. Marin Benito<sup>39,42</sup>, J. Marks<sup>17</sup>, A.M. Marshall<sup>48</sup>, P.J. Marshall<sup>54</sup>, G. Martelli<sup>72,o</sup>, G. Martellotti<sup>30</sup>, L. Martinazzoli<sup>42,m</sup>, M. Martinelli<sup>26,m</sup>, D. Martinez Santos<sup>40</sup>, F. Martinez Vidal<sup>41</sup>, A. Massafferri<sup>1</sup>, M. Materok<sup>14</sup>, R. Matev<sup>42</sup>, A. Mathad<sup>44</sup>, V. Matiunin<sup>38</sup>, C. Matteuzzi<sup>26</sup>, K.R. Mattioli<sup>12</sup>, A. Mauri<sup>32</sup>, E. Maurice<sup>12</sup>, J. Mauricio<sup>39</sup>, M. Mazurek<sup>42</sup>, M. McCann<sup>55</sup>, L. Mcconnell<sup>18</sup>, T.H. McGrath<sup>56</sup>, N.T. McHugh<sup>53</sup>, A. McNab<sup>56</sup>, R. McNulty<sup>18</sup>, J.V. Mead<sup>54</sup>, B. Meadows<sup>59</sup>, G. Meier<sup>15</sup>, D. Melnychuk<sup>36</sup>, S. Meloni<sup>26,m</sup>, M. Merk<sup>32,74</sup>, A. Merli<sup>25,l</sup>, L. Meyer Garcia<sup>2</sup>, D. Miao<sup>4,6</sup>, M. Mikhasenko<sup>70,d</sup>, D.A. Milanes<sup>69</sup>, E. Millard<sup>50</sup>, M. Milovanovic<sup>42</sup>, M.-N. Minard<sup>8,†</sup>, A. Minotti<sup>26,m</sup>, T. Miralles<sup>9</sup>, S.E. Mitchell<sup>52</sup>, B. Mitreska<sup>56</sup>, D.S. Mitzel<sup>15</sup>, A. Mödden<sup>15</sup>, R.A. Mohammed<sup>57</sup>, R.D. Moise<sup>14</sup>, S. Mokhnenko<sup>38</sup>, T. Mombächer<sup>40</sup>, M. Monk<sup>50,63</sup>, I.A. Monroy<sup>69</sup>, S. Monteil<sup>9</sup>, G. Morello<sup>23</sup>, M.J. Morello<sup>29,p</sup>, J. Moron<sup>34</sup>, A.B. Morris<sup>42</sup>, A.G. Morris<sup>50</sup>, R. Mountain<sup>62</sup>, H. Mu<sup>3</sup>, E. Muhammad<sup>50</sup>, F. Muheim<sup>52</sup>, M. Mulder<sup>73</sup>, K. Müller<sup>44</sup>, C.H. Murphy<sup>57</sup>, D. Murray<sup>56</sup>, R. Murta<sup>55</sup>, P. Muzzetto<sup>27,h</sup>, P. Naik<sup>48</sup>, T. Nakada<sup>43</sup>, R. Nandakumar<sup>51</sup>, T. Nanut<sup>42</sup>, I. Nasteva<sup>2</sup>, M. Needham<sup>52</sup>, N. Neri<sup>25,l</sup>, S. Neubert<sup>70</sup>, N. Neufeld<sup>42</sup>, P. Neustroev<sup>38</sup>, R. Newcombe<sup>55</sup>, J. Nicolini<sup>15,11</sup>, E.M. Niel<sup>43</sup>, S. Nieswand<sup>14</sup>, N. Nikitin<sup>38</sup>, N.S. Nolte<sup>58</sup>, C. Normand<sup>8,h,27</sup>, J. Novoa Fernandez<sup>40</sup>, G.N Nowak<sup>59</sup>,

