

Improving the sound absorption performance of sustainable thermal insulation materials: Natural hemp fibres.

Andrea Santoni^{a,*}, Paolo Bonfiglio^b, Patrizio Fausti^a, Cristina Marescotti^a, Valentina Mazzanti^a,
Francesco Mollica^a, Francesco Pompoli^a

^aDepartment of Engineering, University of Ferrara, via G. Saragat 1, 44122 Ferrara, Italy

^bMateriacustica s.r.l., via Ravera 15/A, 44122 Ferrara, Italy

Abstract

Compared to the traditional synthetic fibrous materials, natural fibres represent sustainable solution to be used either in building construction and noise control engineering and acoustic treatments. Natural fibres are mainly employed in the building industry for their hygrothermal properties, however the possibility to also use them for acoustic purposes would greatly increase their appeal to the market. While synthetic fibres have been studied for almost fifty years, the knowledge of natural fibres is still limited and needs to be expanded. Natural fibres are affected by a large variability of the physical properties, which consequently causes great uncertainty in numerical modelling and difficulties during the design process of acoustics treatments. This study highlights the possibility to enhance the acoustic performance of hemp fibrous materials through the manufacturing process, investigating how each treatment affects the material's physical characteristics and its sound absorption coefficient. Moreover, a simplified model to evaluate the acoustic performance of hemp fibrous materials as a function of their density is proposed, in order to provide a practical tool to investigate and compare different solutions. The physical parameters numerically evaluated for a varying compression rate have been compared with the experimental results, measured at each stage of the production process on samples with a different density and thickness. The global reliability of the proposed approach is finally investigated by comparing the experimental sound absorption for normal incidence with the results obtained from the Johnson-Champoux-Allard model.

Keywords: natural-fibres, sustainable building industry, hemp-fibrous material, sound absorption coefficient, material characterisation

1. Introduction

As sound insulation and noise control have become a primary concern in many industrial fields, several porous and fibrous materials have been developed, by using polymers and petroleum-based products. In recent years, in an attempt to reduce energy consumption in order to encourage sustainable development and preserve resources for future generations, sustainable materials and

*Corresponding author

Email address: andrea.santoni@unife.it (Andrea Santoni)

6 their application have been increasingly studied. The sustainability of a material is generally
7 considered in terms of resources usage, environmental impact, human health, and social equity.
8 Moreover, the production process should require the least amount of energy consumption – green
9 energy is preferable to non-renewable resources – and have the minimum manufacturing waste,
10 or provide second life options for the waste products. Besides traditional petroleum-based ma-
11 terials such as melamine and polyurethane foams, or polyester fibres, there are other materials,
12 commonly used for hygrothermal and acoustic treatments, which need to be sealed although they
13 are derived from natural or recycled products, since they may contain substances impacting on
14 human health if they come into contact with the skin or are inhaled. On the other hand, natu-
15 ral materials such as wood, coconut, kenaf and hemp fibres, obtained from renewable resources
16 by means of a manufacturing process with a reduced impact on the environment [1], are also
17 harmless for human health, since they do not contain toxic substances [2, 3]. For these reasons,
18 compared with traditional materials employed to improve the thermal and acoustic performances
19 in buildings, natural fibres represent an eco-friendly sustainable solution. Natural fibres, such
20 as hemp and kenaf, have been thoroughly studied and used for different applications in building
21 construction [4]. For example in fibre-reinforced concrete or other bio-composites [5] and espe-
22 cially as fibrous thermal insulation materials [6]. Several studies can be found in the literature
23 regarding both the characterisation of hygrothermal and physical properties of these fibres [7, 8]
24 and their applications [9, 10, 11]. Recent studies have shown that natural fibres can also be em-
25 ployed for both for noise control applications providing either a good sound insulation or sound
26 absorption performance [12, 13, 14, 15]. However, a systematic characterisation of their acous-
27 tic behaviour is still lacking. Nevertheless, the possibility to also provide an adequate acoustic
28 performance by employing natural fibres, would certainly increase the market interest, helping
29 to move towards a more sustainable building industry. Nowadays in fact, multi-purposes solu-
30 tions are commonly implemented in building construction, providing for example thermal and
31 acoustic protection [16], or combining light shading with thermal and noise protection [17, 18].

32 In the literature several acoustic models to investigate different kinds of fibrous materials
33 can be found. The most widely used empirical model for mineral and glass wool is probably
34 the one developed by Delany and Bazley [19, 20]. Other models have also been developed in
35 order to investigate different fibrous, porous or granular materials, like, for example, the widely
36 used Johnson-Champoux-Allard equivalent fluid model [21, 22], or the alternative six-parameter
37 model proposed by Lafarge [23], or again the model specifically developed for polyester fibres
38 by Garai and Pompoli [24]. While conventional fibrous materials have been thoroughly investi-
39 gated and the physical parameters which affect their acoustic performance are well known, the
40 knowledge of the physical characteristics of natural fibrous materials and their possible influence
41 on their acoustic performance is still limited. Moreover, an optimisation of the manufacturing
42 process of natural fibres in order to increase their acoustic performance has never been investi-
43 gated. Being generally produced in a lower amount compared to synthetic fibres, either by small
44 enterprises, or sometimes even in artisan workshops, natural fibres are characterised by a signifi-
45 cant variability of their diameter and other physical properties. The fact that the fibres' diameter
46 distribution cannot be arbitrarily controlled like in the production process of synthetic fibres re-
47 presents an issue for the evaluation of their acoustic performance, since all the acoustic models
48 have been developed for homogeneous fibrous materials with a constant diameter [25, 26].

49 In this study an experimental investigation on the variety *Liptko* of the plant species *Cannabis*
50 *Sativa* has been performed. Only few studies, recently published, on the characterisation of hemp
51 fibres can be found in the literature [27, 28]. The aim of this work is to investigate some aspects
52 which, to the authors' best knowledge, have never been analysed before for this kind of fibrous

53 material, even though they may have a significant influence on its acoustic behaviour. As was
54 observed, rough hemp-fibrous materials exhibited a relevantly lower sound absorption compared
55 to synthetic materials such as polyester fibres. Therefore, the influence that the different stages of
56 the manufacturing process have on the physical characteristics and how those affect the material's
57 sound absorption have been analysed, in order to define the mechanical and chemical treatments
58 that optimises the acoustic performance of such eco-friendly material. Furthermore, we propose
59 a simplified approach to characterise the physical parameters of hemp-fibrous material, which
60 requires the evaluation of its sound absorption coefficient, based on the knowledge of the effective
61 fibres' radius and of the material's apparent density. The methodology makes use of a fluid
62 dynamic model [29] to determine, from the experimental air flow resistivity, the effective radius
63 of the equivalent homogeneous fibrous-material. Moreover, from the air flow resistivity and the
64 absorption coefficient for normal incidence of the hemp fibrous material, measured at a given
65 compression rate, it is possible to define all the macroscopic parameters [30] required as input
66 data in the Johnson-Champoux-Allard (JCA) acoustic model. From this initial set of parameters,
67 defined for a given density and thickness, the proposed methodology, inspired by Castagnède's
68 model [31], allows to investigate the material's physical parameters and consequently the sound
69 absorption coefficient for any compression rate. The method was validated by comparing the
70 numerical results evaluated at four different stages of the manufacturing process with the experi-
71 mental data. In the next section the investigated material is introduced and the different stages of
72 the manufacturing process are described. In section 3 the basics of Johnson-Champoux-Allard
73 model are summarised. The methodology proposed to characterise the material's physical pa-
74 rameter is given in detail in section 4, while the main results are validated and discussed in
75 section 5.