- C. Nunez  $\textcolor{blue}{ID}^{78}$ , A. Oblakowska-Mucha  $\textcolor{blue}{ID}^{34}$ , V. Obraztsov  $\textcolor{blue}{ID}^{38}$ , T. Oeser  $\textcolor{blue}{ID}^{14}$ , D.P. O'Hanlon  $\textcolor{blue}{ID}^{48}$ , S. Okamura  $\textcolor{blue}{ID}^{21,i}$ , R. Oldeman  $\textcolor{blue}{ID}^{27,h}$ , F. Oliva  $\textcolor{blue}{ID}^{52}$ , C.J.G. Onderwater  $\textcolor{blue}{ID}^{73}$ , R.H. O'Neil  $\textcolor{blue}{ID}^{52}$ , J.M. Otalora Goicochea  $\textcolor{blue}{ID}^2$ , T. Ovsiannikova  $\textcolor{blue}{ID}^{38}$ , P. Owen  $\textcolor{blue}{ID}^{44}$ , A. Oyanguren  $\textcolor{blue}{ID}^{41}$ , O. Ozcelik  $\textcolor{blue}{ID}^{52}$ , K.O. Paddeken  $\textcolor{blue}{ID}^{70}$ , B. Pagare  $\textcolor{blue}{ID}^{50}$ , P.R. Pais  $\textcolor{blue}{ID}^{42}$ , T. Pajero  $\textcolor{blue}{ID}^{57}$ , A. Palano  $\textcolor{blue}{ID}^{19}$ , M. Palutan  $\textcolor{blue}{ID}^{23}$ , Y. Pan  $\textcolor{blue}{ID}^{56}$ , G. Panshin  $\textcolor{blue}{ID}^{38}$ , L. 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Raven  $\textcolor{blue}{ID}^{33,42}$ , M. Rebollo De Miguel  $\textcolor{blue}{ID}^{41}$ , F. Redi  $\textcolor{blue}{ID}^{42}$ , J. Reich  $\textcolor{blue}{ID}^{48}$ , F. Reiss  $\textcolor{blue}{ID}^{56}$ , C. Remon Alepuz  $\textcolor{blue}{ID}^{41}$ , Z. Ren  $\textcolor{blue}{ID}^3$ , P.K. Resmi  $\textcolor{blue}{ID}^{10}$ , R. Ribatti  $\textcolor{blue}{ID}^{29,p}$ , A.M. Ricci  $\textcolor{blue}{ID}^{27}$ , S. Ricciardi  $\textcolor{blue}{ID}^{51}$ , K. Richardson  $\textcolor{blue}{ID}^{58}$ , M. Richardson-Slipper  $\textcolor{blue}{ID}^{52}$ , K. Rinnert  $\textcolor{blue}{ID}^{54}$ , P. Robbe  $\textcolor{blue}{ID}^{11}$ , G. Robertson  $\textcolor{blue}{ID}^{52}$ , A.B. Rodrigues  $\textcolor{blue}{ID}^{43}$ , E. Rodrigues  $\textcolor{blue}{ID}^{54}$ , E. Rodriguez Fernandez  $\textcolor{blue}{ID}^{40}$ , J.A. Rodriguez Lopez  $\textcolor{blue}{ID}^{69}$ , E. Rodriguez Rodriguez  $\textcolor{blue}{ID}^{40}$ , D.L. Rolf  $\textcolor{blue}{ID}^{42}$ , A. Rollings  $\textcolor{blue}{ID}^{57}$ , P. Roloff  $\textcolor{blue}{ID}^{42}$ , V. Romanovskiy  $\textcolor{blue}{ID}^{38}$ , M. Romero Lamas  $\textcolor{blue}{ID}^{40}$ , A. Romero Vidal  $\textcolor{blue}{ID}^{40}$ , J.D. Roth  $\textcolor{blue}{ID}^{78,\dagger}$ , M. Rotondo  $\textcolor{blue}{ID}^{23}$ , M.S. Rudolph  $\textcolor{blue}{ID}^{62}$ , T. Ruf  $\textcolor{blue}{ID}^{42}$ , R.A. Ruiz Fernandez  $\textcolor{blue}{ID}^{40}$ , J. Ruiz Vidal  $\textcolor{blue}{ID}^{41}$ , A. Ryzhikov  $\textcolor{blue}{ID}^{38}$ , J. Ryzka  $\textcolor{blue}{ID}^{34}$ , J.J. Saborido Silva  $\textcolor{blue}{ID}^{40}$ , N. Sagidova  $\textcolor{blue}{ID}^{38}$ , N. Sahoo  $\textcolor{blue}{ID}^{47}$ , B. Saitta  $\textcolor{blue}{ID}^{27,h}$ , M. Salomoni  $\textcolor{blue}{ID}^{42}$ , C. Sanchez Gras  $\textcolor{blue}{ID}^{32}$ , I. Sanderswood  $\textcolor{blue}{ID}^{41}$ , R. Santacesaria  $\textcolor{blue}{ID}^{30}$ , C. Santamarina Rios  $\textcolor{blue}{ID}^{40}$ , M. Santimaria  $\textcolor{blue}{ID}^{23}$ , E. Santovetti  $\textcolor{blue}{ID}^{31,s}$ , D. Saranin  $\textcolor{blue}{ID}^{38}$ , G. Sarpis  $\textcolor{blue}{ID}^{14}$ , M. Sarpis  $\textcolor{blue}{ID}^{70}$ , A. Sarti  $\textcolor{blue}{ID}^{30}$ , C. Satriano  $\textcolor{blue}{ID}^{30,r}$ , A. Satta  $\textcolor{blue}{ID}^{31}$ , M. Saur  $\textcolor{blue}{ID}^{15}$ , D. Savrina  $\textcolor{blue}{ID}^{38}$ , H. Sazak  $\textcolor{blue}{ID}^9$ , L.G. Scantlebury Smead  $\textcolor{blue}{ID}^{57}$ , A. Scarabotto  $\textcolor{blue}{ID}^{13}$ , S. Schael  $\textcolor{blue}{ID}^{14}$ , S. Scherl  $\textcolor{blue}{ID}^{54}$ , M. Schiller  $\textcolor{blue}{ID}^{53}$ , H. Schindler  $\textcolor{blue}{ID}^{42}$ , M. Schmelling  $\textcolor{blue}{ID}^{16}$ , B. Schmidt  $\textcolor{blue}{ID}^{42}$ , S. Schmitt  $\textcolor{blue}{ID}^{14}$ , O. Schneider  $\textcolor{blue}{ID}^{43}$ , A. Schopper  $\textcolor{blue}{ID}^{42}$ , M. Schubiger  $\textcolor{blue}{ID}^{32}$ , S. Schulte  $\textcolor{blue}{ID}^{43}$ , M.H. Schune  $\textcolor{blue}{ID}^{11}$ , R. Schwemmer  $\textcolor{blue}{ID}^{42}$ , B. Sciascia  $\textcolor{blue}{ID}^{23,42}$ , A. Sciuccati  $\textcolor{blue}{ID}^{42}$ , S. Sellam  $\textcolor{blue}{ID}^{40}$ , A. Semennikov  $\textcolor{blue}{ID}^{38}$ , M. Senghi Soares  $\textcolor{blue}{ID}^{33}$ , A. Sergi  $\textcolor{blue}{ID}^{24,k}$ , N. Serra  $\textcolor{blue}{ID}^{44}$ , L. Sestini  $\textcolor{blue}{ID}^{28}$ , A. Seuthe  $\textcolor{blue}{ID}^{15}$ , Y. Shang  $\textcolor{blue}{ID}^5$ , D.M. Shangase  $\textcolor{blue}{ID}^{78}$ , M. Shapkin  $\textcolor{blue}{ID}^{38}$ , I. Shchemberov  $\textcolor{blue}{ID}^{38}$ , L. Shchutska  $\textcolor{blue}{ID}^{43}$ , T. Shears  $\textcolor{blue}{ID}^{54}$ , L. Shekhtman  $\textcolor{blue}{ID}^{38}$ , Z. Shen  $\textcolor{blue}{ID}^5$ , S. Sheng  $\textcolor{blue}{ID}^{4,6}$ , V. Shevchenko  $\textcolor{blue}{ID}^{38}$ , B. Shi  $\textcolor{blue}{ID}^6$ , Q. Shi  $\textcolor{blue}{ID}^6$ , E.B. Shields  $\textcolor{blue}{ID}^{26,m}$ , Y. Shimizu  $\textcolor{blue}{ID}^{11}$ , E. Shmanin  $\textcolor{blue}{ID}^{38}$ , R. Shorkin  $\textcolor{blue}{ID}^{38}$ , J.D. Shupperd  $\textcolor{blue}{ID}^{62}$ , B.G. Siddi  $\textcolor{blue}{ID}^{21,i}$ , R. Silva Coutinho  $\textcolor{blue}{ID}^{62}$ , G. Simi  $\textcolor{blue}{ID}^{28}$ , S. Simone  $\textcolor{blue}{ID}^{19,f}$ , M. Singla  $\textcolor{blue}{ID}^{63}$ , N. Skidmore  $\textcolor{blue}{ID}^{56}$ , R. Skuza  $\textcolor{blue}{ID}^{17}$ , T. Skwarnicki  $\textcolor{blue}{ID}^{62}$ , M.W. Slater  $\textcolor{blue}{ID}^{47}$ , J.C. Smallwood  $\textcolor{blue}{ID}^{57}$ , J.G. Smeaton  $\textcolor{blue}{ID}^{49}$ , E. Smith  $\textcolor{blue}{ID}^{44}$ , K. Smith  $\textcolor{blue}{ID}^{61}$ , M. Smith  $\textcolor{blue}{ID}^{55}$ , A. Snoch  $\textcolor{blue}{ID}^{32}$ , L. Soares Lavra  $\textcolor{blue}{ID}^9$ , M.D. Sokoloff  $\textcolor{blue}{ID}^{59}$ , F.J.P. Soler  $\textcolor{blue}{ID}^{53}$ , A. Solomin  $\textcolor{blue}{ID}^{38,48}$ , A. Solovev  $\textcolor{blue}{ID}^{38}$ , I. Solovyev  $\textcolor{blue}{ID}^{38}$ , R. Song  $\textcolor{blue}{ID}^{63}$ , F.L. Souza De Almeida  $\textcolor{blue}{ID}^2$ , B. Souza De Paula  $\textcolor{blue}{ID}^2$ , B. Spaan  $\textcolor{blue}{ID}^{15,\dagger}$ , E. Spadaro Norella  $\textcolor{blue}{ID}^{25,l}$ , E. Spedicato  $\textcolor{blue}{ID}^{20}$ , E. Spiridenkov  $\textcolor{blue}{ID}^{38}$ , P. Spradlin  $\textcolor{blue}{ID}^{53}$ , V. Sriskaran  $\textcolor{blue}{ID}^{42}$ , F. Stagni  $\textcolor{blue}{ID}^{42}$ , M. Stahl  $\textcolor{blue}{ID}^{42}$ , S. Stahl  $\textcolor{blue}{ID}^{42}$ , S. Stanislaus  $\textcolor{blue}{ID}^{57}$ , E.N. Stein  $\textcolor{blue}{ID}^{42}$ , O. Steinkamp  $\textcolor{blue}{ID}^{44}$ , O. Stenyakin  $\textcolor{blue}{ID}^{38}$ , H. Stevens  $\textcolor{blue}{ID}^{15}$ , S. Stone  $\textcolor{blue}{ID}^{62,\dagger}$ ,