76 2. Material and manufacturing process description

77 The investigated hemp material, provided by EmilCanapa s.r.l. (Reggio Emilia, Italy), was
78 of the species *Cannabis Sativa L* and Lipko variety. The innermost part of the hemp stems, often
79 referred to as the pith, is surrounded by woody-fibres known as hurds or shives. The outermost
80 layer, which encloses the hurds, is constituted by bast fibres. These are the fibres that were used
81 in this investigation. After harvesting, the leaves were manually separated from the hemp stems,
82 which were macerated in the field for two months. Subsequently, the hemp stems were worked
83 first with a farm shredder and then with a rotary sieve to divide the bast fibres from the hurds.
84 These separated fibres subsequently underwent four treatments:

- 85 1. **01.CAR – Carding:** this mechanical process is performed in order to break down and
86 untangle long fibres and to remove the remaining traces of dirt and the shortest fibres [32].
- 87 2. **02.NaOH – Alkaline treatment (5%/h):** : the fibres are soaked in a 5 wt.% sodium
88 hydroxide (NaOH) solution at room temperature for 60 minutes. After treatment, the fibres
89 were washed with distilled water, neutralized with 1 wt.% acetic acid solution, and then
90 dried overnight in a vacuum oven at $T = 80^{\circ}\text{C}$. This chemical treatment removed non-
91 cellulosic components and facilitated the extraction of the fibrils, thus improving quality
92 of the fibres [33, 34].
- 93 3. **03.WTC – Wide tooth combing:** a wide tooth comb is used in order to align the fibres,
94 open the fibre bundles and release the fibrils thanks to the shearing stresses induced by the
95 combing process.

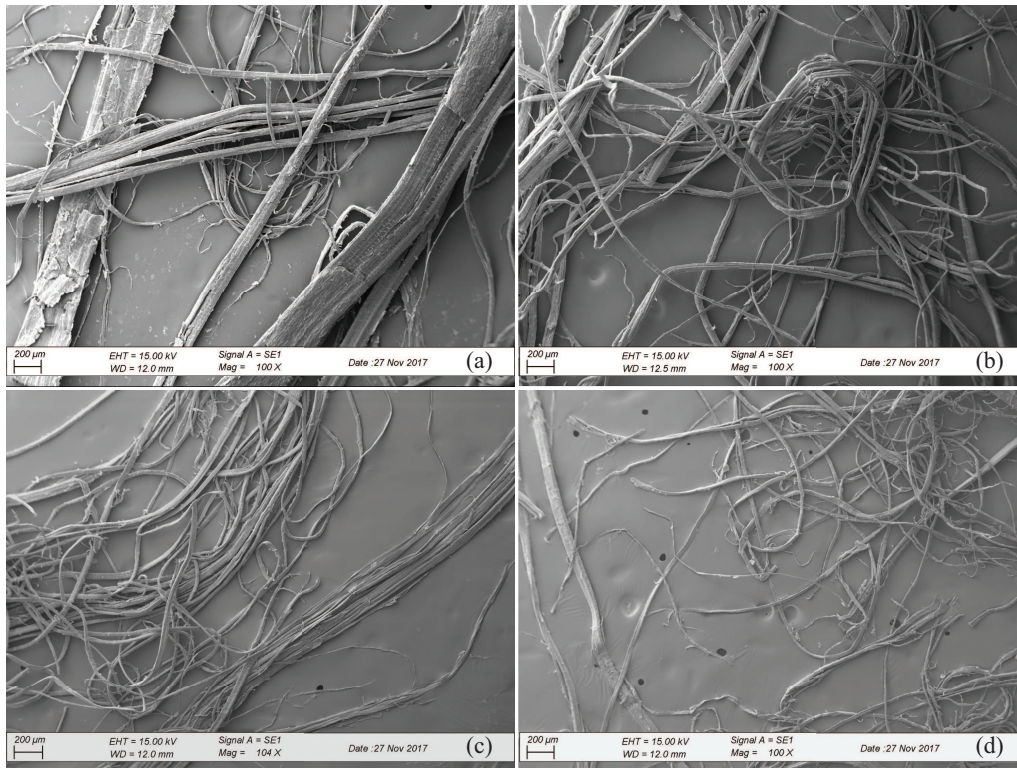


Figure 1: SEM images of hemp fibres: a) hemp fibres after carding process: 01.CAR; b) hemp fibres after alkaline treatment: 02.NaOH; c) hemp fibres after wide tooth combing: 03.WTC; d) hemp fibres after fine tooth combing: 04.FTC.

96 **4. 04.FTC – fine tooth combing:** this last mechanical process is a further refinement of the
97 previous step 03.WTC; in this case a fine tooth comb is used instead of the wide tooth one.

98 After each step of the manufacturing process, samples of hemp fibres were collected and sputter-
99 coated in gold in order to be analysed by a scanning electronic microscopy (SEM). A SEM ZEISS
100 EVO M15 microscope with accelerating voltage of 15 kV was used. Figure 1 shows how each
101 stage of the manufacturing process drastically changed the fibres' morphology. Raw hemp fibres
102 exhibited a complex hierarchical structure constituted by bundles of fibrils glued together by an
103 interphase consisting mainly of non-cellulosic components. The alkalization (02.NaOH) com-
104 bined with mechanical treatments (03.WTC and 04.FTC) removed such cementing substances,
105 unveiling the fibres. For this reason, a significant reduction of the mean diameter of the fibres is
106 shown from step 01.CAR to step 04.FTC, and the amount of fibre bundles decreases [35]. Even
107 though, for all four materials a large dispersion was found in diameters distribution, as shown
108 in Figure 2. Therefore, the determination of a weighted averaged radius [36] is non-trivial and
109 it would certainly be affected by a significant uncertainty. The average radius can be used to
110 compute the physical parameters of a fibrous material. However, an uncertain estimation would
111 provide highly inaccurate results.

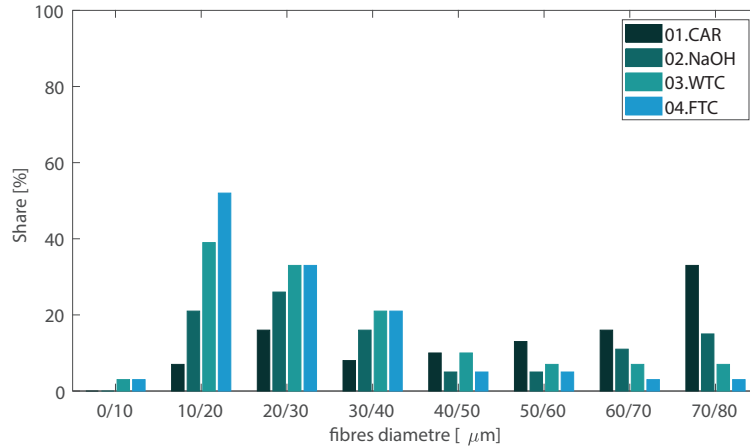


Figure 2: Diameter distribution of hemp fibres after each manufacturing process.

112 3. Acoustical model for sound absorption

113 For a great variety of fibrous and porous materials excited by an incident acoustic wave, it is
 114 possible to assume their solid frame to be rigid, either due to the high value of the elastic modulus,
 115 or the high density, or again because of special test conditions. These media can be described as
 116 equivalent fluids, characterised by an effective density ρ and an effective bulk modulus K . The
 117 well known Johnson-Champoux-Allard (JCA) equivalent fluid model [21, 22] describes visco-
 118 inertial and thermal dissipative effects inside a porous material having a rigid and motionless
 119 frame. Such model involves, in the computation of the effective density ρ and the effective bulk
 120 modulus K , five macroscopic parameters:

- 121 • *airflow resistivity* - σ [Ns/s⁴]: it is the resistance of the material to an airflow passing
 122 through it. Airflow resistivity is determined as:

$$123 \quad \sigma = \frac{\Delta p}{h v_{airflow}} \quad (1)$$

124 where Δp is the pressure drop across the medium while $v_{airflow}$ is the airflow rate passing
 125 through the medium of thickness h .
 126

- 127 • *open porosity* - ϕ [-]: it represents the fractional amount of air volume within the inter-
 128 connected pores in the medium. It can be evaluated as the ratio between the air volume
 129 V_{fluid} and the total volume V_{total} of the investigated material.

$$130 \quad \phi = \frac{V_{fluid}}{V_{tot}} \quad (2)$$

- 131 • *tortuosity* - α_∞ [-]: this dimensionless quantity evaluates the sinuous fluid paths through
 132 the material. It is defined as:

$$133 \alpha_\infty = \frac{\frac{1}{V} \int_V v^2 dV}{\left| \frac{1}{V} \int_V \vec{v} dV \right|^2} \quad (3)$$

134 where v is the microscopic velocity of an inviscid fluid within the pores and V is the
 135 equivalent homogeneous fluid volume.