- D. Strekalina  $\text{\texttt{ID}}^{38}$ , Y.S Su  $\text{\texttt{ID}}^6$ , F. Suljik  $\text{\texttt{ID}}^{57}$ , J. Sun  $\text{\texttt{ID}}^{27}$ , L. Sun  $\text{\texttt{ID}}^{68}$ , Y. Sun  $\text{\texttt{ID}}^{60}$ , P. Svihra  $\text{\texttt{ID}}^{56}$ , P.N. Swallow  $\text{\texttt{ID}}^{47}$ , K. Swientek  $\text{\texttt{ID}}^{34}$ , A. Szabelski  $\text{\texttt{ID}}^{36}$ , T. Szumlak  $\text{\texttt{ID}}^{34}$ , M. Szymanski  $\text{\texttt{ID}}^{42}$ , Y. Tan  $\text{\texttt{ID}}^3$ , S. Taneja  $\text{\texttt{ID}}^{56}$ , M.D. Tat  $\text{\texttt{ID}}^{57}$ , A. Terentev  $\text{\texttt{ID}}^{38}$ , F. Teubert  $\text{\texttt{ID}}^{42}$ , E. Thomas  $\text{\texttt{ID}}^{42}$ , D.J.D. Thompson  $\text{\texttt{ID}}^{47}$ , K.A. Thomson  $\text{\texttt{ID}}^{54}$ , H. Tilquin  $\text{\texttt{ID}}^{55}$ , V. Tisserand  $\text{\texttt{ID}}^9$ , S. T'Jampens  $\text{\texttt{ID}}^8$ , M. Tobin  $\text{\texttt{ID}}^4$ , L. Tomassetti  $\text{\texttt{ID}}^{21,i}$ , G. Tonani  $\text{\texttt{ID}}^{25,l}$ , X. Tong  $\text{\texttt{ID}}^5$ , D. Torres Machado  $\text{\texttt{ID}}^1$ , D.Y. Tou  $\text{\texttt{ID}}^3$ , S.M. Trilov  $\text{\texttt{ID}}^{48}$ , C. Tripli  $\text{\texttt{ID}}^{43}$ , G. Tuci  $\text{\texttt{ID}}^6$ , A. Tully  $\text{\texttt{ID}}^{43}$ , N. Tuning  $\text{\texttt{ID}}^{32}$ , A. Ukleja  $\text{\texttt{ID}}^{36}$ , D.J. Unverzagt  $\text{\texttt{ID}}^{17}$ , A. Usachov  $\text{\texttt{ID}}^{33}$ , A. Ustyuzhanin  $\text{\texttt{ID}}^{38}$ , U. Uwer  $\text{\texttt{ID}}^{17}$ , A. Vagner<sup>38</sup>, V. Vagnoni  $\text{\texttt{ID}}^{20}$ , A. Valassi  $\text{\texttt{ID}}^{42}$ , G. Valenti  $\text{\texttt{ID}}^{20}$ , N. Valls Canudas  $\text{\texttt{ID}}^{76}$ , M. Van Dijk  $\text{\texttt{ID}}^{43}$ , H. Van Hecke  $\text{\texttt{ID}}^{61}$ , E. van Herwijnen  $\text{\texttt{ID}}^{55}$ , C.B. Van Hulse  $\text{\texttt{ID}}^{40,u}$ , M. van Veghel  $\text{\texttt{ID}}^{73}$ , R. Vazquez Gomez  $\text{\texttt{ID}}^{39}$ , P. Vazquez Regueiro  $\text{\texttt{ID}}^{40}$ , C. Vázquez Sierra  $\text{\texttt{ID}}^{42}$ , S. Vecchi  $\text{\texttt{ID}}^{21}$ , J.J. Velthuis  $\text{\texttt{ID}}^{48}$ , M. Veltre  $\text{\texttt{ID}}^{22,t}$ , A. Venkateswaran  $\text{\texttt{ID}}^{43}$ , M. Veronesi  $\text{\texttt{ID}}^{32}$ , M. Vesterinen  $\text{\texttt{ID}}^{50}$ , D. Vieira  $\text{\texttt{ID}}^{59}$ , M. Vieites Diaz  $\text{\texttt{ID}}^{43}$ , X. Vilasis-Cardona  $\text{\texttt{ID}}^{76}$ , E. Vilella Figueras  $\text{\texttt{ID}}^{54}$ , A. Villa  $\text{\texttt{ID}}^{20}$ , P. Vincent  $\text{\texttt{ID}}^{13}$ , F.C. Volle  $\text{\texttt{ID}}^{11}$ , D. vom Bruch  $\text{\texttt{ID}}^{10}$ , A. Vorobyev<sup>38</sup>, V. Vorobyev<sup>38</sup>, N. Voropaev  $\text{\texttt{ID}}^{38}$ , K. Vos  $\text{\texttt{ID}}^{74}$ , C. Vradas  $\text{\texttt{ID}}^{52}$ , R. Waldi  $\text{\texttt{ID}}^{17}$ , J. Walsh  $\text{\texttt{ID}}^{29}$ , G. Wan  $\text{\texttt{ID}}^5$ , C. Wang  $\text{\texttt{ID}}^{17}$ , G. Wang  $\text{\texttt{ID}}^7$ , J. Wang  $\text{\texttt{ID}}^5$ , J. Wang  $\text{\texttt{ID}}^4$ , J. Wang  $\text{\texttt{ID}}^3$ , J. Wang  $\text{\texttt{ID}}^{68}$ , M. Wang  $\text{\texttt{ID}}^{25}$ , R. Wang  $\text{\texttt{ID}}^{48}$ , X. Wang  $\text{\texttt{ID}}^{66}$ , Y. Wang  $\text{\texttt{ID}}^7$ , Z. Wang  $\text{\texttt{ID}}^{44}$ , Z. Wang  $\text{\texttt{ID}}^3$ , Z. Wang  $\text{\texttt{ID}}^6$ , J.A. Ward  $\text{\texttt{ID}}^{50,63}$ , N.K. Watson  $\text{\texttt{ID}}^{47}$ , D. Websdale  $\text{\texttt{ID}}^{55}$ , Y. Wei  $\text{\texttt{ID}}^5$ , B.D.C. Westhenry  $\text{\texttt{ID}}^{48}$ , D.J. White  $\text{\texttt{ID}}^{56}$ , M. Whitehead  $\text{\texttt{ID}}^{53}$ , A.R. Wiederhold  $\text{\texttt{ID}}^{50}$ , D. Wiedner  $\text{\texttt{ID}}^{15}$ , G. Wilkinson  $\text{\texttt{ID}}^{57}$ , M.K. Wilkinson  $\text{\texttt{ID}}^{59}$ , I. Williams<sup>49</sup>, M. Williams  $\text{\texttt{ID}}^{58}$ , M.R.J. Williams  $\text{\texttt{ID}}^{52}$ , R. Williams  $\text{\texttt{ID}}^{49}$ , F.F. Wilson  $\text{\texttt{ID}}^{51}$ , W. Wislicki  $\text{\texttt{ID}}^{36}$ , M. Witek  $\text{\texttt{ID}}^{35}$ , L. Witola  $\text{\texttt{ID}}^{17}$ , C.P. Wong  $\text{\texttt{ID}}^{61}$ , G. Wormser  $\text{\texttt{ID}}^{11}$ , S.A. Wotton  $\text{\texttt{ID}}^{49}$ , H. Wu  $\text{\texttt{ID}}^{62}$ , J. Wu  $\text{\texttt{ID}}^7$ , K. Wyllie  $\text{\texttt{ID}}^{42}$ , Z. Xiang  $\text{\texttt{ID}}^6$ , D. Xiao  $\text{\texttt{ID}}^7$ , Y. Xie  $\text{\texttt{ID}}^7$ , A. Xu  $\text{\texttt{ID}}^5$ , J. Xu  $\text{\texttt{ID}}^6$ , L. Xu  $\text{\texttt{ID}}^3$ , L. Xu  $\text{\texttt{ID}}^3$ , M. Xu  $\text{\texttt{ID}}^{50}$ , Q. Xu  $\text{\texttt{ID}}^6$ , Z. Xu  $\text{\texttt{ID}}^9$ , Z. Xu  $\text{\texttt{ID}}^6$ , D. Yang  $\text{\texttt{ID}}^3$ , S. Yang  $\text{\texttt{ID}}^6$ , X. Yang  $\text{\texttt{ID}}^5$ , Y. Yang  $\text{\texttt{ID}}^6$ , Z. Yang  $\text{\texttt{ID}}^5$ , Z. Yang  $\text{\texttt{ID}}^{60}$ , L.E. Yeomans  $\text{\texttt{ID}}^{54}$ , V. Yeroshenko  $\text{\texttt{ID}}^{11}$ , H. Yeung  $\text{\texttt{ID}}^{56}$ , H. Yin  $\text{\texttt{ID}}^7$ , J. Yu  $\text{\texttt{ID}}^{65}$ , X. Yuan  $\text{\texttt{ID}}^{62}$ , E. Zaffaroni  $\text{\texttt{ID}}^{43}$ , M. Zavertyaev  $\text{\texttt{ID}}^{16}$ , M. Zdybal  $\text{\texttt{ID}}^{35}$ , O. Zenaiev  $\text{\texttt{ID}}^{42}$ , M. Zeng  $\text{\texttt{ID}}^3$ , C. Zhang  $\text{\texttt{ID}}^5$ , D. Zhang  $\text{\texttt{ID}}^7$ , L. Zhang  $\text{\texttt{ID}}^3$ , S. Zhang  $\text{\texttt{ID}}^{65}$ , S. Zhang  $\text{\texttt{ID}}^5$ , Y. Zhang  $\text{\texttt{ID}}^5$ , Y. Zhang<sup>57</sup>, A. Zharkova  $\text{\texttt{ID}}^{38}$ , A. Zhelezov  $\text{\texttt{ID}}^{17}$ , Y. Zheng  $\text{\texttt{ID}}^6$ , T. Zhou  $\text{\texttt{ID}}^5$ , X. Zhou  $\text{\texttt{ID}}^6$ , Y. Zhou  $\text{\texttt{ID}}^6$ , V. Zhovkovska  $\text{\texttt{ID}}^{11}$ , X. Zhu  $\text{\texttt{ID}}^3$ , X. Zhu  $\text{\texttt{ID}}^7$ , Z. Zhu  $\text{\texttt{ID}}^6$ , V. Zhukov  $\text{\texttt{ID}}^{14,38}$ , Q. Zou  $\text{\texttt{ID}}^{4,6}$ , S. Zucchelli  $\text{\texttt{ID}}^{20,g}$ , D. Zuliani  $\text{\texttt{ID}}^{28}$ , G. Zunica  $\text{\texttt{ID}}^{56}$