- 136 • *viscous and thermal characteristic lengths* - Λ, Λ' [μm]: these two quantities describe the
 137 viscous forces and the thermal exchanges between the solid frame and the saturated fluid
 138 contained in it. Their influence is significant at high frequencies. They are defined as:

$$139 \Lambda = 2 \frac{\int_V |v(r)|^2 dV}{\int_S |v(r)|^2 dS} \quad (4)$$

$$140 \Lambda' = 2 \frac{\int_V dV}{\int_S dS} \quad (5)$$

142 where V is the volume of the fluid contained within the pores and S is the interface surface
 143 between the solid frame and the fluid.

144 The effective density ρ of the porous material, which is associated to the inertial and viscous
 145 forces, can be determined as:

$$146 \rho = \frac{\alpha_\infty \rho_0}{\phi} + \frac{\sigma}{j\omega} \sqrt{1 + \frac{4j\alpha_\infty^2 \eta \rho_0 \omega}{\sigma^2 \Lambda^2 \phi^2}} \quad (6)$$

147 where ρ_0 is the air density and η its viscosity, while ω represents the angular frequency and
 148 $j = \sqrt{-1}$ is the imaginary unit. The effective bulk modulus K takes into account the thermal
 149 exchanges between the frame and fluid, and it can be determined as:

$$150 K = \frac{\gamma P_0}{\phi} \left[\gamma - (\gamma - 1) \left(1 + \frac{8\eta}{j\rho_0 \omega N_{Pr} \Lambda'^2} \sqrt{1 + \frac{j\rho_0 \omega N_{Pr} \Lambda'^2}{16\eta}} \right)^{-1} \right] \quad (7)$$

151 being N_{Pr} the Prandtl number, γ the specific heat ratio and P_0 the static pressure. From the
 152 complex effective density ρ and the complex effective bulk modulus, given in Eq. (6) and (7)
 153 respectively, the characteristic impedance Z_c and the complex wavenumber k_c can be computed
 154 as:

$$155 Z_c = \sqrt{\rho K} \quad (8)$$

$$156 k_c = \omega \sqrt{\frac{\rho}{K}} \quad (9)$$

158 Considering a porous material of thickness h placed on a rigid reflecting boundary, the surface
 159 impedance for normal incidence Z_s can be determined as:

$$160 Z_s = -jZ_c \cot(k_c h) \quad (10)$$

Table 1: Effective fluid dynamic radius of the hemp fibres determined at each stage of the manufacturing process.

	01.CAR	02.NaOH	03.WTC	04.FTC
$r_{effect,FD} [\mu\text{m}]$	27.3	29.3	22.7	18.4

161 the normal incidence sound absorption coefficient α_n is finally evaluated as:

$$162 \quad \alpha_n = \frac{4\text{Re}\left\{\frac{Z_s}{\rho_0 c_0}\right\}}{\left|\frac{Z_s}{\rho_0 c_0}\right|^2 + 2\text{Re}\left\{\frac{Z_s}{\rho_0 c_0}\right\} + 1} \quad (11)$$

163 where c_0 represents the speed of sound in air.

164 4. Proposed methodology

165 4.1. Single density material characterisation

166 In order to determine all five physical parameters of the hemp fibrous material at each stage
 167 of the manufacturing process, which are required as input data in the JCA model as described
 168 in the previous paragraph Eq. (1)–(5), an approach based on an equivalent effective radius has
 169 been used. As already mentioned in section 2, the large dispersion which characterises the fibres’
 170 diameter distribution obtained from SEM images would lead to highly inaccurate estimates of a
 171 weighted average radius. Moreover, it should also be considered that with a non-normal distri-
 172 bution, as we obtained for each material, the average value has no physical meaning and is not
 173 a representative value of the distribution. Therefore, in the proposed methodology, an equivalent
 174 effective radius was determined by using a different approach. Based on a fluid dynamic analysis,
 175 Tarnow [37] derived a relationship between the constant fibre radius r and the air flow resistivity
 176 σ , considering randomly distributed fibres and an air flow perpendicular to their axes:

$$177 \quad \sigma = \frac{4\pi\eta}{b^2 \left[0.64 \ln\left(\frac{b^2}{\pi r^2}\right) - 0.737 + \frac{\pi r^2}{b^2}\right]} \quad (12)$$

178 the coefficient b can be computed from the ratio between the fibres’ density ρ_g and the density of
 179 the fibrous material ρ_w as [29]:

$$180 \quad b = r \sqrt{\pi \frac{\rho_g}{\rho_w}} \quad (13)$$

181 Therefore, being the air flow resistivity an easily measurable quantity, the effective radius r of
 182 an equivalent homogeneous material with single diameter fibres can be computed using a sim-
 183 ple minimisation algorithm based on Eq. (12). To this purpose, samples of loose hemp fibres,
 184 resulting from each manufacturing process described in section 2, were tested in the acoustic labo-
 185 ratories of the University of Ferrara. For each sample, the air flow resistivity σ was measured
 186 by means of the alternate flow method, as described in the EN 29053:1993 standard [38]. The
 187 effective radius r obtained for each material is reported in Table 1. Moreover, the normal inci-
 188 dence sound absorption coefficient α_n was measured by using a well-established transfer function
 189 method in an impedance tube, according to the ISO 10534-2:1998 standard [39]. The cylindrical



Figure 3: Investigated hemp fibres: a) loose hemp fibres; b) a fibrous material sample is created within the test rig; c) metallic mesh used to restrain the fibres in the sample holder.

190 samples were prepared by compressing the loose hemp fibres into the sample holder of each mea-
 191 surement test rig, as shown in Figure 3. A coarse metallic mesh was used to restrain the fibres
 192 when compressed at a chosen rate, in order to obtain samples with a constant density and also
 193 to prevent any leakage around the edge. For each processing stage, three different measurements
 194 were performed; each time the sample was removed and reinserted in the test rig. The porosity
 195 ϕ of the fibrous material was computed from the mass of loose hemp fibres, the volume of the
 196 sample holder it was contained within and the density of the fibres, according to the relationship:

$$197 \quad \phi = 1 - \frac{\rho_w}{\rho_g} \quad (14)$$

198 The density of the fibres was estimated from the literature: $\rho_g = 1300 \text{ kg/m}^3$. The tortuosity can
 199 be evaluated as a function of the material porosity as [40]:

$$200 \quad \alpha_\infty = \left(\frac{1}{\phi}\right)^{0.7659} \quad (15)$$

201 It is possible to estimate the thermal characteristic length Λ' as a function of the mean square
 202 radius of the hemp fibres r [31], evaluated from the measured air flow resistivity using Eq. 12 as:

$$203 \quad \Lambda' = b - r \quad (16)$$

204 By knowing these four physical parameters, the viscous characteristic length Λ can be finally
 205 evaluated from the measured absorption coefficient α_n by using a well established inversion
 206 method [41].

207 4.2. Material characterisation varying the compression rate

208 The analysis was further extended to the investigation of the possibility to compute the phys-
 209 ical parameters of hemp fibrous material with an arbitrary compression rate. For example, the
 210 formulations developed by Castagnède et al. [31], both for 1D and 2D compression, allow to