<sup>1</sup> Centro Brasileiro de Pesquisas Físicas (CBPF), Rio de Janeiro, Brazil<sup>2</sup> Universidade Federal do Rio de Janeiro (UFRJ), Rio de Janeiro, Brazil<sup>3</sup> Center for High Energy Physics, Tsinghua University, Beijing, China<sup>4</sup> Institute Of High Energy Physics (IHEP), Beijing, China<sup>5</sup> School of Physics State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China<sup>6</sup> University of Chinese Academy of Sciences, Beijing, China<sup>7</sup> Institute of Particle Physics, Central China Normal University, Wuhan, Hubei, China<sup>8</sup> Université Savoie Mont Blanc, CNRS, IN2P3-LAPP, Annecy, France<sup>9</sup> Université Clermont Auvergne, CNRS/IN2P3, LPC, Clermont-Ferrand, France<sup>10</sup> Aix Marseille Univ, CNRS/IN2P3, CPPM, Marseille, France<sup>11</sup> Université Paris-Saclay, CNRS/IN2P3, IJCLab, Orsay, France<sup>12</sup> Laboratoire Leprince-Ringuet, CNRS/IN2P3, Ecole Polytechnique, Institut Polytechnique de Paris, Palaiseau, France<sup>13</sup> LPNHE, Sorbonne Université, Paris Diderot Sorbonne Paris Cité, CNRS/IN2P3, Paris, France<sup>14</sup> I. Physikalisches Institut, RWTH Aachen University, Aachen, Germany<sup>15</sup> Fakultät Physik, Technische Universität Dortmund, Dortmund, Germany