211 evaluate all five physical parameters for any compression rate $n = \rho_{w,(n)}/\rho_{w,(0)}$ by knowing the
 212 material's characteristics for a given density $\rho_{w,(0)}$ and thickness $h_{(0)}$. The great advantage of this
 213 approach is certainly represented by the fact that it can be easily applied in practical contexts, re-
 214 quiring few simple experimental measurements. However, it should be mentioned that the linear
 215 relationships, developed for 1D compression, have been validated only for a material with high
 216 internal porosity within the range $\phi = 0.944 \div 0.995$. In fact, it has been shown [42, 43] that
 217 Castagnède et al.'s formulas for 1D compression provide reliable results when a small compression
 218 rate and highly porous materials are considered, although it may be inaccurate to investigate
 219 high density materials. Moreover, this approach does not take into account the variation of the
 220 fibres' orientation due to the compression, which has been proven to have a significant influence
 221 on the physical properties of the material [43]. Nevertheless, Castagnède's model, due to its
 222 simplicity and straightforward applicability, represents a good starting point to define an empiri-
 223 cal tool to be used in order to characterise hemp fibres, which have never been systematically
 224 analysed before from an acoustic point of view. To this purpose, the linear relationships associ-
 225 ated to mono-axial 1D compression, to define the physical parameters of the JCA model, have
 226 been used. However, a correction coefficient has been introduced as exponential of the compression
 227 rate in the air flow resistivity formulation. Instead of the linear relationship proposed by
 228 Castagnède for 1D compression, the compression rate n is raised to the power of $A = 2.1337$.
 229 Such coefficient was determined through a least square minimisation of the percentage error be-
 230 tween the computed air flow resistivity and the experimental data, measured on samples of five
 231 different thicknesses for each investigated material. Indicating with the subscript (0) the initial set
 232 of physical parameters, determined as described in the previous section, the air flow resistivity
 233 of hemp fibrous material at any given compression rate n , indicated with the subscript (n) , was
 234 evaluated as:

$$\sigma_{(n)} = n^A \sigma_{(0)} \quad (17)$$

236 It should be mentioned that an alternative set of equations, developed using numerical fluid
 237 analyses, assuming that only the air filling the porosity changes during compression while the
 238 cross-sectional area of the fibres remains undeformed (which is equivalent to the assumptions
 239 adopted by Castagnède for 1D compression), was also recently presented by Hirosawa and Nak-
 240 agawa [44]. In this case the airflow resistivity does not vary linearly with the compression rate,
 241 even though, as will be shown in the next section, this equation does not provide accurate results
 242 at the higher densities either. The other physical parameters of the hemp fibrous material, such
 243 as its tortuosity α_∞ , and its viscous and thermal characteristic lengths Λ and Λ' , were computed
 244 for a varying compression rate by using Castagnède's formulations:

$$\alpha_{\infty,(n)} = 1 - n(1 - \alpha_{\infty,(0)}) \quad (18)$$

$$\Lambda_{(n)} = \frac{\Lambda_{(0)}}{n^{1/2}} + r \left(\frac{1}{n^{1/2}} - 1 \right) \quad (19)$$

$$\Lambda'_{(n)} = \frac{\Lambda'_{(0)}}{n^{1/2}} + r \left(\frac{1}{n^{1/2}} - 1 \right) \quad (20)$$

250 The material's porosity was determined as the ratio between the material's density and the fibres'
 251 density, as provided in Eq. (14); alternatively it can also be expressed as:

$$\phi_{(n)} = 1 - n(1 - \phi_{(0)}) \quad (21)$$

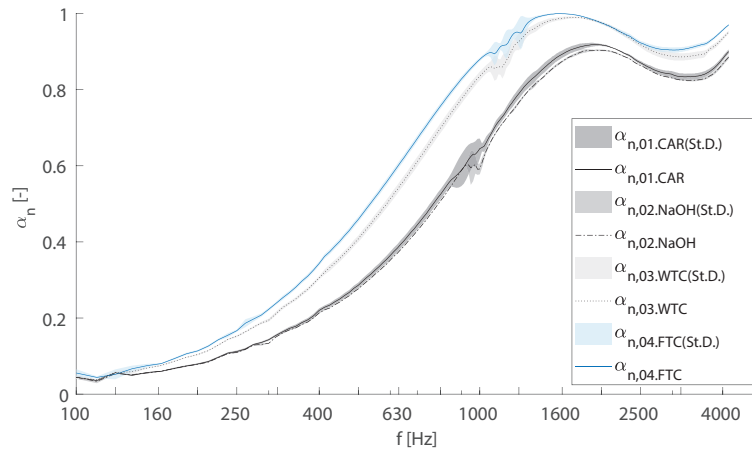


Figure 4: Normal incidence sound absorption coefficient measured on hemp fibre samples, with thickness $h = 40$ mm and density $\rho_w = 88 \text{ kg/m}^3$ at the four different stages of the manufacturing process.

253 The hemp fibrous material was characterised for any compression rate within the range $n = 0.5 \div$
254 2, starting from the physical parameters determined as a function of the fluid dynamic effective
255 radius for the material with density $\rho_{w,(0)} = 88 \text{ kg/m}^3$ and thickness $h_{(0)} = 40$ mm.

256 5. Results and validation

257 The way in which each manufacturing process, presented in section 2, affects the acoustic
258 performance of the material was analysed by comparing the normal incidence sound absorption
259 coefficient α_n , measured on samples of loose hemp fibres at the four different stages of the process,
260 with identical density $\rho = 88 \text{ kg/m}^3$ and thickness $h = 40$ mm. The sound absorption
261 measurements were performed, as described in section 4, three times for each sample; the exper-
262 imental standard deviation is reported as a shaded area around the average curve of the associated
263 absorption coefficient. As shown in Figure 4, there is no relevant difference between the fibres
264 only carded 01.CAR and the fibres which also went through the alkaline treatment 02.NaOH. In
265 fact, even though the sound absorption coefficient associated with the 02.NaOH fibres is slightly
266 lower than the values related to the carded fibres 01.CAR, these differences are barely noticeable
267 and the two curves almost superpose within the entire frequency range. On the other hand, the
268 alkaline treatment 02.NaOH combined with the wide tooth combing 03.WTC significantly in-
269 creases the sound absorption of the material. This effect is due to the noteworthy reduction of the
270 mean diameter caused by a substantial morphological change of the fibre structure. In fact, only
271 the combination of the elimination of non-cellulosic components (02.NaOH) combined with a
272 mechanical treatment (03.WTC) allows the opening of the bundles and the release of the fibrils
273 that have a smaller diameter. The acoustic performance of the hemp fibrous material is slightly
274 further enhanced after the fine tooth combing treatment 04.FTC, which allows the natural fibrous
275 material to match the normal incidence sound absorption of traditional synthetic fibres.

Table 2: Physical parameters of the hemp fibrous materials used as initial data set to characterise the material properties at varying compression rate

	01.CAR	02.NaOH	03.WTC	04.FTC
$h_{(0)}$ [mm]	40	40	40	40
$\rho_{(0)}$ [kg/m ³]	88	88	88	88
$\phi_{(0)}$ [-]	0.93	0.93	0.93	0.93
$\sigma_{(0)}$ [Pas/m ²]	5536	4920	7883	12503
$\alpha_{\infty,(0)}$ [-]	1.05	1.05	1.05	1.05
$\Lambda_{(0)}$ [μ m]	109	115	59	50
$\Lambda'_{(0)}$ [μ m]	160	170	135	109

276 The characterisation methodology described in the previous section was validated by compar-
 277 ing the physical parameters, computed for a varying compression rate n , with experimental
 278 results. The initial set of material properties, provided in Table 2, was computed as a func-
 279 tion of the effective radius, derived from the experimental air flow resistivity $\sigma_{(0)}$, or from the
 280 normal incidence sound absorption coefficient $\alpha_{n,(0)}$; both measured on samples of hemp fibres
 281 40 mm thick, with a density of $\rho_{(0)} = 88 \text{ kg/m}^3$. In order to validate the results, additional
 282 measurements were made on the hemp fibrous material at each stage of the manufacturing pro-
 283 cess. In particular, a total of five different densities ρ_w were experimentally tested within the
 284 range $\rho_{w,i} = 59 \div 141 \text{ kg/m}^3$, obtaining hemp fibrous samples with a thickness varying from
 285 $h = 60 \text{ mm}$ to $h = 25 \text{ mm}$. For each sample the normal incidence absorption coefficient and
 286 the air flow resistivity were experimentally measured as described in section 4. Moreover, the
 287 materials' tortuosity α_{∞} was assessed from an ultrasonic experimental method [45], while the
 288 viscous Λ and thermal characteristic lengths Λ' were evaluated by using well consolidated inver-
 289 sion techniques based on a least mean square algorithm [41].

290 Figure 5 shows the experimental air flow resistivity, measured on samples of the fibrous mater-
 291 ial, with five different densities, after each processing step. Results were found to be consistent
 292 with the sound absorption coefficients, shown in Figure 4. In fact, for all the investigated den-
 293 sities the fibres resulting from the process 04.FTC exhibit the higher air flow resistivity, while the
 294 lowest values are associated with the fibres resulting from the alkaline treatment 02.NaOH; this
 295 behaviour is more emphasised as the density increases. An increase in air flow resistivity is thus
 296 associated with an improvement of the material's absorption coefficient. The experimental air
 297 flow resistivity, plotted together with the error bars representing the experimental standard devi-
 298 ation, are compared in Figure 5, with the air flow resistivity computed for a varying compression
 299 rate $n = 0.5 \div 2$, according to Eq. (17). For each material a good agreement is found between the
 300 experimental results and the curve computed using the formulation proposed in this paper $\sigma_{i,(n)}$.
 301 The results obtained from the linear model proposed by Castagnède for 1D compression is also
 302 reported as a dotted line $\sigma_{i,(n,1D)}$, in order to demonstrate how it is not suitable for this kind of ma-
 303 terial, providing inaccurate results which progressively deviate from the experimental evidence
 304 as the density increases. The air flow resistivity $\sigma_{i,(n,H)}$ computed using the non-linear equation

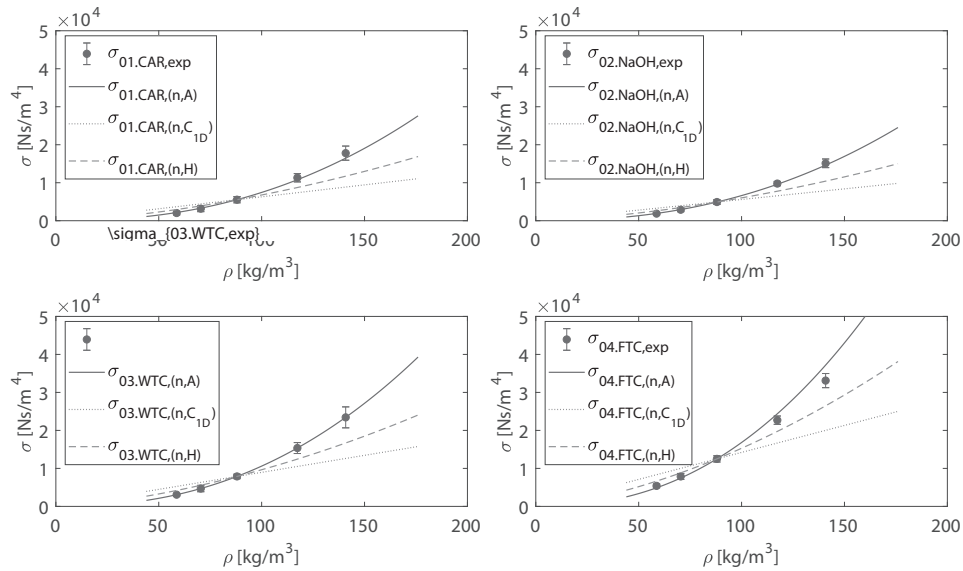


Figure 5: Experimental and estimated air flow resistivity evaluated at each stage of the manufacturing process on samples of various densities at the four manufacturing stages: (a) carding process: 01.CAR; (b) alkaline treatment: 02.NaOH; (c) wide tooth combing: 03.WTC; (d) fine tooth combing: 04.FTC.

305 developed by Hirosawa et al. is also reported. A slightly better agreement is found for a low
 306 compression rate. Nevertheless, at the higher densities it clearly deviates from the experimental
 307 results.

308 On the other hand, Castagnède's equation for 1D compression, which in this case is equiv-
 309 alent to the one developed by Hirosawa, provides a good approximation of the experimental
 310 tortuosity, as shown in Figure 6. However, the curve obtained from the linear relationship for
 311 1D compression slightly deviates from the experimental results at the highest densities. It should
 312 be considered that such differences are comparable with the experimental standard deviation,
 313 reported as error bars; besides, such variations have a negligible influence on sound absorption
 314 performance.

315 It was not possible to measure the viscous and the thermal characteristics lengths. These
 316 quantities were thus determined, for each material, by using a consolidated inversion technique,
 317 minimising the difference between the experimental sound absorption coefficient and the results
 318 obtain from the JCA model [41].

319 However, keeping in mind that at the low frequencies, or for materials with a low density,
 320 both the viscous and the thermal characteristics lengths do not have a significant influence on the
 321 sound absorption coefficient, it is evident in these cases the inversion technique would not nec-
 322 essarily provide the best physical solution. A sensitivity analysis of these quantities on the JCA
 323 model, was carried out as a function of the material density. Numerical results were compared
 324 with all the values of Λ and Λ' , obtained from the minimisation algorithm, which guarantee an
 325 absorption coefficient within 3% of error with respect to the experimental sound absorption coef-
 326 ficient. . As shown in Figure 7, the standard deviation associated with the characteristic viscous
 327 length, plotted as a shaded area around the best fit results, is very limited for all four materials,

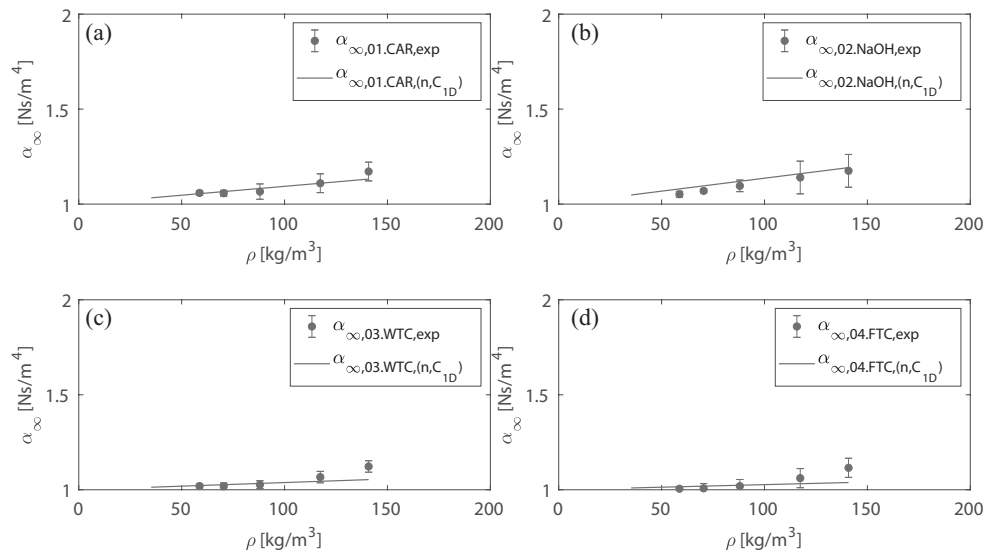


Figure 6: Experimental and estimated tortuosity evaluated on samples of various densities at the four manufacturing stages: (a) carding process: 01.CAR; (b) alkaline treatment: 02.NaOH; (c) wide tooth combing: 03.WTC; (d) fine tooth combing: 04.FTC.

328 both for high and low densities. However, as the porosity increases, the influence of the fluid
 329 viscosity is more significant and the standard deviation decreases further. The numerical curve
 330 for a varying compression rate $\Lambda_{(n,1D)}$, given in Figure 7, is in good agreement with the experi-
 331 mental values of the material resulting from the two last manufacturing processes: 03.WTC and
 332 04.FTC. However, it significantly deviates from the characteristic viscous length evaluated from
 333 experimental data for the fibres which underwent only the first two manufacturing treatments:
 334 01.CAR and 02.NaOH. As shown in Figure 8, an analogous picture can be drawn comparing
 335 the thermal characteristic length derived from the experimental data set and the numerical curve
 336 evaluated for a varying compression rate. As was found for the viscous characteristic length,
 337 Castagnède's equation for 1D compression provides an accurate approximation for the materials
 338 obtained from the last two manufacturing processes: 03.WTC and 04.FTC, while relevant dis-
 339 crepancies are found for the materials at the first two stages: 01.CAR and 02.NaOH. However,
 340 by looking at results obtained from the inversion algorithm, considering all the solutions which
 341 provide an absorption coefficient within 3% error, a huge standard deviation is found. This means
 342 that the thermal characteristic length does not significantly affect the sound absorption coefficient
 343 of materials with high porosity or with rough fibres. Although the numerical results deviate from
 344 the best value solution, represented by black circles, were well within the shaded area which
 345 represents the standard deviation obtained in the minimisation approach. Both characteristics
 346 lengths were also computed at a varying compression rate by using the formulations provided
 347 by Hirose, which, unlike in Castagnède's model, do not depend on the effective radius of the
 348 fibres. These results are indicated in Figures 7 and 8 as $\Lambda_{(n,H)}$ and $\Lambda'_{(n,H)}$ respectively. Even
 349 though small differences are observed, this model does not provide a significant better accuracy
 350 compared to experimental results.