- <sup>16</sup> Max-Planck-Institut für Kernphysik (MPIK), Heidelberg, Germany  
<sup>17</sup> Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany  
<sup>18</sup> School of Physics, University College Dublin, Dublin, Ireland  
<sup>19</sup> INFN Sezione di Bari, Bari, Italy  
<sup>20</sup> INFN Sezione di Bologna, Bologna, Italy  
<sup>21</sup> INFN Sezione di Ferrara, Ferrara, Italy  
<sup>22</sup> INFN Sezione di Firenze, Firenze, Italy  
<sup>23</sup> INFN Laboratori Nazionali di Frascati, Frascati, Italy  
<sup>24</sup> INFN Sezione di Genova, Genova, Italy  
<sup>25</sup> INFN Sezione di Milano, Milano, Italy  
<sup>26</sup> INFN Sezione di Milano-Bicocca, Milano, Italy  
<sup>27</sup> INFN Sezione di Cagliari, Monserrato, Italy  
<sup>28</sup> Università degli Studi di Padova, Università e INFN, Padova, Padova, Italy  
<sup>29</sup> INFN Sezione di Pisa, Pisa, Italy  
<sup>30</sup> INFN Sezione di Roma La Sapienza, Roma, Italy  
<sup>31</sup> INFN Sezione di Roma Tor Vergata, Roma, Italy  
<sup>32</sup> Nikhef National Institute for Subatomic Physics, Amsterdam, Netherlands  
<sup>33</sup> Nikhef National Institute for Subatomic Physics and VU University Amsterdam, Amsterdam, Netherlands  
<sup>34</sup> AGH - University of Science and Technology, Faculty of Physics and Applied Computer Science, Kraków, Poland  
<sup>35</sup> Henryk Niewodniczanski Institute of Nuclear Physics Polish Academy of Sciences, Kraków, Poland  
<sup>36</sup> National Center for Nuclear Research (NCBJ), Warsaw, Poland  
<sup>37</sup> Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest-Magurele, Romania  
<sup>38</sup> Affiliated with an institute covered by a cooperation agreement with CERN  
<sup>39</sup> ICCUB, Universitat de Barcelona, Barcelona, Spain  
<sup>40</sup> Instituto Galego de Física de Altas Enerxías (IGFAE), Universidade de Santiago de Compostela, Santiago de Compostela, Spain  
<sup>41</sup> Instituto de Fisica Corpuscular, Centro Mixto Universidad de Valencia - CSIC, Valencia, Spain  
<sup>42</sup> European Organization for Nuclear Research (CERN), Geneva, Switzerland  
<sup>43</sup> Institute of Physics, Ecole Polytechnique Fédérale de Lausanne (EPFL), Lausanne, Switzerland  
<sup>44</sup> Physik-Institut, Universität Zürich, Zürich, Switzerland  
<sup>45</sup> NSC Kharkiv Institute of Physics and Technology (NSC KIPT), Kharkiv, Ukraine  
<sup>46</sup> Institute for Nuclear Research of the National Academy of Sciences (KINR), Kyiv, Ukraine  
<sup>47</sup> University of Birmingham, Birmingham, United Kingdom  
<sup>48</sup> H.H. Wills Physics Laboratory, University of Bristol, Bristol, United Kingdom  
<sup>49</sup> Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom  
<sup>50</sup> Department of Physics, University of Warwick, Coventry, United Kingdom  
<sup>51</sup> STFC Rutherford Appleton Laboratory, Didcot, United Kingdom  
<sup>52</sup> School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom  
<sup>53</sup> School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom  
<sup>54</sup> Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom  
<sup>55</sup> Imperial College London, London, United Kingdom  
<sup>56</sup> Department of Physics and Astronomy, University of Manchester, Manchester, United Kingdom  
<sup>57</sup> Department of Physics, University of Oxford, Oxford, United Kingdom  
<sup>58</sup> Massachusetts Institute of Technology, Cambridge, MA, United States  
<sup>59</sup> University of Cincinnati, Cincinnati, OH, United States  
<sup>60</sup> University of Maryland, College Park, MD, United States  
<sup>61</sup> Los Alamos National Laboratory (LANL), Los Alamos, NM, United States  
<sup>62</sup> Syracuse University, Syracuse, NY, United States  
<sup>63</sup> School of Physics and Astronomy, Monash University, Melbourne, Australia, associated to <sup>50</sup>  
<sup>64</sup> Pontifícia Universidade Católica do Rio de Janeiro (PUC-Rio), Rio de Janeiro, Brazil, associated to <sup>2</sup>