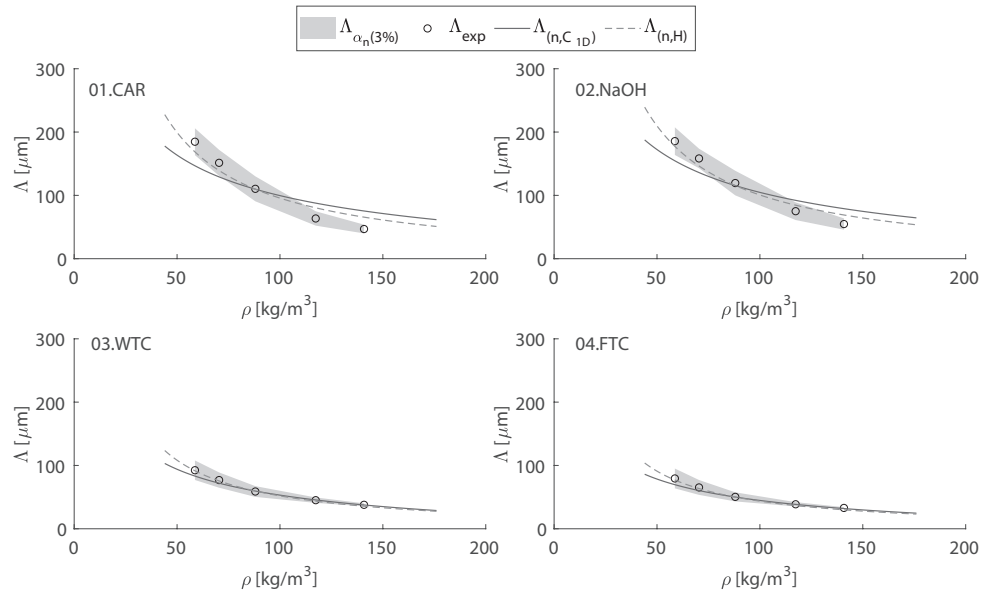


Figure 7: Experimental determined from inversion technique and estimated characteristic viscous length evaluated on samples of various densities at the four manufacturing stages: (a) carding process: 01.CAR; (b) alkaline treatment: 02.NaOH; (c) wide tooth combing: 03.WTC; (d) fine tooth combing: 04.FTC.

351 Finally, in order to determine whether the accuracy provided by this approach may guarantee
 352 a good approximation of the material's acoustic performance, the experimental sound absorption
 353 form normal incidence α_n , measured at each stage of the manufacturing process for five differ-
 354 ent compression rates, was compared to the results obtained from the JCA model. The physical
 355 parameters, numerically computed for a varying compression rate, were used as input data in
 356 the model. As shown in Figure 9, a good agreement is found between the numerical results and
 357 the experimental sound absorption for all four materials and for each of the investigated com-
 358 pression rates. Since the small discrepancies highlighted in some cases between numerical and
 359 experimental curves are comparable with the experimental standard deviation found between
 360 different measurements of the same sample, shown in Figure 4, it can be concluded that the char-
 361 acterisation approach investigated in this study provides an accurate estimation of the physical
 362 parameters required to describe the acoustic performance of hemp fibre materials using the JCA
 363 model.

364 6. Conclusion

365 A study regarding the acoustic performance of hemp fibrous materials and the physical pa-
 366 rameters by which it is affected has been presented. An analysis was carried out of how to
 367 optimise the manufacturing process in order to obtain a natural, sustainable and renewable fi-
 368 brous material, which can provide a sound absorbing performance comparable that provided by
 369 traditional synthetic fibres. An experimental investigation on hemp-fibres identified the influence
 370 of each stage of the manufacturing process both on the acoustic performance and on the physical

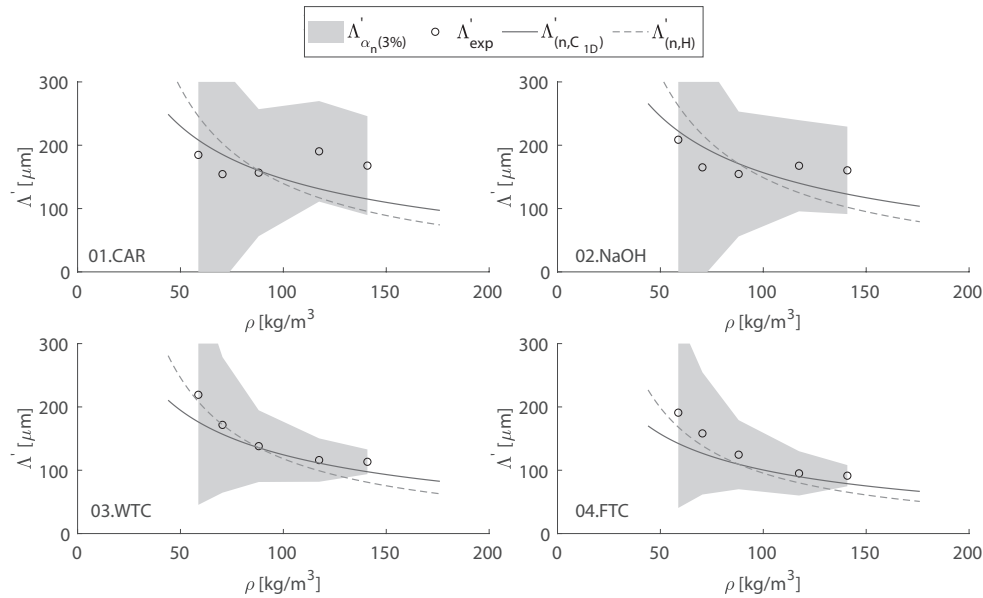


Figure 8: Experimental determined from inversion technique and estimated characteristic thermal length evaluated on samples of various densities at the four manufacturing stages: (a) carding process: 01.CAR; (b) alkaline treatment: 02.NaOH; (c) wide tooth combing: 03.WTC; (d) fine tooth combing: 04.FTC.

371 characteristics of the fibrous material. It was shown that an alkaline treatment (02.NaOH) per-
 372 formed on the material after the carding process (01.CAR) does not significantly affect either its
 373 acoustic performance or its physical properties, such as the airflow resistivity for example, unless
 374 this is followed by two combing processes. The first one is made with a wide tooth comb, while
 375 the second with a finer one. These processes allow to improve the material acoustic performance,
 376 by increasing air flow resistivity and reducing the effective radius of the fibres. The first combing
 377 represents a fundamental step of the process, while the last one is a refinement which allows a
 378 small further improvement of the performance. From a physical point of view, the combing pro-
 379 cess, performed after alkaline treatment, modifies the morphology of the fibre by promoting the
 380 opening of the bundles and freeing the fibrils. These substructures are in fact characterized by
 381 a smaller diameter. The physical parameters required to describe the hemp fibres using the JCA
 382 equivalent fluid model were characterised after each manufacturing procedure. Since natural fi-
 383 bres are characterised by large variability the distribution of diameters, the proposed approach is
 384 based on the concept of effective equivalent fluid-dynamic radius, derived from the experimental
 385 air flow resistivity by using the model developed by Tarnow. Besides, a simplified methodology
 386 to investigate the acoustic performance of hemp fibrous materials, as a function of the material
 387 density, was then proposed, providing a useful tool to compute and compare the sound propaga-
 388 tion into the material with different degrees of compression. The properties of the hemp fibrous
 389 material for a varying compression rate were evaluated with good accuracy by using the simple
 390 model developed by Castagnède for mono-axial compression, from the parameters characterised
 391 for a given density, by means of a fluid-dynamic approach. However, in order to consider also
 392 materials with a high compression rate, i.e. high density and low porosity, or aspects which

393 are not taken into account in this simplified model, such as the fibres' orientation, the 1D linear
394 equation provided to compute the air flow resistivity was modified by introducing an exponen-
395 tial correction coefficient. Such term was determined from the airflow resistivity measured on
396 hemp fibre samples at different compression rates by means of a minimisation algorithm. At each
397 manufacturing stage, all of the investigated parameters numerically evaluated were validated by
398 comparison with the experimental results, measured on hemp fibre samples with five different
399 densities and thicknesses. More specifically, the air flow resistivity, the tortuosity and the normal
400 incidence sound absorption were directly measured, while the viscous and thermal character-
401 istic lengths were determined from the experimental data set using an inversion technique. The
402 validation highlighted a good agreement between the results obtained from the proposed method-
403 ology and the experimental physical parameters. Due to the simplified and practical nature of
404 the proposed approach some discrepancies were found, especially at the highest densities and
405 for the materials with rough fibres, such as the ones obtained from the first two stages of the
406 manufacturing process. However, the differences found between the experimental and numerical
407 results were comparable with the experimental standard deviation obtained from different mea-
408 surements of the same sample. Moreover, numerically evaluated parameters used as input data in
409 the Johnson-Champoux-Allard model provided a very good approximation of the experimental
410 sound absorption coefficient measured at each stage of the manufacturing process on materials at
411 five different compression rates. The proposed model can certainly be refined in the follow-up of
412 this project, even though this simplified approach represents a simple and reliable tool to analyse
413 the acoustic performance of hemp fibrous materials, which has not been investigated before.

414 **Acknowledgement**

415 The authors are grateful to Mr. Jonathan Vallari of EmilCanapa for providing the hemp fibres
416 used in this study.

417 **References**

418 **References**

- 419 [1] F. Ardente, M. Beccali, M. Cellura, M. Mistretta, Building energy performance: a lca case study of kenaf-fibres
420 insulation board, *Energy and Buildings* 40 (1) (2008) 1–10.
- 421 [2] S. V. Joshi, L. Drzal, A. Mohanty, S. Arora, Are natural fiber composites environmentally superior to glass fiber
422 reinforced composites?, *Composites Part A: Applied science and manufacturing* 35 (3) (2004) 371–376.
- 423 [3] C. Zhou, S. Q. Shi, Z. Chen, L. Cai, L. Smith, Comparative environmental life cycle assessment of fiber reinforced
424 cement panel between kenaf and glass fibers, *Journal of Cleaner Production* 200 (2018) 196–204.
- 425 [4] C. Ingrao, A. L. Giudice, J. Bacenetti, C. Tricase, G. Dotelli, M. Fiala, V. Siracusa, C. Mbohwa, Energy and envi-
426 ronmental assessment of industrial hemp for building applications: A review, *Renewable and Sustainable Energy*
427 *Reviews* 51 (2015) 29–42.
- 428 [5] E. Awwad, M. Mabsout, B. Hamad, M. T. Farran, H. Khatib, Studies on fiber-reinforced concrete using industrial
429 hemp fibers, *Construction and Building Materials* 35 (2012) 710–717.
- 430 [6] H.-R. Kymäläinen, A.-M. Sjöberg, Flax and hemp fibres as raw materials for thermal insulations, *Building and*
431 *environment* 43 (7) (2008) 1261–1269.
- 432 [7] A. Bourdot, T. Moussa, A. Gacoin, C. Maalouf, P. Vazquez, C. Thomachot-Schneider, C. Bliard, A. Merabtime,
433 M. Lachi, O. Douzane, et al., Characterization of a hemp-based agro-material: Influence of starch ratio and hemp
434 shive size on physical, mechanical, and hygrothermal properties, *Energy and Buildings* 153 (2017) 501–512.
- 435 [8] F. Iucolano, B. Liguori, P. Aprea, D. Caputo, Thermo-mechanical behaviour of hemp fibers-reinforced gypsum
436 plasters, *Construction and Building Materials* 185 (2018) 256–263.

- 437 [9] E. Sassoni, S. Manzi, A. Motori, M. Montecchi, M. Canti, Novel sustainable hemp-based composites for applica-
438 tion in the building industry: Physical, thermal and mechanical characterization, *Energy and Buildings* 77 (2014)
439 219–226.
- 440 [10] G. Costantine, C. Maalouf, T. Moussa, G. Polidori, Experimental and numerical investigations of thermal perfor-
441 mance of a hemp lime external building insulation, *Building and Environment* 131 (2018) 140–153.
- 442 [11] B. Moujalled, Y. A. Oumeziane, S. Moissette, M. Bart, C. Lanos, D. Samri, Experimental and numerical evalu-
443 ation of the hygrothermal performance of a hemp lime concrete building: A long term case study, *Building and*
444 *Environment* 136 (2018) 11–27.
- 445 [12] A. Santoni, P. Bonfiglio, F. Mollica, P. Fausti, F. Pompoli, V. Mazzanti, Vibro-acoustic optimisation of wood plastic
446 composite systems, *Construction and Building Materials* 174 (2018) 730–740.
- 447 [13] F. Asdrubali, S. Schiavoni, K. Horoshenkov, A review of sustainable materials for acoustic applications, *Building*
448 *Acoustics* 19 (4) (2012) 283–311.
- 449 [14] F. Asdrubali, F. D’Alessandro, S. Schiavoni, A review of unconventional sustainable building insulation materials,
450 *Sustainable Materials and Technologies* 4 (2015) 1–17.
- 451 [15] U. Berardi, G. Iannace, Predicting the sound absorption of natural materials: Best-fit inverse laws for the acoustic
452 impedance and the propagation constant, *Applied Acoustics* 115 (2017) 131–138.
- 453 [16] A. Santoni, P. Bonfiglio, J. L. Davy, P. Fausti, F. Pompoli, L. Pagnoncelli, Sound transmission loss of ETICS
454 cladding systems considering the structure-borne transmission via the mechanical fixings: Numerical prediction
455 model and experimental evaluation, *Applied acoustics* 122 (2017) 88–97.
- 456 [17] S. Secchi, G. Cellai, P. Fausti, A. Santoni, N. Zuccherini Martello, Sound transmission between rooms with curtain
457 wall façades: a case study, *Building Acoustics* 22 (3-4) (2015) 193–207.
- 458 [18] N. Zuccherini Martello, F. Aletta, P. Fausti, J. Kang, S. Secchi, A psychoacoustic investigation on the effect of
459 external shading devices on building façades, *Applied Sciences* 6 (12) (2016) 429.
- 460 [19] M. Delany, E. Bazley, Acoustical properties of fibrous absorbent materials, *Applied acoustics* 3 (2) (1970) 105–116.
- 461 [20] Y. Miki, Acoustical properties of porous materials-modifications of delany-bazley models, *Journal of the Acoustical*
462 *Society of Japan* (E) 11 (1) (1990) 19–24.
- 463 [21] D. L. Johnson, J. Koplik, R. Dashen, Theory of dynamic permeability and tortuosity in fluid-saturated porous
464 media, *Journal of fluid mechanics* 176 (1987) 379–402.
- 465 [22] Y. Champoux, J.-F. Allard, Dynamic tortuosity and bulk modulus in air-saturated porous media, *Journal of applied*
466 *physics* 70 (4) (1991) 1975–1979.
- 467 [23] D. Lafarge, P. Lemarinier, J. F. Allard, V. Tarnow, Dynamic compressibility of air in porous structures at audible
468 frequencies, *The Journal of the Acoustical Society of America* 102 (4) (1997) 1995–2006.
- 469 [24] M. Garai, F. Pompoli, A simple empirical model of polyester fibre materials for acoustical applications, *Applied*
470 *Acoustics* 66 (12) (2005) 1383–1398.
- 471 [25] U. Berardi, G. Iannace, Acoustic characterization of natural fibers for sound absorption applications, *Building and*
472 *Environment* 94 (2015) 840–852.
- 473 [26] C. Piégay, P. Gle, E. Gourdon, E. Gourlay, S. Marceau, Acoustical model of vegetal wools including two types of
474 fibers, *Applied Acoustics* 129 (2018) 36–46.
- 475 [27] H. Mamtaz, M. H. Fouladi, M. Al-Atabi, S. Narayana Namasivayam, Acoustic absorption of natural fiber compos-
476 ites, *Journal of Engineering* 2016.
- 477 [28] Z. Lim, A. Putra, M. Nor, M. Yaakob, Sound absorption performance of natural kenaf fibres, *Applied Acoustics*
478 130 (2018) 107–114.
- 479 [29] V. Tarnow, Measurement of sound propagation in glass wool, *The Journal of the Acoustical Society of America*
480 97 (4) (1995) 2272–2281.
- 481 [30] P. Shravage, P. Bonfiglio, F. Pompoli, et al., Hybrid inversion technique for predicting geometrical parameters of
482 porous materials, *Journal of the Acoustical Society of America* 123 (5) (2008) 3284.
- 483 [31] B. Castagnede, A. Aknine, B. Brouard, V. Tarnow, Effects of compression on the sound absorption of fibrous
484 materials, *Applied Acoustics* 61 (2) (2000) 173–182.
- 485 [32] M. Horne, Bast fibres: hemp cultivation and production, in: *Handbook of Natural Fibres: Types, Properties and*
486 *Factors Affecting Breeding and Cultivation*, Elsevier, 2012, pp. 114–145.
- 487 [33] L. Y. Mwaikambo, M. P. Ansell, Chemical modification of hemp, sisal, jute, and kapok fibers by alkalization,
488 *Journal of applied polymer science* 84 (12) (2002) 2222–2234.
- 489 [34] J. Zhang, H. Zhang, J. Zhang, Effect of alkali treatment on the quality of hemp fiber., *Journal of Engineered Fabrics*
490 *& Fibers (JEFF)* 9 (2).
- 491 [35] V. Mazzanti, A. Bonanno, F. Mollica, G. Filippone, Influence of alkaline treatment on hemp fibers filled poly (lactic
492 acid), in: *AIP Conference Proceedings*, Vol. 1981, AIP Publishing, 2018, p. 020016.
- 493 [36] H. T. Luu, R. Panneton, C. Perrot, Effective fiber diameter for modeling the acoustic properties of polydisperse
494 fiber networks, *The Journal of the Acoustical Society of America* 141 (2) (2017) EL96–EL101.
- 495 [37] V. Tarnow, Airflow resistivity of models of fibrous acoustic materials, *The Journal of the Acoustical Society of*