- <sup>65</sup> Physics and Micro Electronic College, Hunan University, Changsha City, China, associated to <sup>7</sup>
- <sup>66</sup> Guangdong Provincial Key Laboratory of Nuclear Science, Guangdong-Hong Kong Joint Laboratory of Quantum Matter, Institute of Quantum Matter, South China Normal University, Guangzhou, China, associated to <sup>3</sup>
- <sup>67</sup> Lanzhou University, Lanzhou, China, associated to <sup>4</sup>
- <sup>68</sup> School of Physics and Technology, Wuhan University, Wuhan, China, associated to <sup>3</sup>
- <sup>69</sup> Departamento de Fisica, Universidad Nacional de Colombia, Bogota, Colombia, associated to <sup>13</sup>
- <sup>70</sup> Universität Bonn - Helmholtz-Institut für Strahlen und Kernphysik, Bonn, Germany, associated to <sup>17</sup>
- <sup>71</sup> Eotvos Lorand University, Budapest, Hungary, associated to <sup>42</sup>
- <sup>72</sup> INFN Sezione di Perugia, Perugia, Italy, associated to <sup>21</sup>
- <sup>73</sup> Van Swinderen Institute, University of Groningen, Groningen, Netherlands, associated to <sup>32</sup>
- <sup>74</sup> Universiteit Maastricht, Maastricht, Netherlands, associated to <sup>32</sup>
- <sup>75</sup> Tadeusz Kosciuszko Cracow University of Technology, Cracow, Poland, associated to <sup>35</sup>
- <sup>76</sup> DS4DS, La Salle, Universitat Ramon Llull, Barcelona, Spain, associated to <sup>39</sup>
- <sup>77</sup> Department of Physics and Astronomy, Uppsala University, Uppsala, Sweden, associated to <sup>53</sup>
- <sup>78</sup> University of Michigan, Ann Arbor, MI, United States, associated to <sup>62</sup>