- 496 America 100 (6) (1996) 3706–3713.
- 497 [38] EN 29053 – Acoustics - Materials for acoustical applications - Determination of airflow resistance, Standard,
498 European Committee for Standardisation, Brussels, Belgium (1993).
- 499 [39] ISO 10534-2 – Acoustics – Determination of sound absorption coefficient and impedance in impedance tubes –
500 Part 2: Transfer-function method., Standard, International Organization for Standardization, Geneva, CH (1998).
- 501 [40] H. T. Luu, C. Perrot, R. Panneton, Influence of porosity, fiber radius and fiber orientation on the transport and
502 acoustic properties of random fiber structures, *Acta Acustica united with Acustica* 103 (6) (2017) 1050–1063.
- 503 [41] P. Bonfiglio, F. Pompoli, Inversion problems for determining physical parameters of porous materials: Overview
504 and comparison between different methods, *Acta Acustica united with Acustica* 99 (3) (2013) 341–351.
- 505 [42] B. Campolina, N. Dauchez, N. Atalla, O. Doutres, Effect of porous material compression on the sound transmission
506 of a covered single leaf panel, *Applied Acoustics* 73 (8) (2012) 791–797.
- 507 [43] L. Lei, N. Dauchez, J. Chazot, Prediction of the six parameters of an equivalent fluid model for thermocompressed
508 glass wools and melamine foam, *Applied Acoustics* 139 (2018) 44–56.
- 509 [44] K. Hirose, H. Nakagawa, Formulae for predicting non-acoustical parameters of deformed fibrous porous mate-
510 rials, *The Journal of the Acoustical Society of America* 141 (6) (2017) 4301–4313.
- 511 [45] P. Bonfiglio, F. Pompoli, et al., Frequency dependent tortuosity measurement by means of ultrasonic tests, in:
512 Proceeding of the 14th International Congress on Sound and Vibration, Vol. 14, ICSV, Cairns, Australia, 2007, pp.
513 1–6.

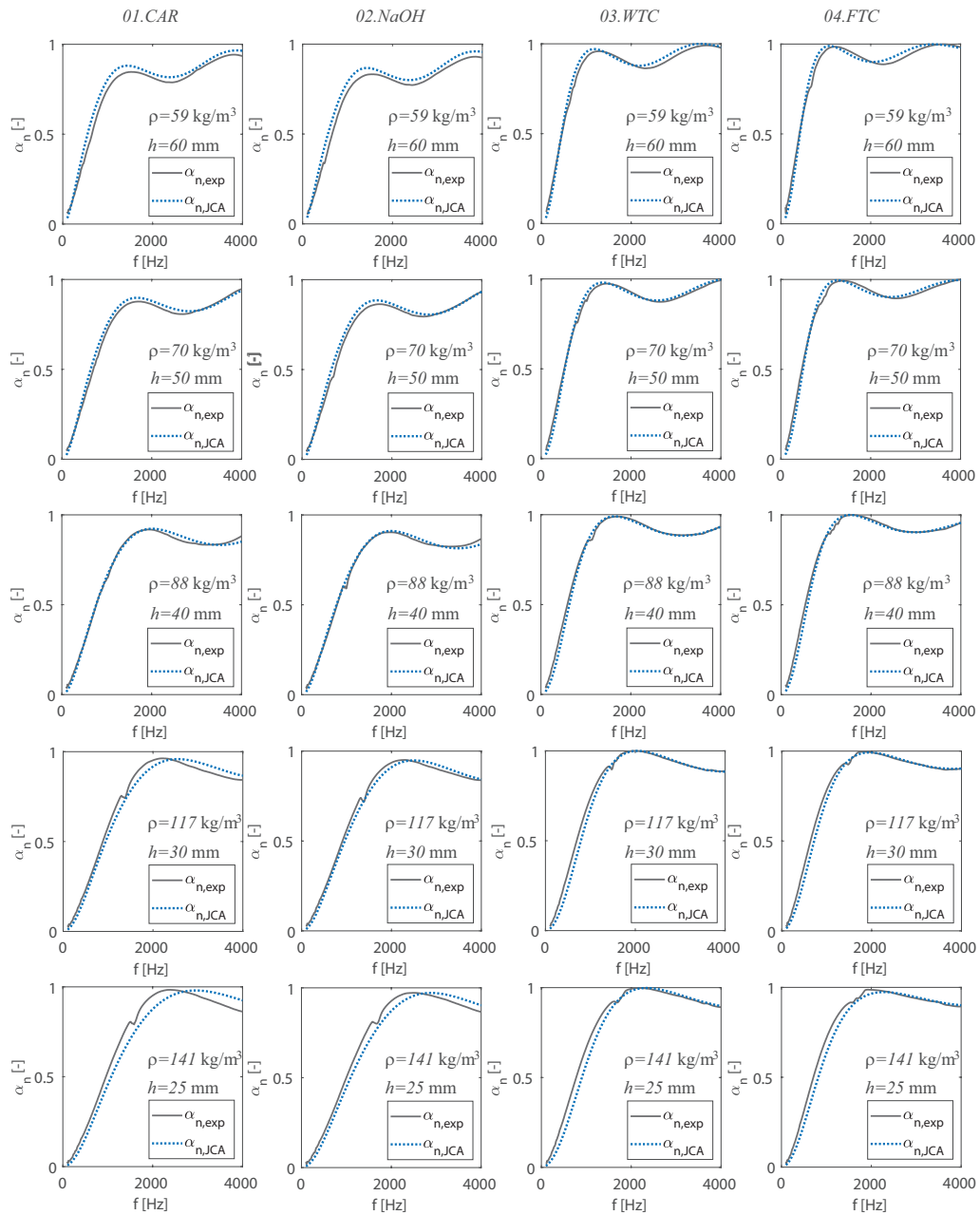


Figure 9: Comparison between numerical and experimental normal incidence sound absorption coefficient of the hemp fibrous material at each stage of the manufacturing process, for different compression rate.