<sup>a</sup> Universidade de Brasília, Brasília, Brazil

<sup>b</sup> Central South U., Changsha, China

<sup>c</sup> Hangzhou Institute for Advanced Study, UCAS, Hangzhou, China

<sup>d</sup> Excellence Cluster ORIGINS, Munich, Germany

<sup>e</sup> Universidad Nacional Autónoma de Honduras, Tegucigalpa, Honduras

<sup>f</sup> Università di Bari, Bari, Italy

<sup>g</sup> Università di Bologna, Bologna, Italy

<sup>h</sup> Università di Cagliari, Cagliari, Italy

<sup>i</sup> Università di Ferrara, Ferrara, Italy

<sup>j</sup> Università di Firenze, Firenze, Italy

<sup>k</sup> Università di Genova, Genova, Italy

<sup>l</sup> Università degli Studi di Milano, Milano, Italy

<sup>m</sup> Università di Milano Bicocca, Milano, Italy

<sup>n</sup> Università di Padova, Padova, Italy

<sup>o</sup> Università di Perugia, Perugia, Italy

<sup>p</sup> Scuola Normale Superiore, Pisa, Italy

<sup>q</sup> Università di Pisa, Pisa, Italy

<sup>r</sup> Università della Basilicata, Potenza, Italy

<sup>s</sup> Università di Roma Tor Vergata, Roma, Italy

<sup>t</sup> Università di Urbino, Urbino, Italy

<sup>u</sup> Universidad de Alcalá, Alcalá de Henares, Spain

<sup>†</sup> Deceased